

CNP research day

Diffuse Boosted Cosmic Neutrino Background

Xiaolin Qi

Department of Physics

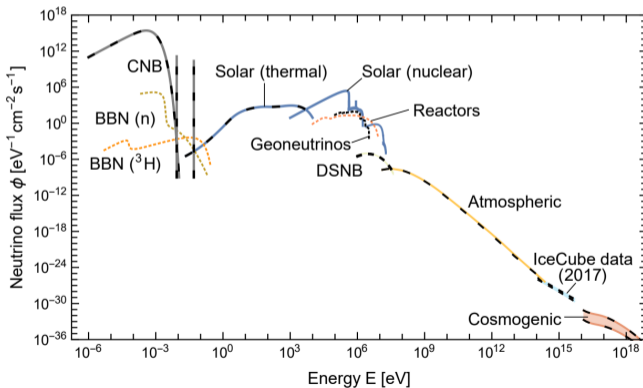
xiaolinq76@vt.edu

Collaborator: Gonzalo Herrera, Shunsaku Horiuchi

arXiv:2405.14946

Motivation

Figure: Vitagliano, Tamborra, and Raffelt 2020



Motivation

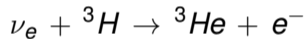
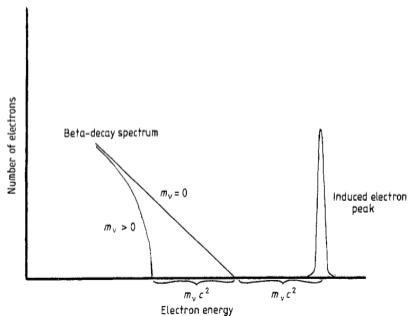
- $C\nu B$ (relic neutrinos) decoupled from matter at the early universe (~ 1 second old).
- As the universe expanded, it further cooled down ($\sim 1.95K$).

$C\nu B$ is extremely difficult to detect due to its low energy!

Previous work

- 'overdensity'
- **KATRIN** placed an upper limit on neutrino overdensity Aker et al. 2022: 10^{11}

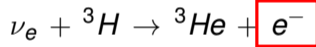
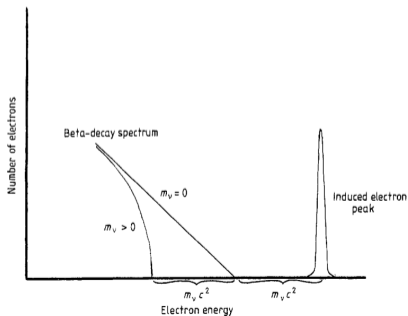
Figure: Irvine and Humphreys 1983



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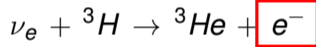
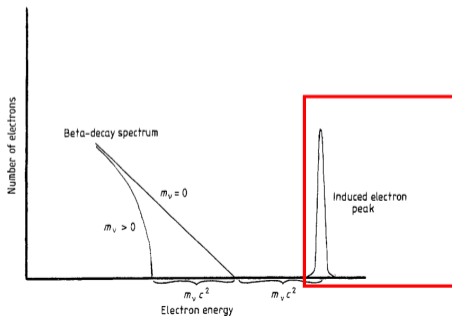
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Is it possible for $C\nu B$ to be boosted to higher energies?

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YES!

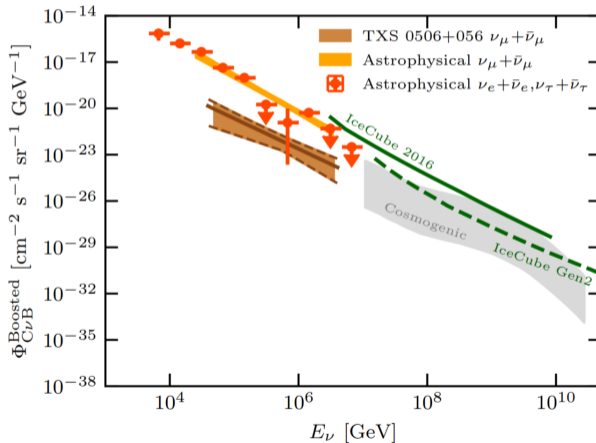
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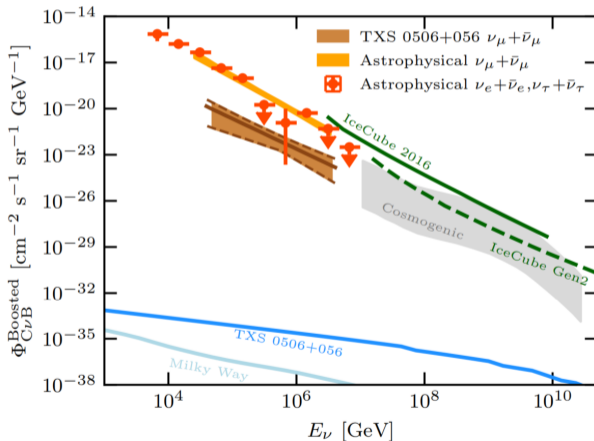
YES!

Cosmic rays scattering off relic neutrinos \rightarrow boost!

Previous work



Previous work

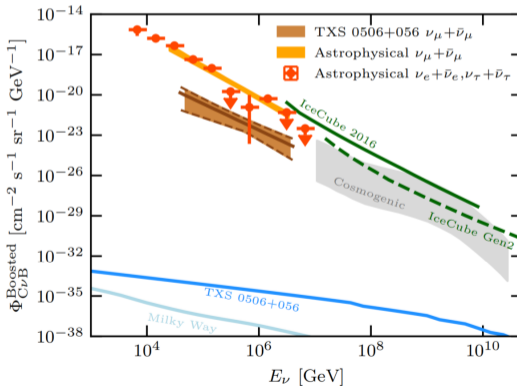


Previous work

Figure: Císcar-Monsalvatje, Herrera, and Shoemaker 2024

Milky Way $\rightarrow \sim 10^{13}$

TXS 0506+056 $\rightarrow \sim 10^{10}$



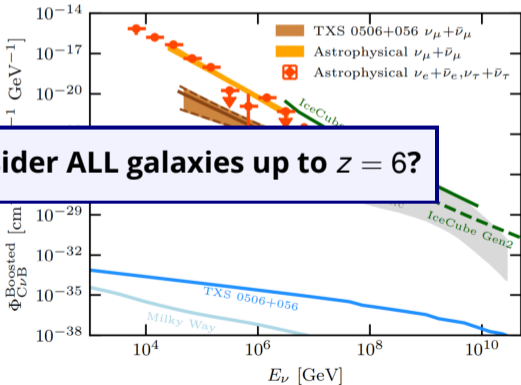
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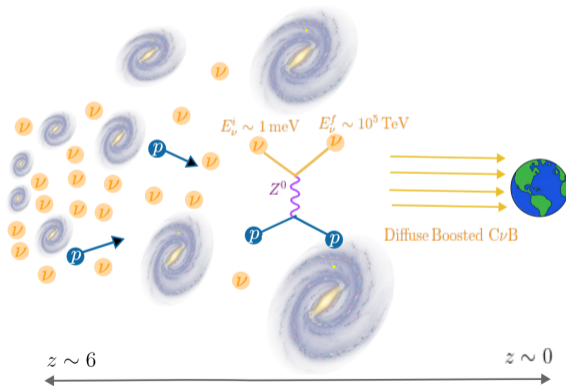
TXS 0506+056 \rightarrow

What if we consider ALL galaxies up to $z = 6$?



Schematic plot

Figure: Herrera, Horiuchi, and Qi 2024



Double integral

$$\frac{d\phi_\nu}{dT_\nu} = \int_{z_{\min}}^{z_{\max}} dz \frac{c}{H_0} \frac{1}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} f_i(z) n_\nu (1+z)^3 \int_0^\infty dT_p \sigma_{p\nu}(T_p) \frac{d\phi_p}{dT_p} \frac{1}{T_\nu^{\max}(T_p)} \Theta [T_\nu^{\max}(T_p) - T_\nu(1+z)]$$

Assumptions

CR composition

consider only protons

scattering process

consider only neutral current interactions

Double integral

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Cosmological terms

Double integral

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proton-neutrino scattering cross section

Double integral

$$\frac{d\phi_\nu}{dT_\nu} = \int_{z_{\min}}^{z_{\max}} dz \frac{c}{H_0} \frac{1}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} f_i(z) n_\nu (1+z)^3 \int_0^\infty dT_p \sigma_{p\nu}(T_p) \boxed{\frac{d\phi_p}{dT_p}} \frac{1}{T_\nu^{\max}(T_p)} \Theta [T_\nu^{\max}(T_p) - T_\nu(1+z)]$$

Cosmic Ray spectrum in the Milky Way

Double integral

$$\frac{d\phi_\nu}{dT_\nu} = \int_{z_{\min}}^{z_{\max}} dz \frac{c}{H_0} \frac{1}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} f_i(z) n_\nu (1+z)^3 \int_0^\infty dT_p \sigma_{p\nu}(T_p) \frac{d\phi_p}{dT_p} \frac{1}{T_\nu^{\max}(T_p)} \Theta [T_\nu^{\max}(T_p) - T_\nu(1+z)]$$

maximal energy transferred to a neutrino in one scattering

Double integral

$$\frac{d\phi_\nu}{dT_\nu} = \int_{z_{\min}}^{z_{\max}} dz \frac{c}{H_0} \frac{1}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} \boxed{f_i(z)} n_\nu (1+z)^3 \int_0^\infty dT_p \sigma_{p\nu}(T_p) \frac{d\phi_p}{dT_p} \frac{1}{T_\nu^{\max}(T_p)} \Theta [T_\nu^{\max}(T_p) - T_\nu(1+z)]$$

redshift evolution of Cosmic Ray flux

▶ Go to detailed explanation

Double integral

$$\frac{d\phi_\nu}{dT_\nu} = \int_{z_{\min}}^{z_{\max}} dz \frac{c}{H_0} \frac{1}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} f_i(z) \boxed{n_\nu(1+z)^3} \int_0^\infty dT_p \sigma_{p\nu}(T_p) \frac{d\phi_p}{dT_p} \frac{1}{T_\nu^{\max}(T_p)} \Theta [T_\nu^{\max}(T_p) - T_\nu(1+z)]$$

neutrino density at redshift z

▶ Check the results!

Redshift evolution of CR

◀ Return to other terms

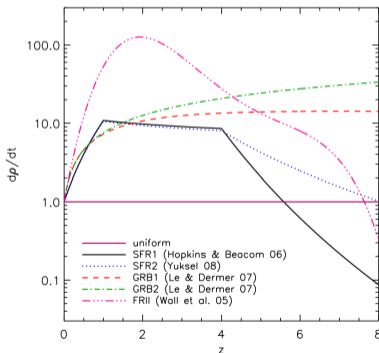
Redshift evolution of Cosmic Ray flux

$$f_i(z) = \frac{N_i(z)}{N_i(z_{min})}$$

$N_i(z)$: distribution of:

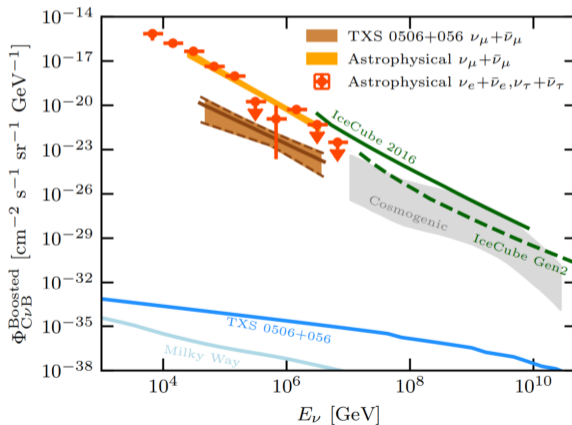
- SFR
- GRB
- QSO

Figure: Kotera, Allard, and Olinto 2010



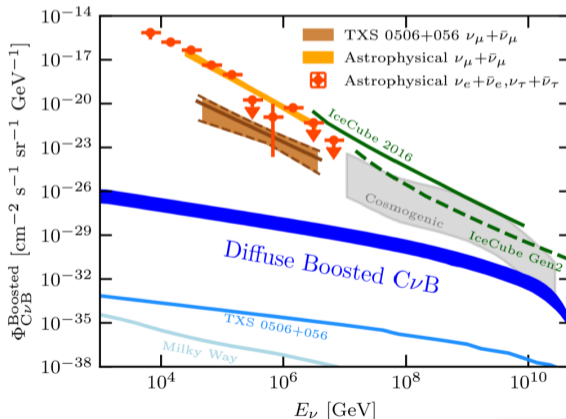
Result

Figure: Císcar-Monsalvatje, Herrera, and Shoemaker 2024



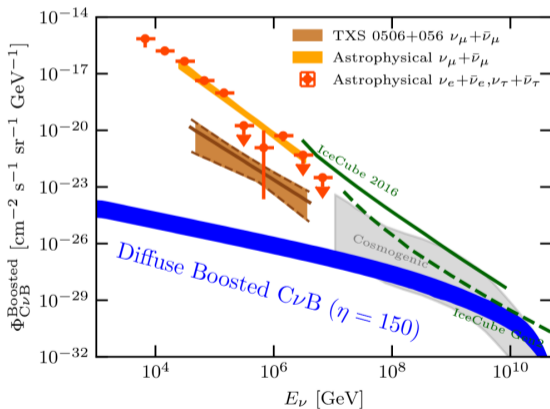
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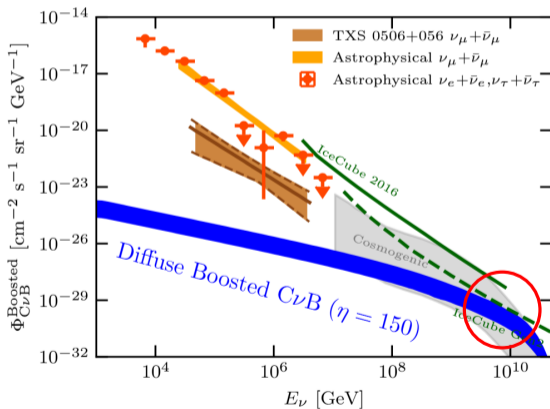
Result - incorporating overdensity

Figure: Herrera, Horiuchi, and Qi 2024



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Conclusion and Outlook

Boosted flux

Relic neutrino flux is remarkably boosted.

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Stringent constraint on overdensity

With a sensitivity increasing of only ~ 150 compared to IceCube Gen2, we have a hope to directly detect $C\nu B$.

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Outlook


Further improvement is possible by taking into account:

- heavy nuclei
- deep inelastic regime
- charged current interaction
→ a program in progress with Shunsaku Horiuchi, Ian Shoemaker, Gonzalo Herrela.

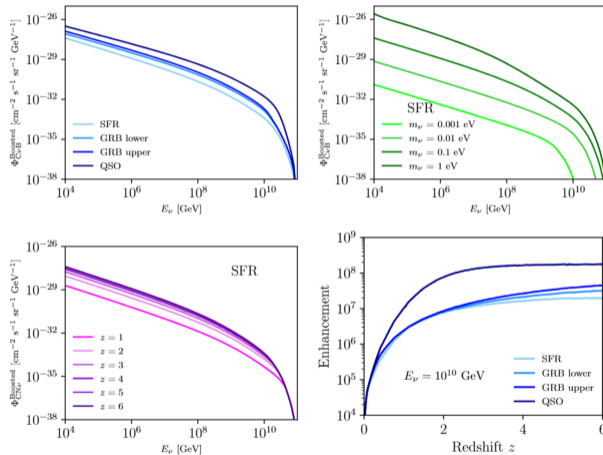
Questions?

References I

-  Aker, M et al. (2022). “New constraint on the local relic neutrino background overdensity with the first KATRIN data runs”. In: *Physical Review Letters* 129.1, p. 011806.
-  Císcar-Monsalvatje, Mar, Gonzalo Herrera, and Ian M Shoemaker (2024). “Upper limits on the cosmic neutrino background from cosmic rays”. In: *Physical Review D* 110.6, p. 063036.
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-  Irvine, JM and R Humphreys (1983). “Neutrino masses and the cosmic-neutrino background”. In: *Journal of Physics G: Nuclear Physics* 9.7, p. 847.
-  Kotera, Kumiko, Denis Allard, and Angela V Olinto (2010). “Cosmogenic neutrinos: parameter space and detectability from PeV to ZeV”. In: *Journal of Cosmology and Astroparticle Physics* 2010.10, p. 013.

-  Vitagliano, Edoardo, Irene Tamborra, and Georg Raffelt (2020). "Grand unified neutrino spectrum at Earth: Sources and spectral components". In: *Reviews of Modern Physics* 92.4, p. 045006.

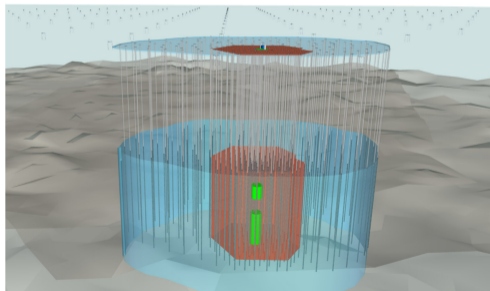
More Results



IceCube Gen2

- next-generation South Pole neutrino observatory
- "A core detector will be the IceCube-Gen2 optical array, with a size eight times larger than the current IceCube and optimized to detect neutrinos with energies ranging from hundreds of TeV to tens or even a few hundreds of PeV."

Figure: <https://icecube-gen2.wisc.edu/about/icecube-gen2/>



A 3D rendering of the planned IceCube-Gen2 extension. IceCube-Gen2 encompasses three new arrays—in-ice optical, surface, and extensive radio—that expand the capabilities of the current IceCube Neutrino Observatory. Image credit: IceCube Collaboration