

# High energy photon polarimetry

B. Wojtsekhowski, JLab

motivated by the DVCS proposal

under development by Maxime Defurne, Saclay

## Photon polarimetry

- Neutral pion:  $\pi^0$  decay to  $\gamma + \gamma$ , polarizations
- Polarimetry for 100s MeV photons beam,  $e^+e^-$
- Astrophysics: Photon of low energy, Compton
- Photon Beam: goniometer – Hall D,  $e^+e^-$
- Astrophysics: Photons of high energy,  $\gamma \Rightarrow e^+e^-$
- DVCS: large solid angle, several GeV photons

# Pion decay polarimetry

## Possible Experimental Determination of Whether the Neutral Meson is Scalar or Pseudoscalar

C. N. YANG

*Institute for Advanced Study, Princeton, New Jersey*

January 16, 1950

NEUTRAL mesons with a mass of  $\sim 300$  Mev which decay into two photons have been reported.<sup>1</sup> It can be proved<sup>2,3</sup> on general grounds of rotation-inversion invariance that a particle which dematerializes into two photons cannot have spin 1. Furthermore,<sup>3</sup> if the neutral meson is a scalar meson, the two photons would always have parallel planes of polarization, while if the neutral meson is a pseudoscalar meson, the two photons would always have perpendicular planes of polarization. One is

PHYSICAL REVIEW

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APRIL 1, 1962

## Measurement of Linear Photon Polarization by Pair Production\*

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(Received November 7, 1961)

# Photon beam polarimetry

## NEW METHOD FOR MEASUREMENT OF $\gamma$ -RAY POLARIZATION BY DETECTION OF ANGULAR CORRELATION IN PAIR PRODUCTION

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Received 5 May 1972

$$R(\eta_-, \eta_+) = \frac{\frac{\partial^3 \sigma_{\parallel}}{\partial E_- \partial \eta_- \partial \eta_+} - \frac{\partial^3 \sigma_{\perp}}{\partial E_- \partial \eta_- \partial \eta_+}}{\frac{\partial^3 \sigma_{\parallel}}{\partial E_- \partial \eta_- \partial \eta_+} + \frac{\partial^3 \sigma_{\perp}}{\partial E_- \partial \eta_- \partial \eta_+}},$$

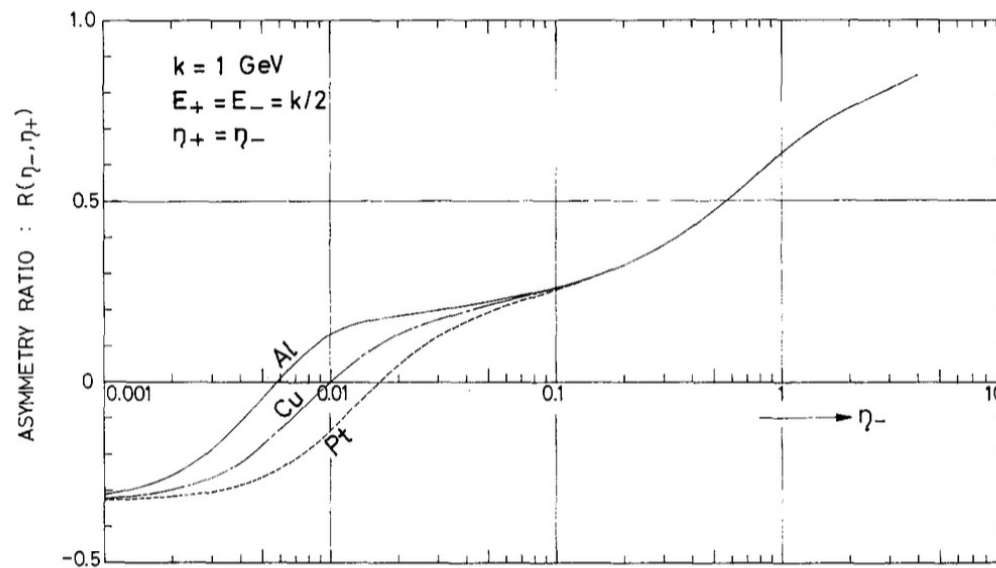
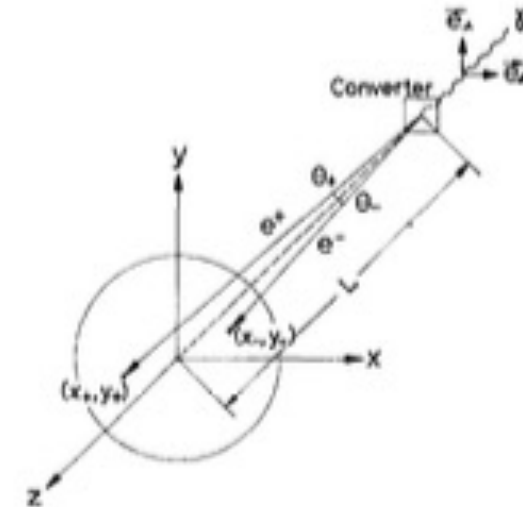
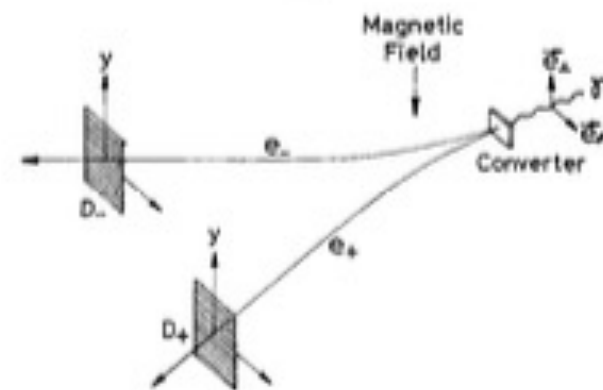


Fig. 6a.  $R(\eta_-, \eta_+)$  calculated on the line  $\eta_+ = \eta_-$  in different materials.



(a)



(b)

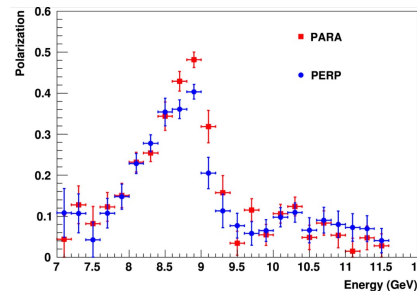
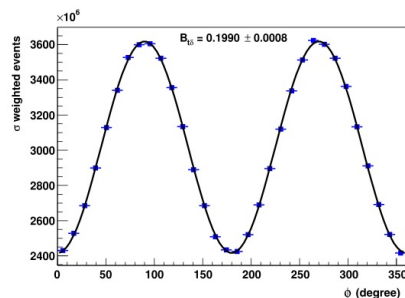
Fig. 1. (a) Coordinates used in the calculation. (b) A sketch of an example of possible experimental arrangements.  $D_{\pm}$  are detectors such as wire chambers or hodoscopes.

# Photon beam polarimetry – triplet

## Design and construction of a high-energy photon polarimeter

M. Dugger<sup>a,\*</sup>, B.G. Ritchie<sup>a</sup>, N. Sparks<sup>a</sup>, K. Moriya<sup>a</sup>, R.J. Tucker<sup>a</sup>, R.J. Lee<sup>a</sup>, B.N. Thorpe<sup>a</sup>,  
 T. Hodges<sup>a</sup>, F.J. Barbosa<sup>b</sup>, N. Sandoval<sup>b</sup>, R.T. Jones<sup>c</sup>

Nuclear Inst. and Methods in Physics Research, A 867 (2017) 115–127



**Triplet production:**  $e^+e^-$  pair  
 produced from an electron  
 Cross section  $\Rightarrow \sim Z$   
 It is a large angle recoil, but  
**low energy electron**

After all cuts have been applied, the detector accepts approximately 2% of the total triplet cross section. The design goal for the GlueX photon beam is to have a polarization of 0.4 and rate of  $10^8$  photons per second within the coherent peak energy range of 8.4–9.0 GeV. Given such a beam incident on a 75  $\mu\text{m}$  thick beryllium converter, the expected TPOL rate surviving all cuts is 75 Hz (and perhaps as high as 10 kHz without cuts for the full incident photon-energy range), leading to an approximate statistical uncertainty in the TPOL determination of polarization of 0.014, after a single hour of production running.

$10^{11}$  photons  $\Rightarrow$  0.014 accuracy

**or interpolate to  $10^9$  for 14%**

# High energy photon beam polarimetry

A pair polarimeter for linearly polarized high-energy photons

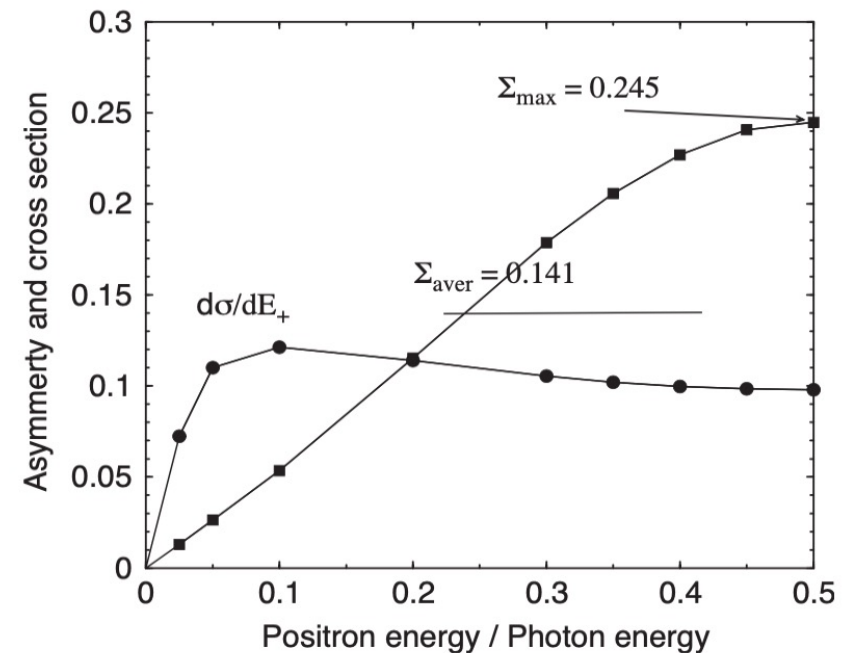
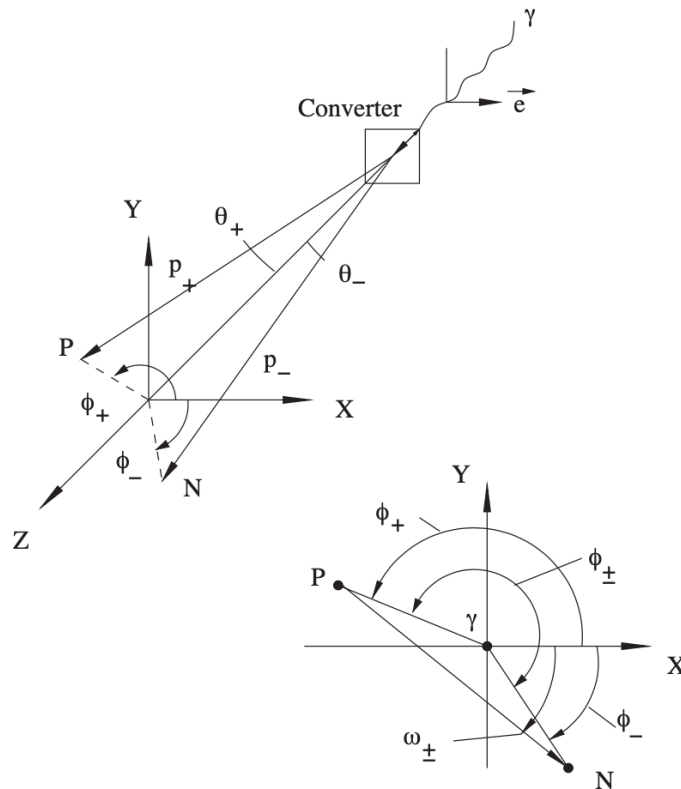
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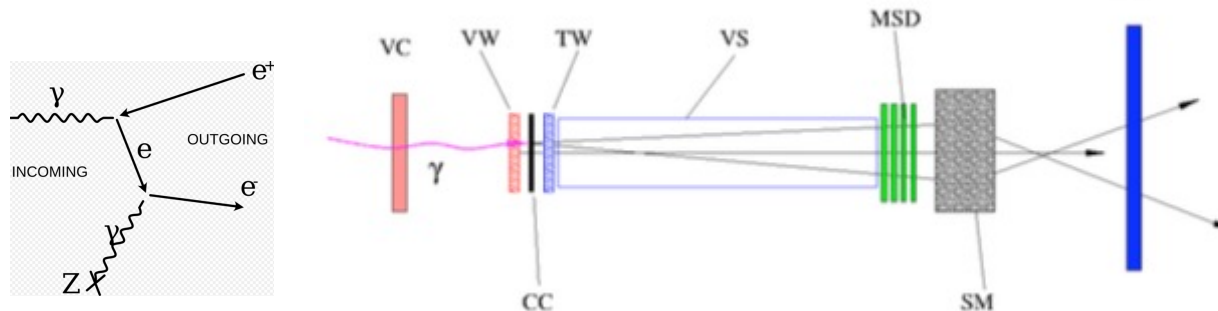


# High energy photon beam polarimetry

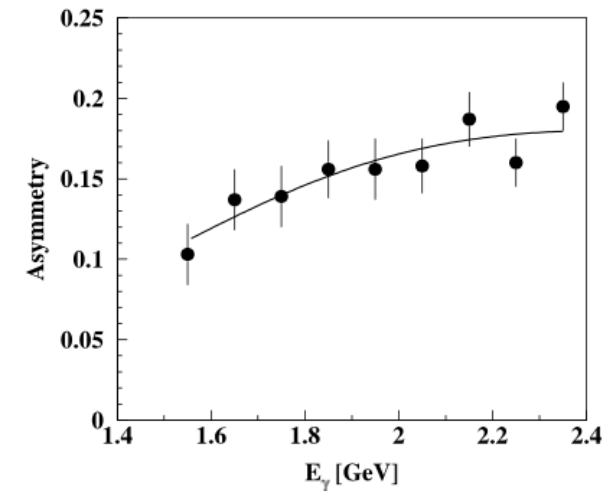
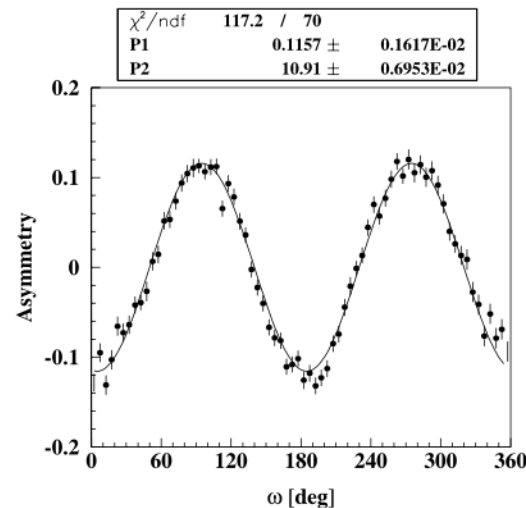
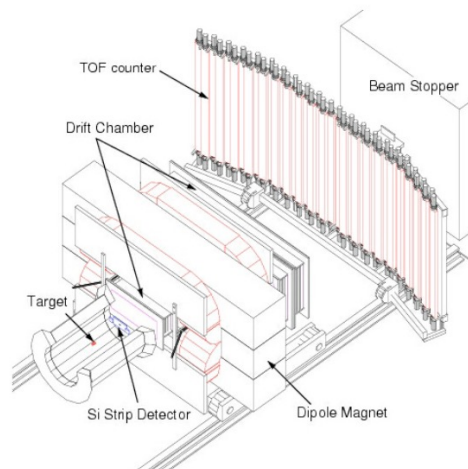
Eur Phys J A (2004) **19**, s01, 275–278

## A pair polarimeter for linearly polarized high energy photons

C. de Jager<sup>1</sup>, B. Wojtsekhowski<sup>1</sup>, D. Tedeschi<sup>2</sup>, B. Vlahovic<sup>1,3</sup>, D. Abbott<sup>1</sup>, J. Asai<sup>4</sup>, G. Feldman<sup>5</sup>, T. Hotta<sup>6</sup>, M. Khadaker<sup>7</sup>, H. Kohri<sup>6</sup>, T. Matsumura<sup>6</sup>, T. Mibe<sup>6</sup>, T. Nakano<sup>6</sup>, V. Nelyubin<sup>1,2</sup>, G. Orielly<sup>5</sup>, A. Rudge<sup>8</sup>, P. Weilhammer<sup>8</sup>, M. Wood<sup>2</sup>, T. Yorita<sup>6</sup>, and R. Zegers<sup>9</sup>



$$E_\gamma = 1.6 - 2.4 \text{ GeV}$$



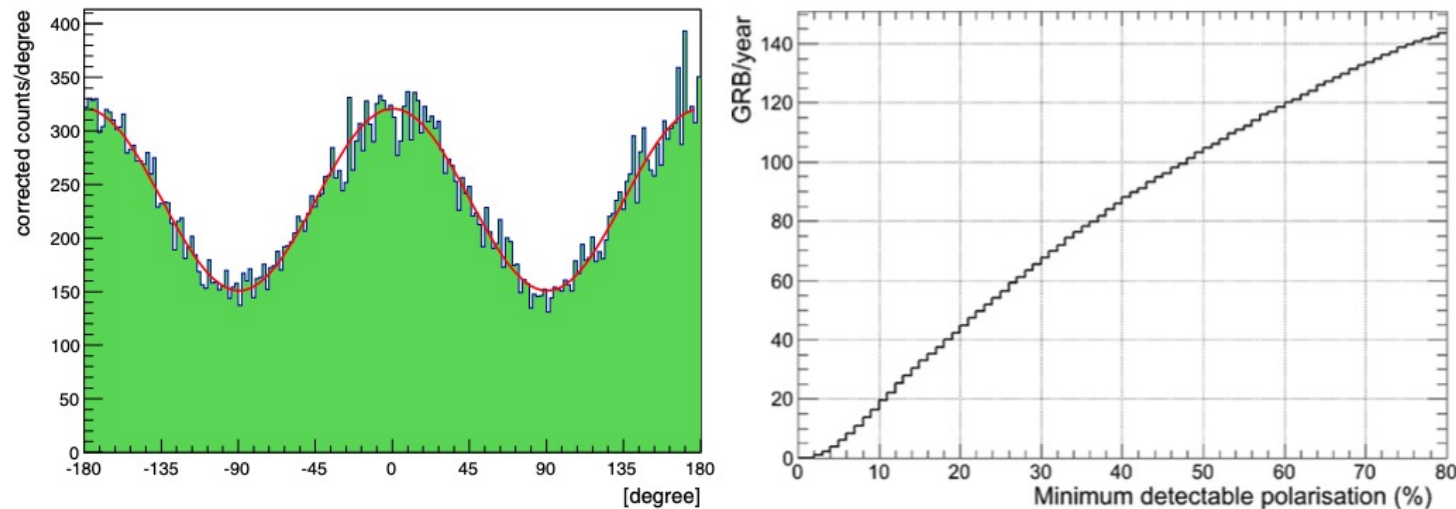


# Astrophysics gamma polarimetry

$$E_\gamma < 2 \text{ MeV}$$

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Exp Astron (2017) 44:25–82



**Fig. 20** *Left panel*—e-ASTROGAM polarization response (polarigramme) in the 0.2–2 MeV range for a 100% polarized, 10 mCrab-like source observed on axis for  $10^6$  s. The corresponding modulation is  $\mu_{100} = 0.36$ . *Right panel*—Cumulative number of GRBs to be detected by e-ASTROGAM as a function of the minimum detectable polarization at the 99% confidence level

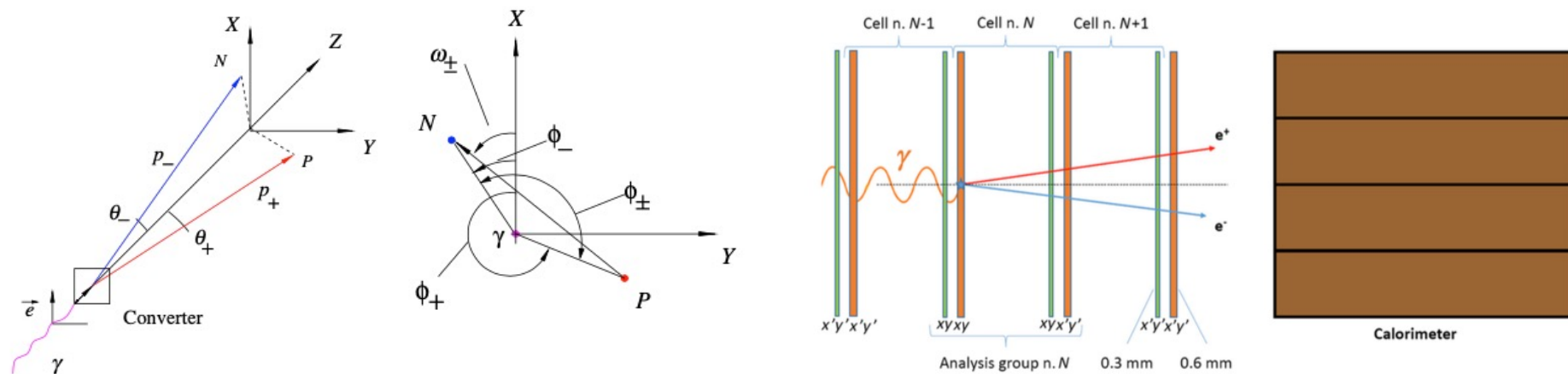


# Astrophysics gamma polarimetry: Si MSD

Journal of Astronomical Telescopes, Instruments, and Systems 4(1), 011006 (Jan–Mar 2018)

## High-energy photon polarimeter for astrophysics

Maxim Eingorn,<sup>a</sup> Lakma Fernando,<sup>a</sup> Branislav Vlahovic,<sup>a,\*</sup> Cosmin Ilie,<sup>b,c</sup> Bogdan Wojtsekhowski,<sup>d,\*</sup>  
 Guido Maria Urciuoli,<sup>e</sup> Fulvio De Persio,<sup>e</sup> Franco Meddi,<sup>e,f</sup> and Vladimir Nelyubin<sup>g</sup>



For high photon energy it is better to use MSD due to its coordinate resolution

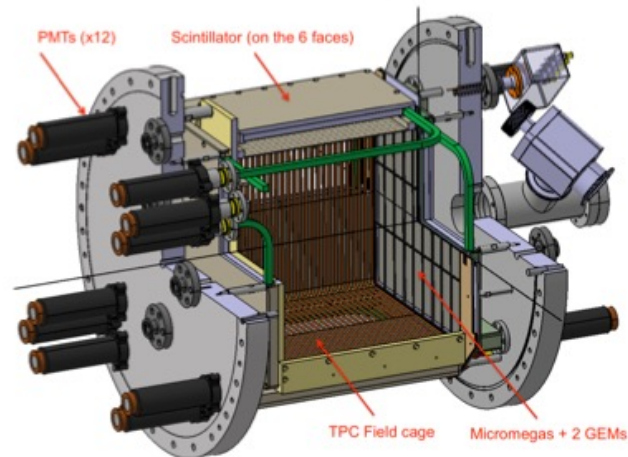
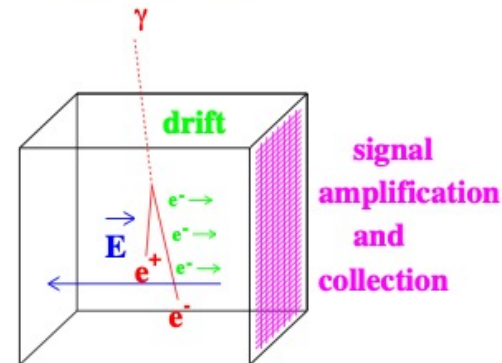
# Astrophysics gamma polarimeter: TPC

D. Bernard et al.

Future Space-based Gamma-ray Observatories Workshop, Goddard Space Flight Center, 2016

## *HARPO: the Demonstrator*

- Time Projection Chamber (TPC)
- $(30\text{cm})^3$  cubic TPC
- Up to 5 bar.
- Micromegas + GEM gas amplification
- Collection on  $x, y$  strips, pitch 1 mm.
- AFTER chip digitization, up to 100 MHz.
- Scintillator / WLS / PMT based trigger



NIM A 695 (2012) 71, NIM A 718 (2013) 395

# Physics: DVCS photon polarimetry

Gluon tomography in nucleons by  $\gamma$ -polarimetry

Maxime Defurne

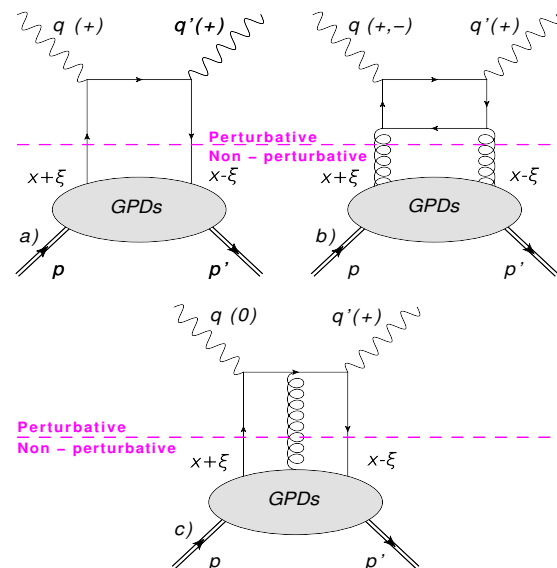
February 11, 2023

We propose to measure the degree of linear polarization of a photon produced by deeply virtual Compton scattering, never considered but rich in terms of new information about the proton structure. Indeed, the photon polarization is rigorously and straightforwardly related to the gluon transversity GPDs, still completely unknown today.

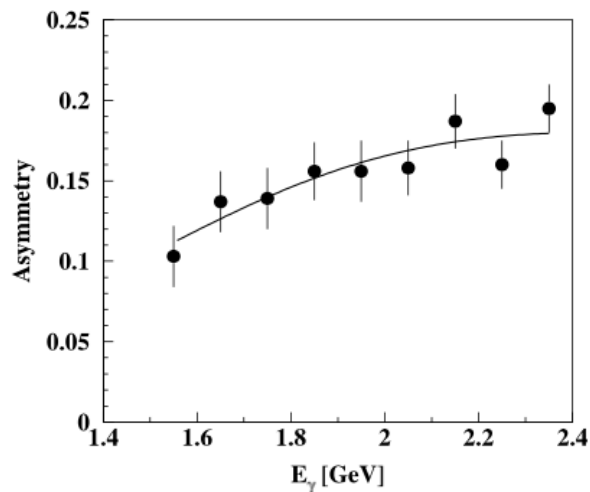
## 1 From gluon transversity GPDs to $\gamma$ -polarization

### 1.1 A *simple* answer to: “Why $\gamma$ -polarization?”

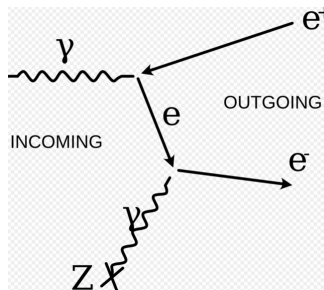
In deeply virtual Compton scattering (DVCS), a highly virtual photon is absorbed by the nucleon and the latter subsequently emits a high-energy real photon. At the partonic level, the process can be described by the diagrams presented in Figure 1. The process can be described in terms of helicity



# DVCS photon polarimetry



For photon energies above 100 MeV, the silicon MSD is a preferable option because of the limit on the apparatus's total length. Indeed, considering a 300-cm total length and an energy of 500 MeV, the number of cells is 5 for the drift chamber option, 6 for the TPC/micromega, and 46 for the MSD option. At the same time, the total amount of matter in the polarimeter should be limited to one radiation length or less because of significant absorption of the incident photons, which will reduce the average efficiency per cell. For example, in the MSD option 54% absorption will occur with 46 cells



$$L_{\text{pair}} = 9/7 X_0$$

$$\Theta_{\text{ms}} \sim 1/p \text{ for } 1\% X_0 \text{ detector}$$

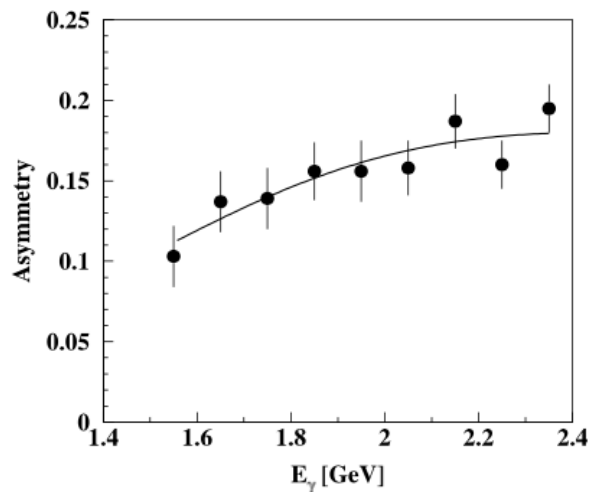
	$X_0$
Scint	400 mm
Si	94 mm
Cu	14 mm

1 GeV, 100  $\mu\text{m}$   $\Rightarrow$  10 cm  
 spacing between layers

$$\theta_{\text{plane}}^{\text{rms}} = \frac{13.6 \text{ MeV}}{p\beta} z \sqrt{x/X_0} [1 + 0.038 \ln(x/X_0)]$$

The key is a two-track  
 space resolution  $\Rightarrow$  MSD

## DVCS photon polarimetry



For photon energies above 100 MeV, the silicon MSD is a preferable option because of the limit on the apparatus's total length. Indeed, considering a 300-cm total length and an energy of 500 MeV, the number of cells is 5 for the drift chamber option, 6 for the TPC/micromegas, and 46 for the MSD option. At the same time, the total amount of matter in the polarimeter should be limited to one radiation length or less because of significant absorption of the incident photons, which will reduce the average efficiency per cell. For example, in the MSD option 54% absorption will occur with 46 cells

- Using projected asymmetry  $A = 0.15$  and desirable accuracy of 10% the statistics of  $e^+e^-$  events of 4k is required.
- With detection efficiency of 1% per detector layer ( $1\%X_0 \sim 100 \mu\text{m Si}$ ) the required number of DVCS events is 0.4M.

Due detector resolution will be some reduction of  $A$ , so statistics of 1M is likely a lower limit.

With 46 layers in the statistics of 25k should be sufficient.



# DVCS photon polarimetry

## Scintillator fiber systems

### Scintillating fiber detectors for the HypHI project at GSI

D. Nakajima<sup>a,b,\*</sup>, B. Özel-Tashenov<sup>a,c,\*\*</sup>, S. Bianchin<sup>a</sup>, O. Borodina<sup>a,d</sup>, V. Bc  
 M. Kavatsyuk<sup>e</sup>, S. Minami<sup>a</sup>, C. Rappold<sup>a,f</sup>, T.R. Saito<sup>a,d</sup>, P. Achenbach<sup>d</sup>, S. P  
 T. Fukuda<sup>h</sup>, Y. Hayashi<sup>i</sup>, T. Hiraiwa<sup>i</sup>, J. Hoffmann<sup>a</sup>, K. Koch<sup>a</sup>, N. Kurz<sup>a</sup>, O. I  
 Y. Mizoi<sup>h</sup>, T. Mochizuki<sup>h</sup>, M. Moritsu<sup>i</sup>, T. Nagae<sup>i</sup>, L. Nungesser<sup>d</sup>, A. Okamui  
 A. Sakaguchi<sup>h</sup>, M. Sako<sup>i</sup>, C.J. Schmidt<sup>a</sup>, H. Sugimura<sup>i</sup>, K. Tanida<sup>i</sup>, M. Träger

A two-dimensional scintillation-based neutron detector with  
 wavelength-shifting fibers and incorporating an interpolation method

T. Nakamura<sup>a,\*</sup>, K. Toh<sup>a</sup>, T. Kawasaki<sup>a</sup>, M. Ebine<sup>b</sup>, A. Birumachi<sup>b</sup>, K. Sakasai<sup>a</sup>, K. Soyama<sup>a</sup>

### 3.3. Spatial resolution

Fig. 7 shows the spatial resolution measured while scanning the collimated beam over the detector. The spatial responses were fitted with a Gaussian function to extract the variance ( $\sigma$ ) for each incidence position. The spatial resolution, which was calculated as the full width at half maximum (FWHM) by  $2.35\sigma$ , was better when the neutron beam was incident on top or near the WLS fiber than when it was incident between the fibers. Measurements made over a distance of 10 mm revealed a periodicity in the spatial resolution of 2.5 mm for all of the MPC logics, which reflected the pitch of the WLS fibers. These observations were consistent with the spatial responses shown in Figs. 4 and 5.

The average FWHM spatial resolutions were  $3.3 \pm 0.3$ ,  $2.7 \pm 0.1$ , and  $2.5 \pm 0.1$  mm for standard-, half-, and quarter-pitch logics, respectively. The spatial resolution improved from 1.2- to 1.5-fold

noise. At a single photon threshold the SiPM suffers from the same problems as a GAPD. Current state of the art SiPMs have thermal noise rates<sup>6</sup> at  $\mathcal{O}(100 \text{ kHz})$ <sup>7</sup> at single photon level. This value depends on the temperature and the bias voltage. The typical pixel sizes of SiPMs are between  $25 \times 25 \mu\text{m}^2$  and  $100 \times 100 \mu\text{m}^2$ . One single device can cover active areas up to  $6 \times 6 \text{ mm}^2$  (fig. 3.9).

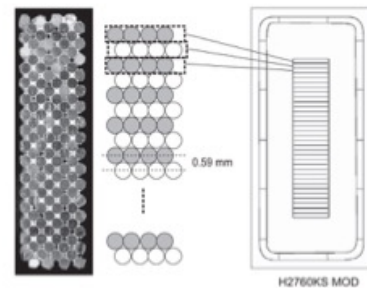
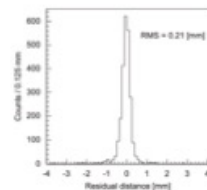


Fig. 2. (a) Cross-section of a 32-channel fiber bundle and (b) corresponding enlarged schematic drawing. The panel (c) shows a scheme of the surface of PMT H2760KS MOD.



Two-track coordinate resolution of the scintillator fiber-based system is not sufficiently good

# DVCS photon polarimetry

*J. Derré et al. / Nuclear Instruments and Methods in Physics Research A 459 (2001) 523–531*

527

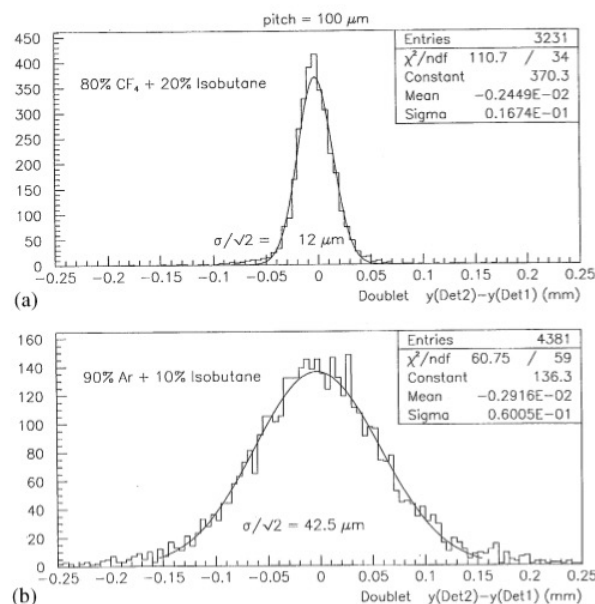


Fig. 3. The difference between the intercept of tracks in the two detectors of a Micromegas doublet for a gas mixture of (a) CF<sub>4</sub> and isobutane, (b) argon and isobutane. The fitted lines are gaussian distributions.

This observation is related to the behavior of the transverse diffusion in the gas. A similar effect has been observed [5] for the longitudinal diffusion.

## 4.2. Other gas mixtures

The spatial resolution of Micromegas has also been studied for other gas mixtures. For comparison with the above results, the corrected position difference distribution in a doublet operated with our standard gas mixture Ar–isobutane (90%,10%) is shown in Fig. 3(b). The measurement shows a clear degradation of the accuracy which is  $42.5 \mu\text{m}$  (standard deviation divided by  $\sqrt{2}$ ).

Table 1 summarizes all the results obtained with a doublet of detectors. For each gas mixture we give the measured resolution and the transverse diffusion coefficient. The transverse diffusion coefficient results from an adjustment in our simulation to reproduce the charge and cluster size distributions shown by the data.

The spatial resolution improves when the amount of quencher is increased in the gas mixture and the amount of CF<sub>4</sub> is large. The best result was

Layers of micromega chamber could be the best detector for the DVCS polarimeter due to high coordinate resolution also good two-track resolution.



## Detector setup for DVCS polarimetry

150 cm x 150 cm area coordinate detector  
25 layers in 500 cm depth – 25% efficiency  
EM calorimeter for event selection behind

Located 5 m from the target: 90 msr solid angle

70 msr SBS spectrometer as an electron arm

With 10  $\mu$ A beam on 15 cm long LH2 target