

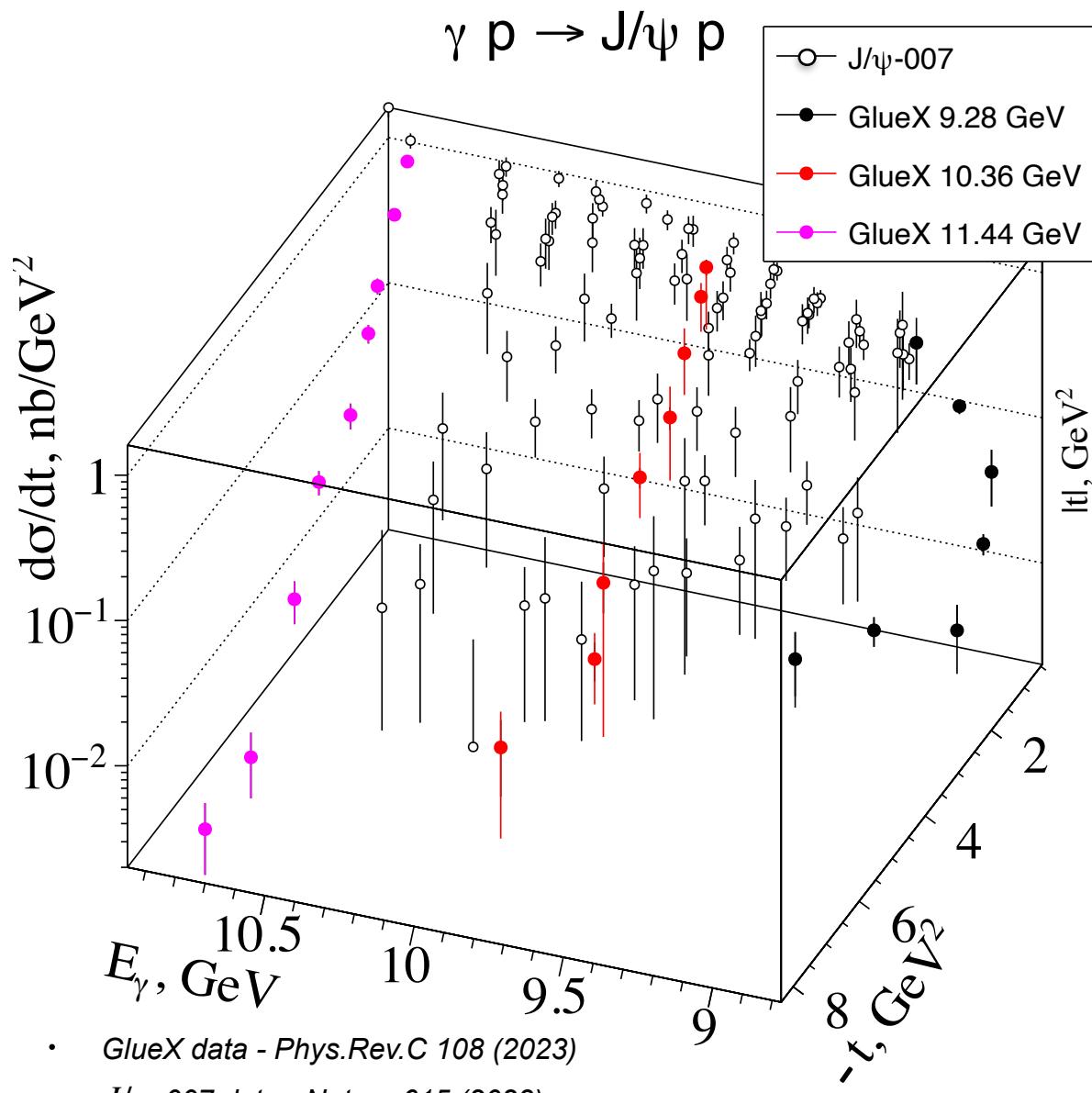
Data-driven gluon GFF extraction

What can we learn from threshold charmonium production?

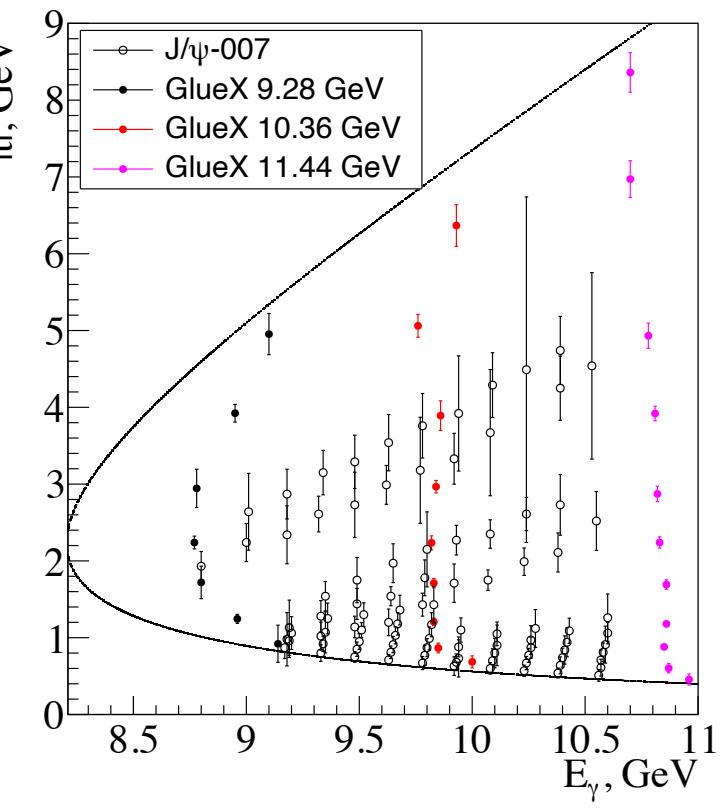
Lubomir Pentchev

- Test GFF models using data from threshold J/ψ photoproduction
 - find features in the data consistent with some general model predictions
- Then we will attempt to extract the gluon GFFs from the data only, without any additional theoretical/lattice constraints
- How the extracted gGFFs compare to lattice results
- Comparing J/ψ and χ_c photoproduction near threshold
 - Why they are different?
 - What can we learn about the production mechanism?
- Why do we need 22 GeV accelerator?

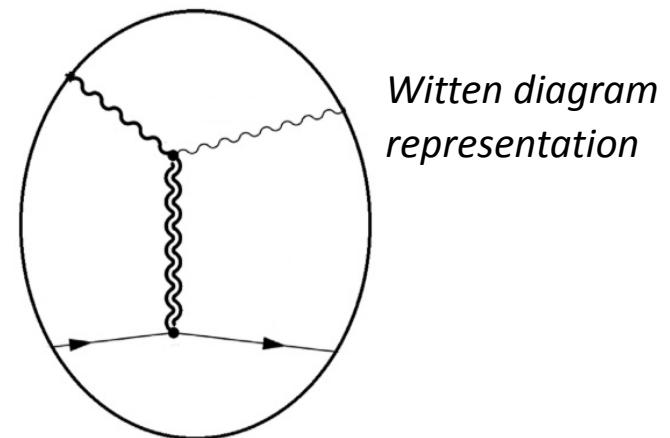
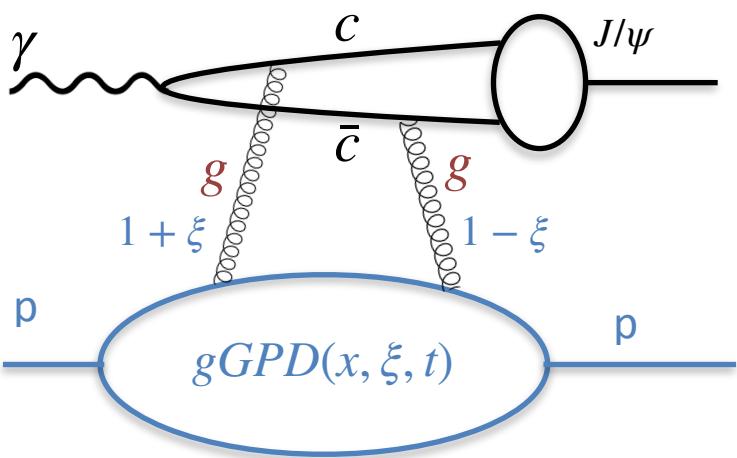
Threshold J/ψ photoproduction - the data



*find features in the data
consistent with some
general model predictions*



Threshold charmonium photoproduction - GPD and holographic approaches



$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = F(E_\gamma) \xi^{-4} [G_0(t) + \xi^2 G_2(t)] + \dots$$

$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t) \right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t) \right)^2$$

$$G_2(t) = 2\mathcal{A}_g^{(2)}(t)\mathcal{C}_g(t) + 2\frac{t}{4m^2}\mathcal{B}_g^{(2)}(t)\mathcal{C}_g(t) - \left(\mathcal{A}_g^{(2)}(t) + \mathcal{B}_g^{(2)}(t) \right)^2$$

$$\mathcal{A}_g^{(2)}(t) = A_g(t)$$

$$\mathcal{B}_g^{(2)}(t) = B_g(t)$$

$$\mathcal{C}_g(t) = 4C_g(t)$$

$$\xi = \frac{M_{J/\psi}^2 - t}{2(s - m^2) - M_{J/\psi}^2 + t}$$

leading moments
for high ξ values

$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = N(E_\gamma) [H_0(t) + \eta^2 H_2(t)] + \dots$$

$$H_0(t) = A_g^2(t)$$

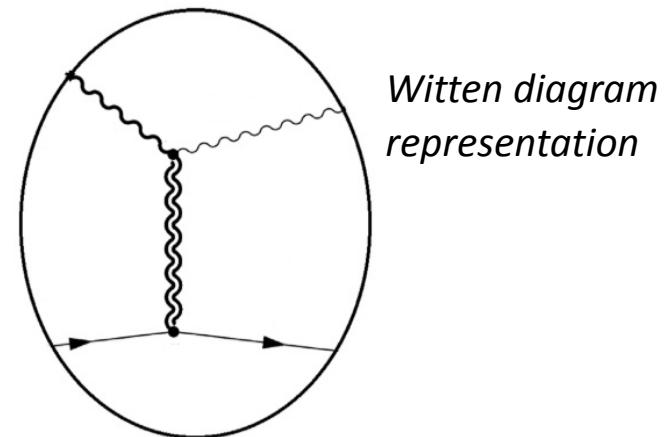
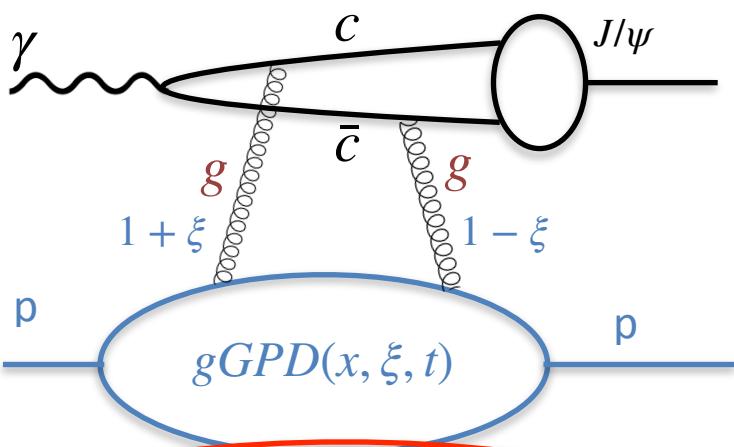
$$H_2(t) = 8A_g(t)C_g(t)$$

for large N_c and
strong α_s

$$\eta = \frac{M_{J/\psi}^2}{2(s - m^2) - M_{J/\psi}^2 + t}$$

Holographic analysis by Mamo and Zahed PRD 106 (2022), PRD, PRD 101 (2020), Hatta and Yang PRD 98 (2018)

Threshold charmonium photoproduction - GPD and holographic approaches



$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = F(E_\gamma) \xi^{-4} [G_0(t) + \xi^2 G_2(t)] + \dots$$

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for high ξ values
leading moments

$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = N(E_\gamma) [H_0(t) + \eta^2 H_2(t)] + \dots$$

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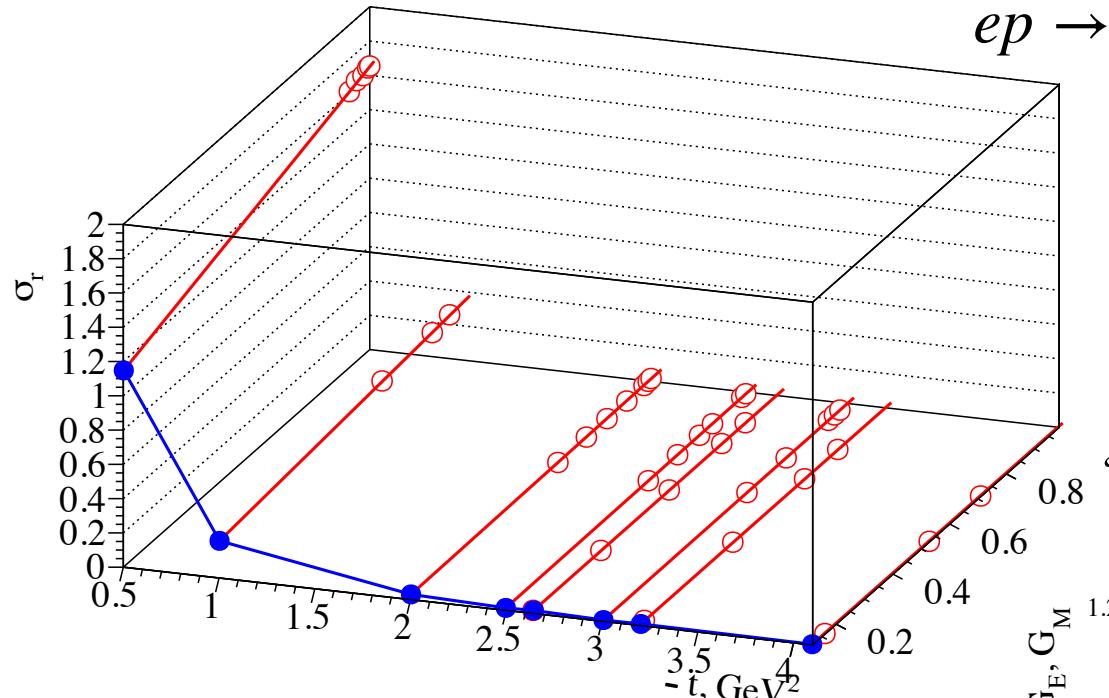
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Holographic analysis by Mamo and Zahed PRD 106 (2022),
PRD, PRD 101 (2020), Hatta and Yang PRD 98 (2018)

Rosenbluth separation

$G_E^2(t)$ and $\tau G_M^2(t)$ extraction

$ep \rightarrow ep$

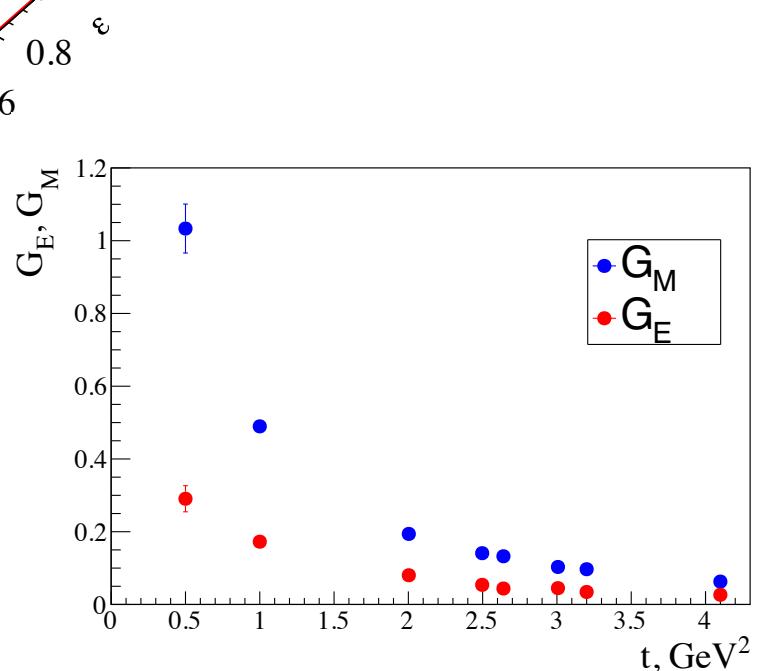


Data from:

R.C.Walker et al. Phys.Rev.D 49, 11 (1994) - SLAC

I.A.Qattan et al. arXiv 2411.05201 (2024) - JLab

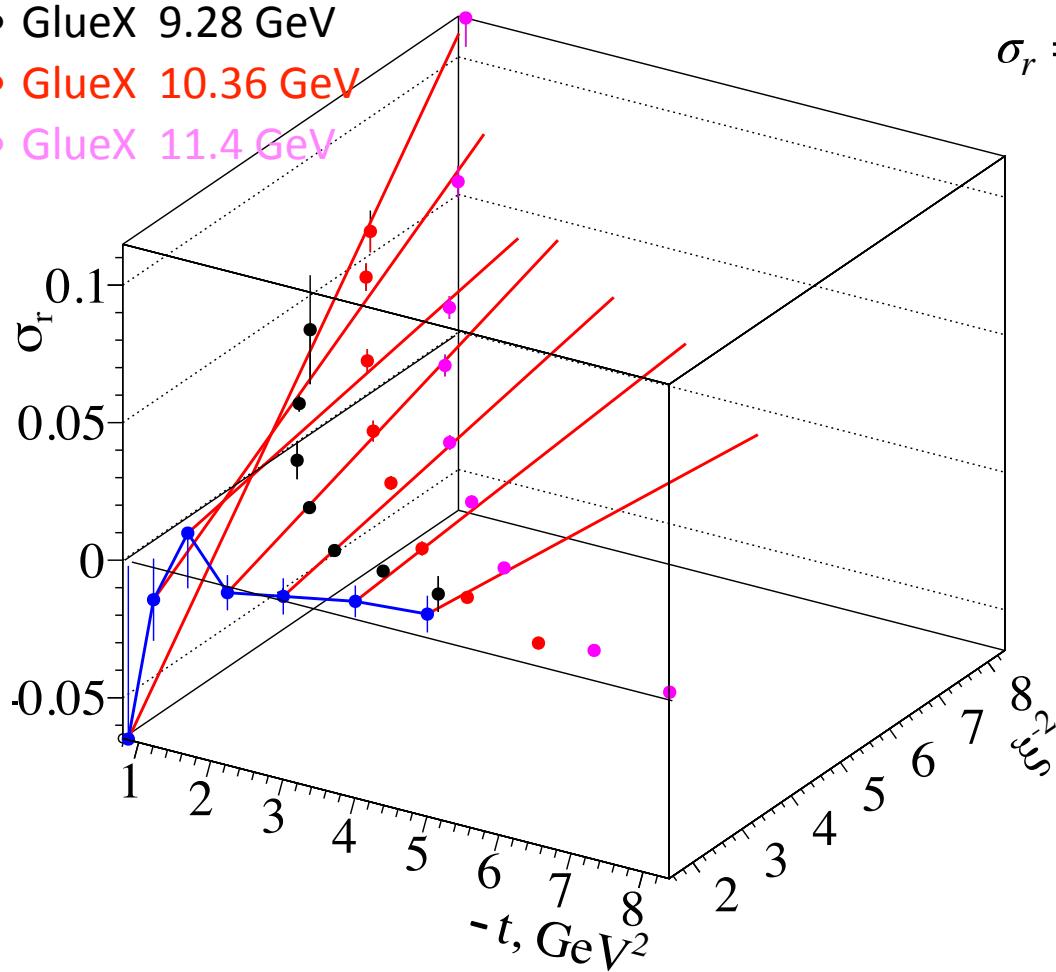
$$\sigma_R = \frac{d\sigma}{d\Omega} / \left(\frac{d\sigma}{d\Omega} \right)_M \quad \epsilon(1 + \tau) = \epsilon G_E^2(t) + \tau G_M^2(t)$$



Threshold J/ψ photoproduction - “Rosenbluth” separation

$G_0(t)$ and $G_2(t)$ extraction

- GlueX 9.28 GeV
- GlueX 10.36 GeV
- GlueX 11.4 GeV



$$\sigma_r = \frac{d\sigma}{dt} \frac{\xi^2}{F(E_\gamma)} = \xi^{-2} G_0(t) + G_2(t)$$

$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t) \right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t) \right)^2 > 0$$

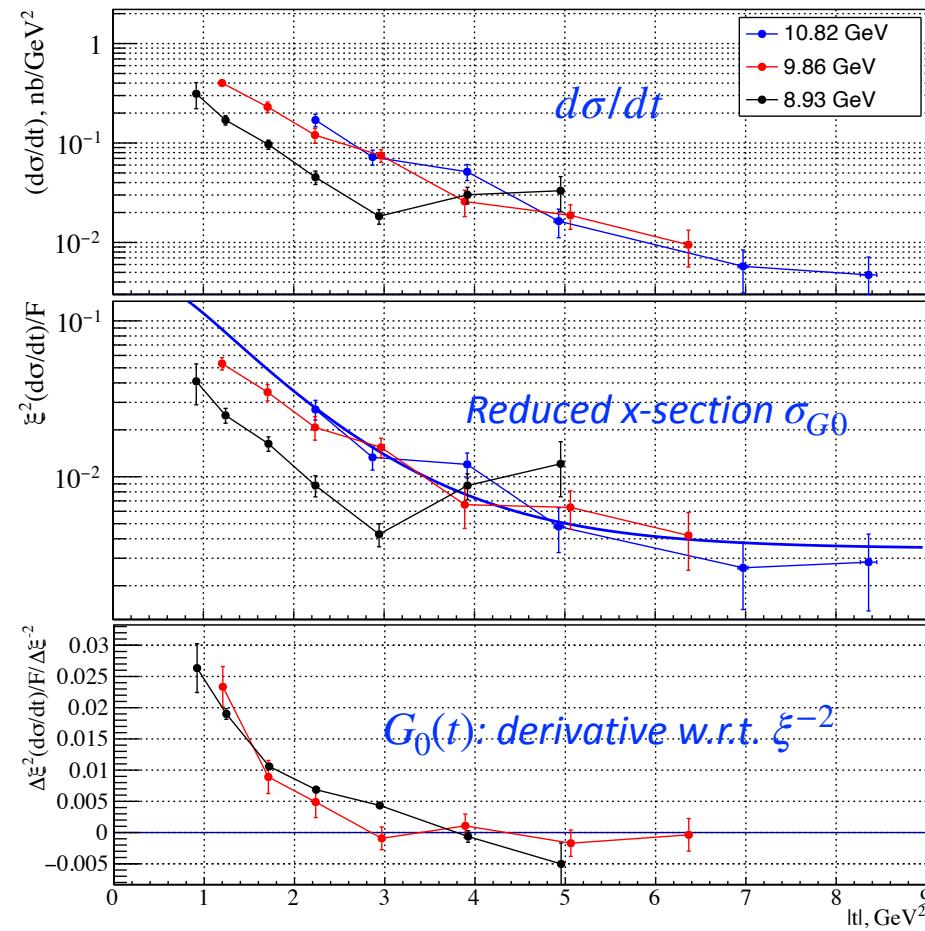
Therefore $d\sigma/dt(E, t)$ at fixed t , must increase with energy, consistent with the data!

Testing Energy Independence of Form Factor Functions

$G_0(t)$ - GPD

$$\sigma_{G0} = \frac{d\sigma}{dt} \frac{\xi^2}{F(E_\gamma)} = \xi^{-2} G_0(t) + G_2(t) \quad \xi > 0.4$$

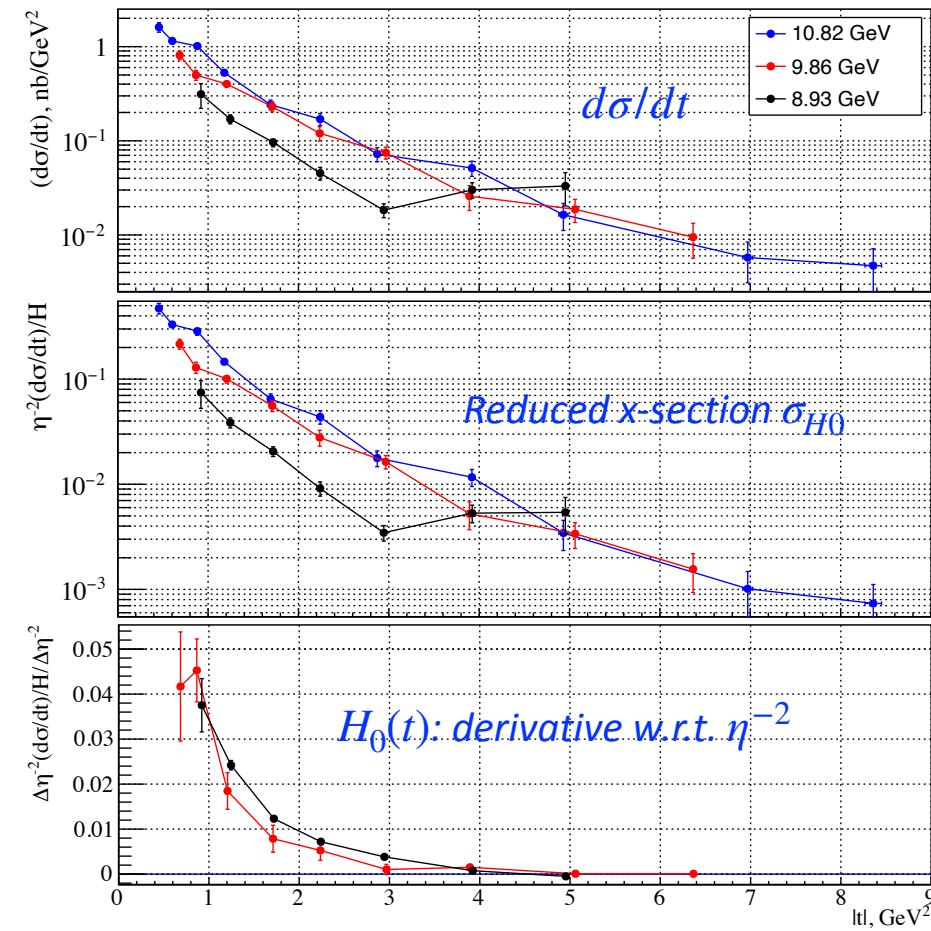
$$G_0(t) = \left[\sigma_{R0}(E_i, t) - \sigma_{R0}(E_j, t) \right] / \left[\xi^{-2}(E_i, t) - \xi^{-2}(E_j, t) \right]$$



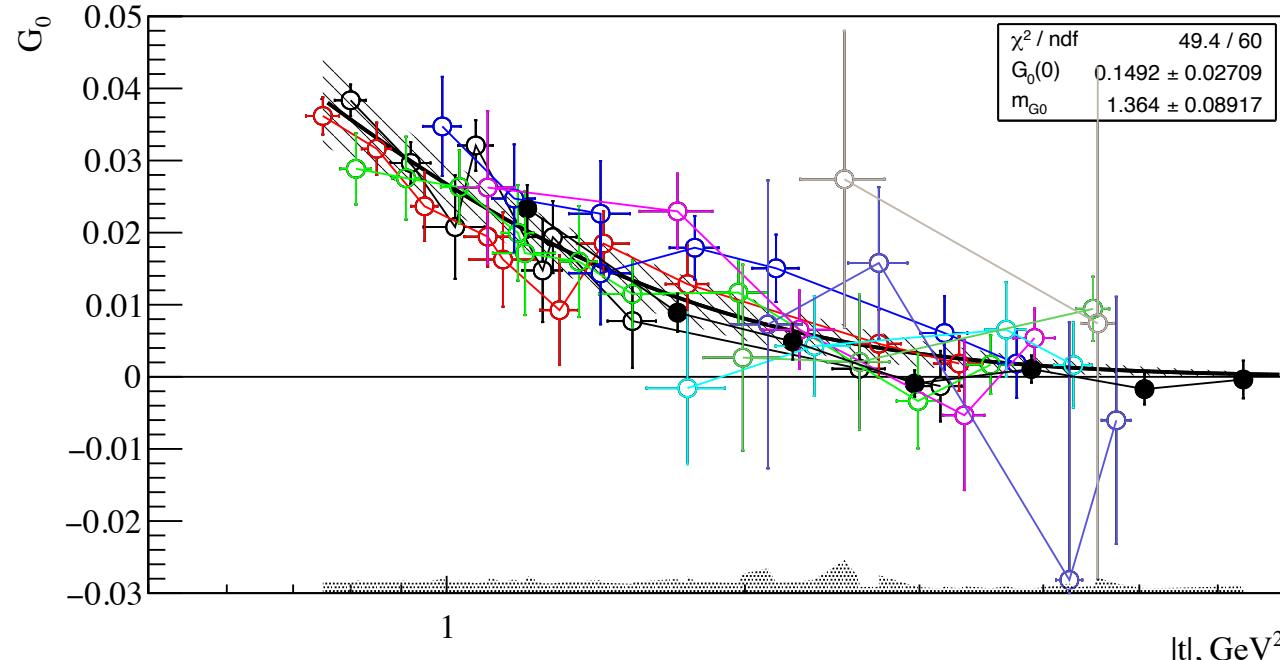
$H_0(t)$ - Holography

$$\sigma_{H0} = \frac{d\sigma}{dt} \frac{\eta^{-2}}{H(E_\gamma)} = \eta^{-2} H_0(t) + 4H_2(t)$$

$$H_0(t) = \left[\sigma_{R0}(E_i, t) - \sigma_{R0}(E_j, t) \right] / \left[\eta^{-2}(E_i, t) - \eta^{-2}(E_j, t) \right]$$



Gluon Form Factors - energy independence

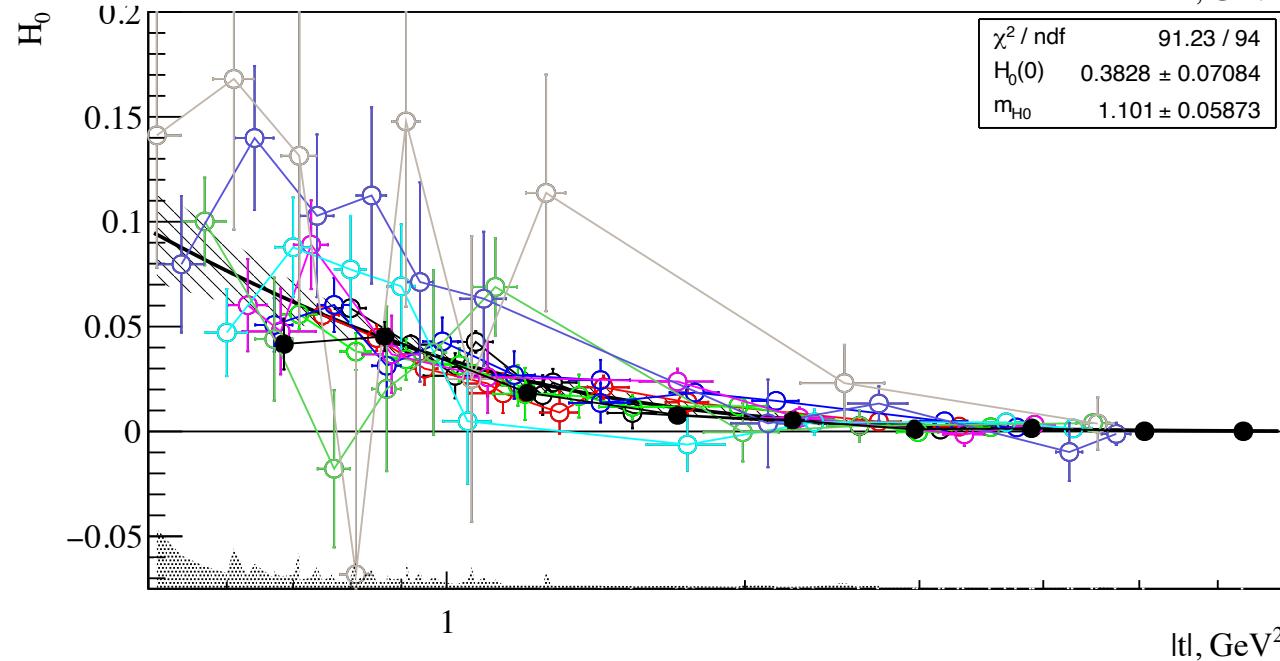


Energy independence of the $G(t)/H(t)$ functions (within errors)

Fits with:

$$\frac{G_0(0)}{(1 - t/m_{G_0}^2)^4}$$

$$\frac{H_0(0)}{(1 - t/m_{H_0}^2)^4}$$



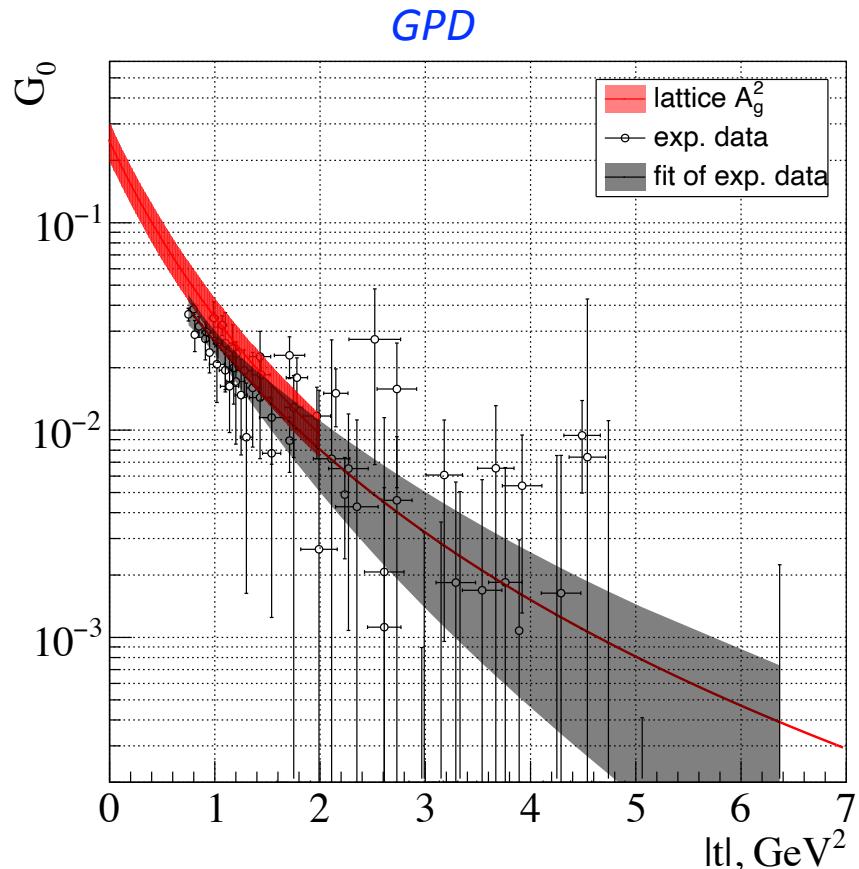
Using GlueX and J/ψ -007 data - different colors - different energies

No theory/lattice constraints used!

Phys.Rev.C 108 (2023)
Nature 615 (2023)

LP and E.Chudakov
arXiv:2404.18776

Assuming leading-term approximation - data vs lattice



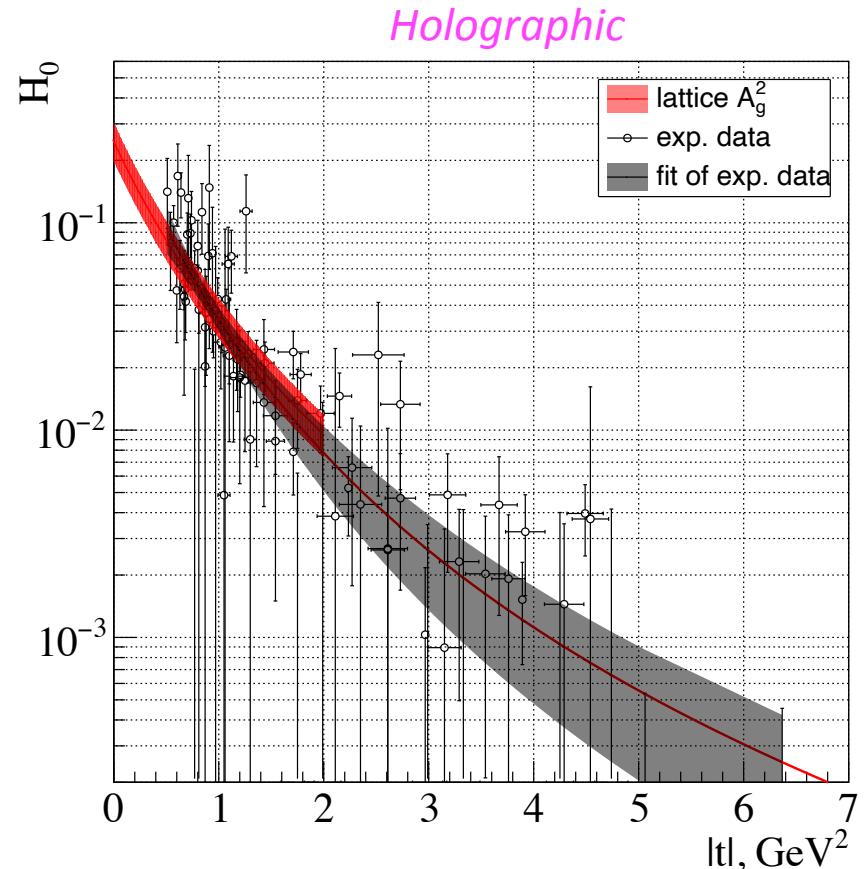
$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t) \right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t) \right)^2 \quad \text{for high } \xi \text{ values}$$

$\mathcal{A}_g^{(2)}(t) = A_g(t)$ *leading-moment*

$\mathcal{B}_g^{(2)}(t) = B_g(t)$ *approximation*

neglecting $B_g(t)$

$G_0(t) = A_g^2(t)$

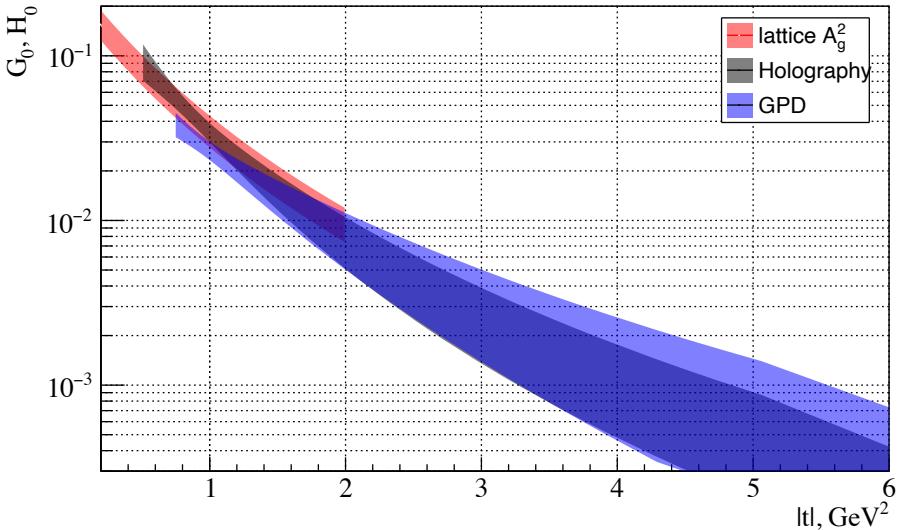


$B_g(t)=0$

$H_0(t) = A_g^2(t)$

for large N_c and
strong α_s

Gluon Gravitational Form Factors - summary



Features in data consistent with the GFF models:

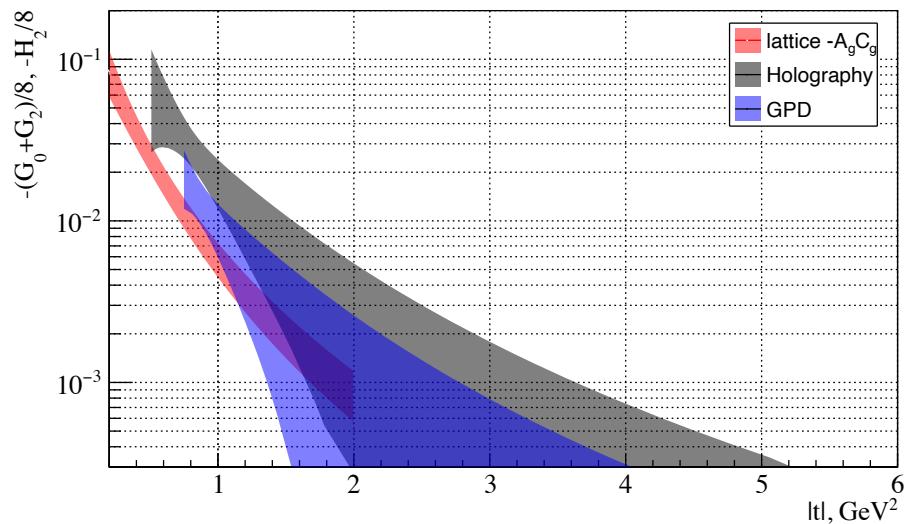
- $d\sigma/dt(E_\gamma, t \rightarrow \text{fixed})$ increases with energy
- $G(t)$ and $H(t)$ form factor functions are energy independent (within the experimental errors)
- In leading-term approximation (and neglecting B_g):

$$G_0(t) = H_0(t) = A_g^2(t) \text{ and}$$

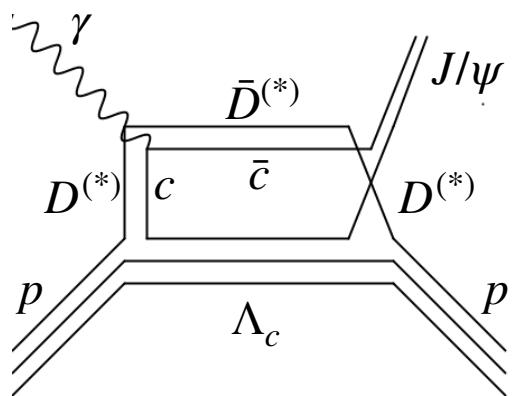
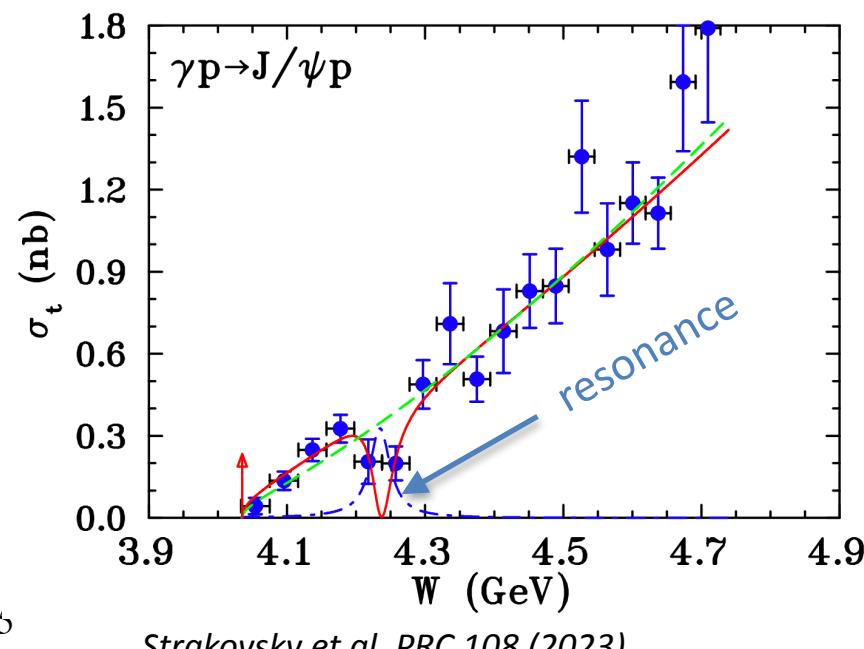
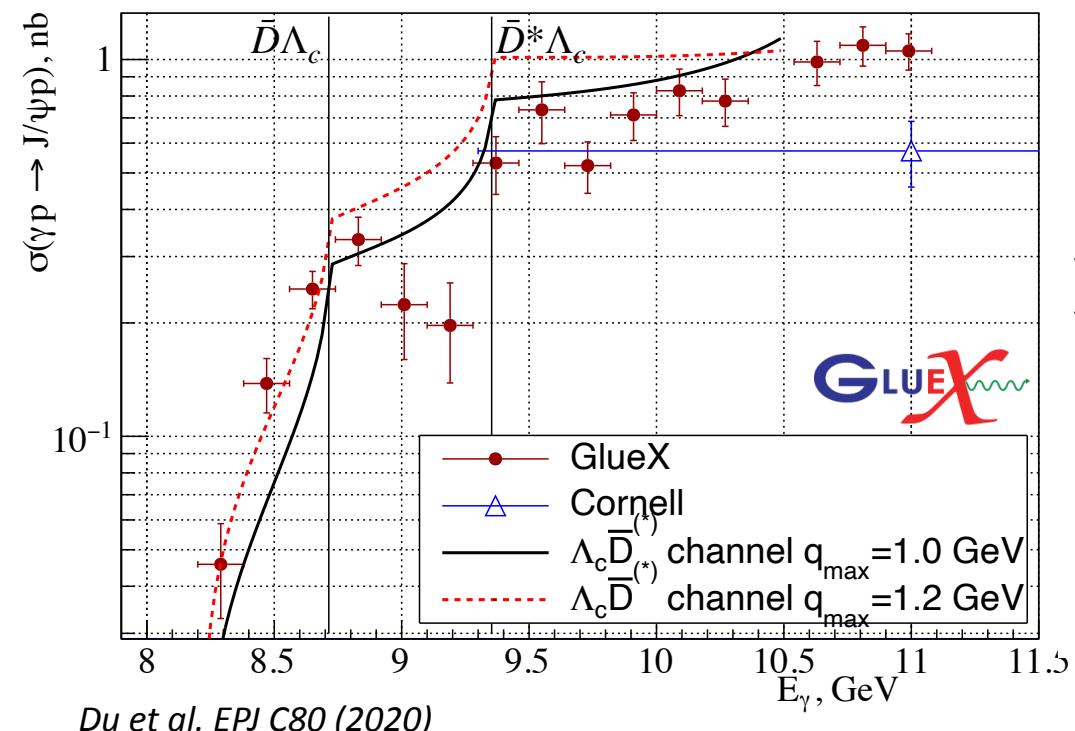
$$G_0(t) + G_2(t) = H_2(t) = 8C_g(t)A_g(t)$$

General agreement b/n extracted FFs using two diametric theories

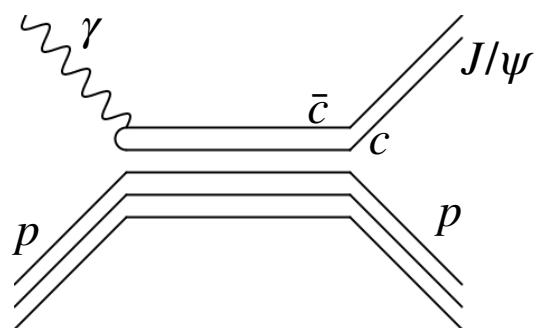
- agreement with lattice
- possible conclusion: the model corrections are not dominant



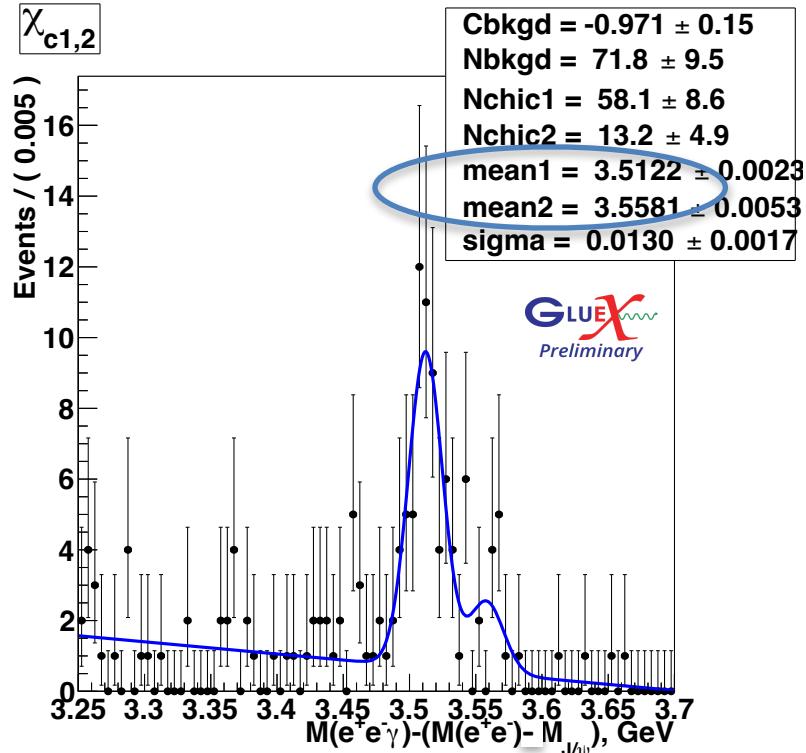
Other reaction mechanisms: open-charm, resonance



JPAC Phys.Rev.D 108 (2023)



Higher-mass charmonium states at threshold in GlueX

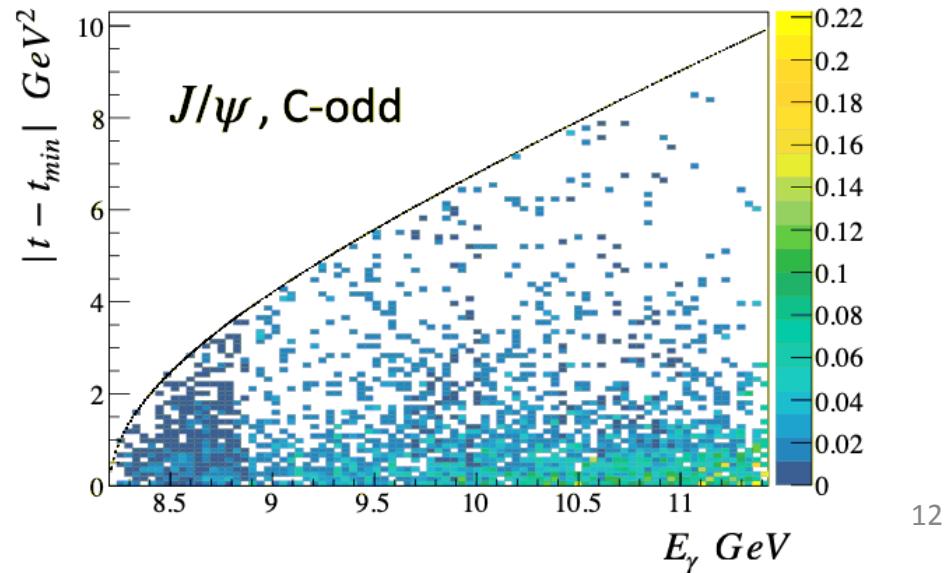
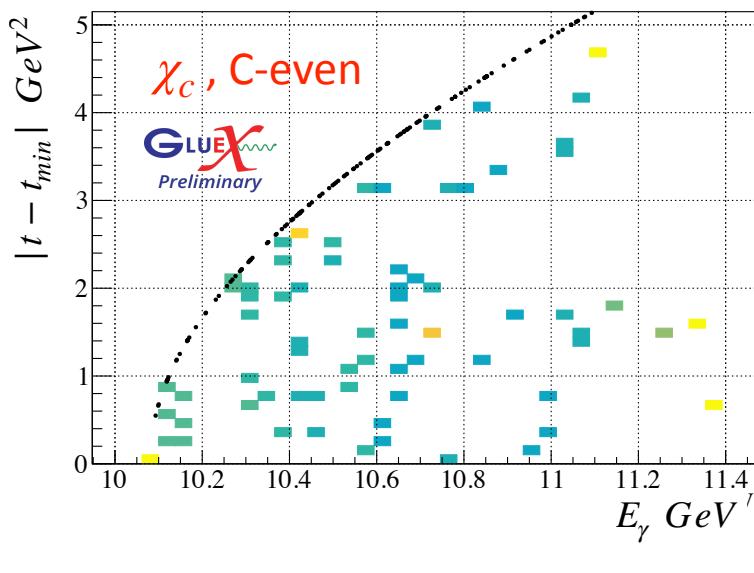


$\chi_c(3511)$ and $\chi_c(3556)$, 1^{++} and 2^{++} ,

$$E_\gamma^{thr} = 10.1 \text{ GeV}$$

- First ever evidence for photoproduction of C-even charmonium
- t -channel is suppressed (requires C-odd exchange), allows to study s -channel reaction mechanisms

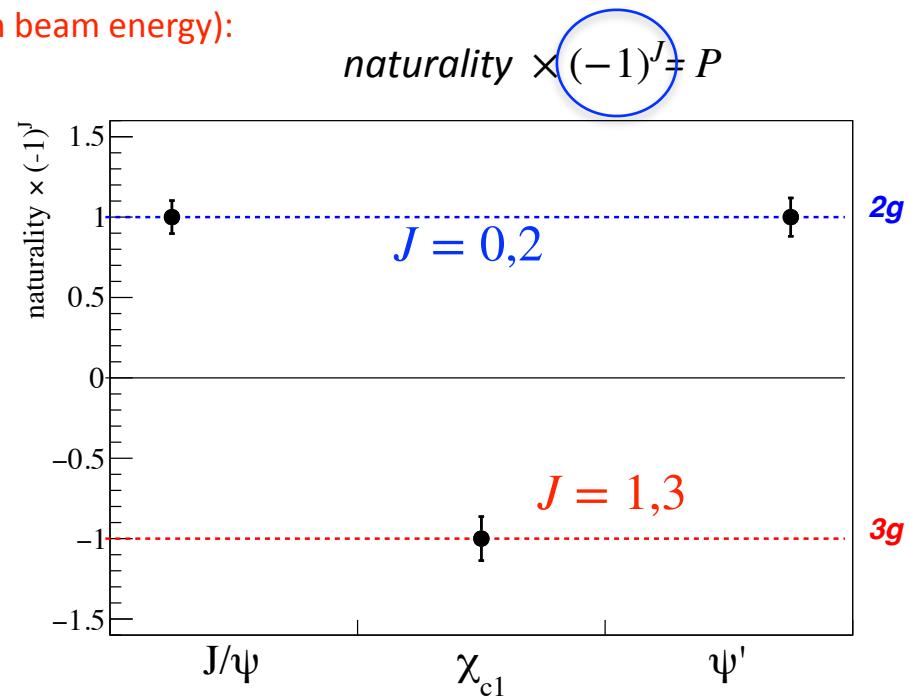
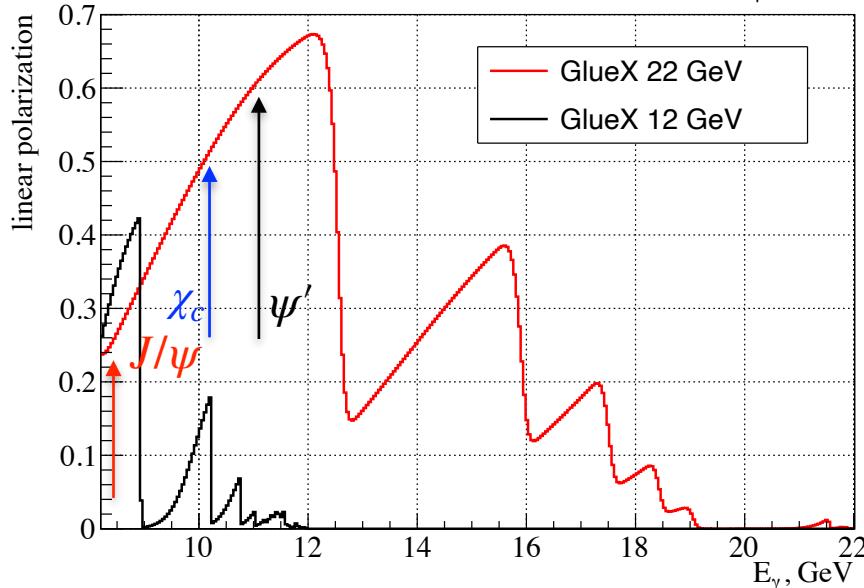
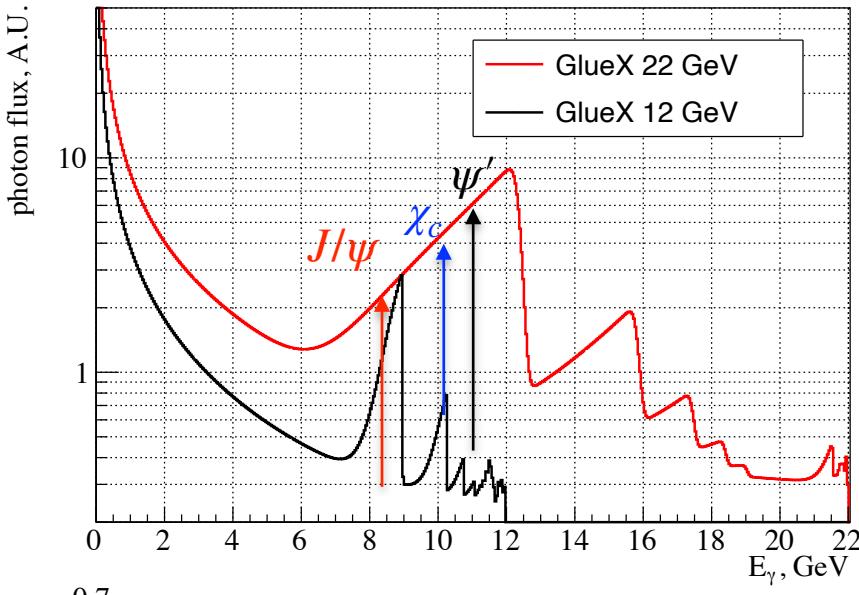
Dramatic difference: χ_c distribution in (E_γ, t) vs J/ψ



Threshold charmonium photoproduction at JLab22 with GlueX

GlueX uses polarized photon beam from coherent Bremsstrahlung

Taking advantage of increased end-point (electron beam energy):



$$asymmetry = \frac{2}{P_\gamma} \frac{Y_{J/\psi}(0) - Y_{J/\psi}(90)}{Y_{J/\psi}(0) + Y_{J/\psi}(90)} = \\ = - (\rho_{1-1}^1 - Im \rho_{1-1}^2) \cos[2(\phi_{hel} - \Phi)]$$

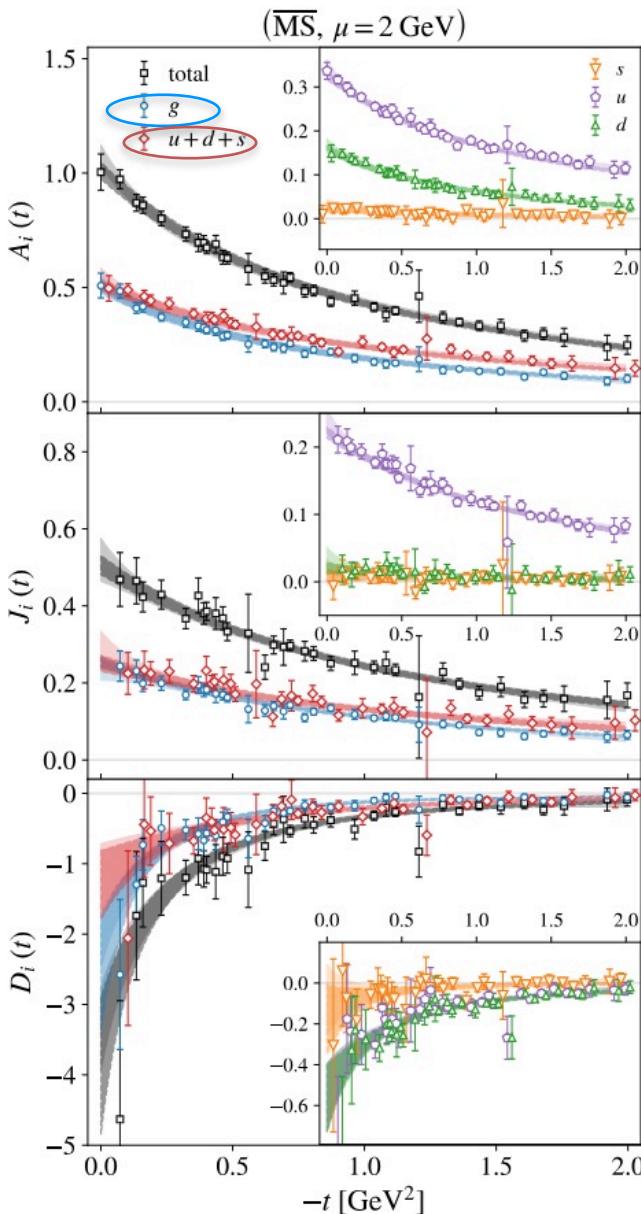
$$naturality = (-1)^J P$$

Conclusions

- Features in the J/ψ data consistent with some general predictions of the GFF theoretical models - **skewness scaling, $d\sigma/dt(E_\gamma, t \rightarrow \text{fix})$ increases with energy**
- Gluon GFFs extracted from J/ψ data **using two diametric approaches (without any external constraints)** on the same scale with the lattice calculations
- Higher mass charmonium states - complementary in understanding the reaction mechanism
- CEBAF energy upgrade adds new dimensions to these studies:
 - Threshold production of higher-mass charmonium states (with different quantum numbers) - $\psi(2S)$, χ_c
 - Polarization measurements with high FOM

Back up slides

Gluonic structure of the proton

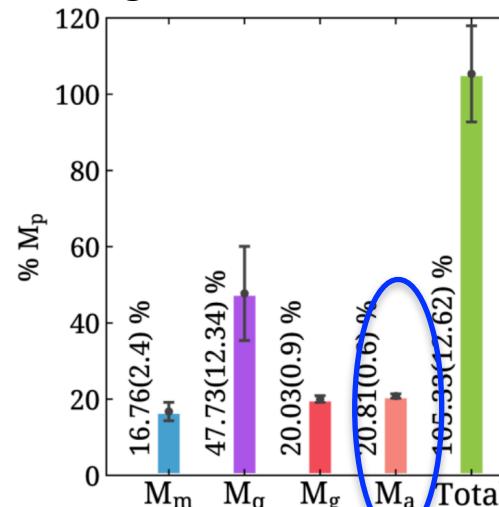


- Gluon contribution to the mechanical properties of the proton equally important as the quark one:

Lattice calculations of Gravitational Form Factors (GFFs) show similar contributions from **gluons (g)** and **quarks (u+d+s)**.

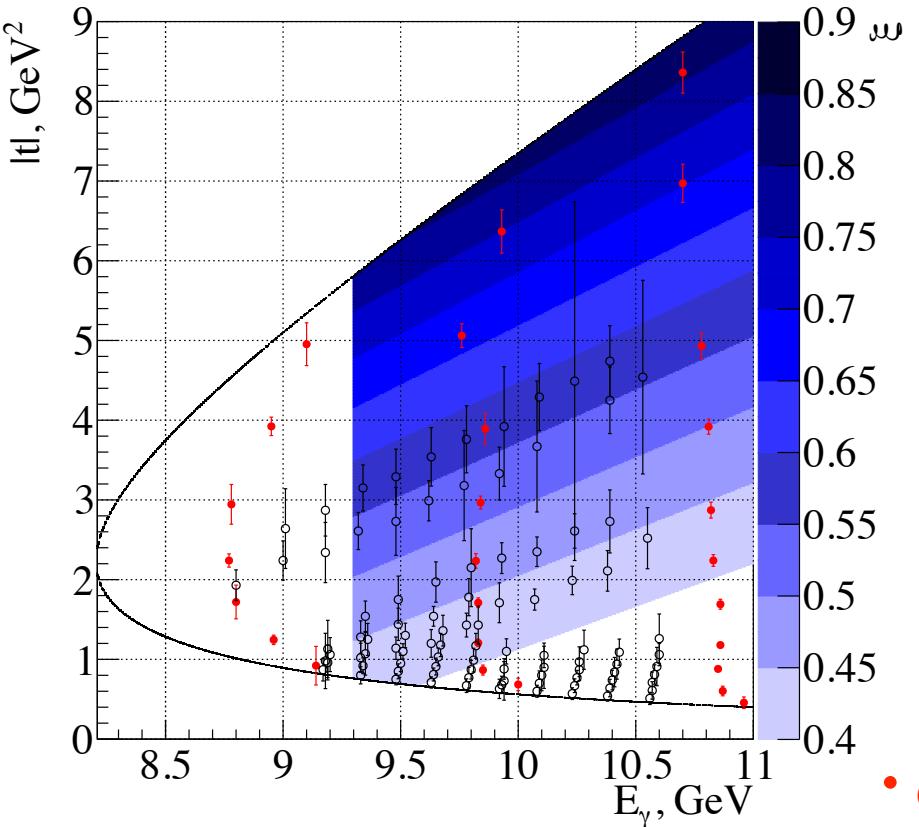
Hackett, Pefkou, Shanahan arxiv:2310.08484 (2023)

- Quark masses and kinetic energies of quarks and gluons are not enough to explain the mass of the proton: gluon condensate, or anomalous contribution to the mass of the proton is significant:



Data used in gluon Form Factor extraction

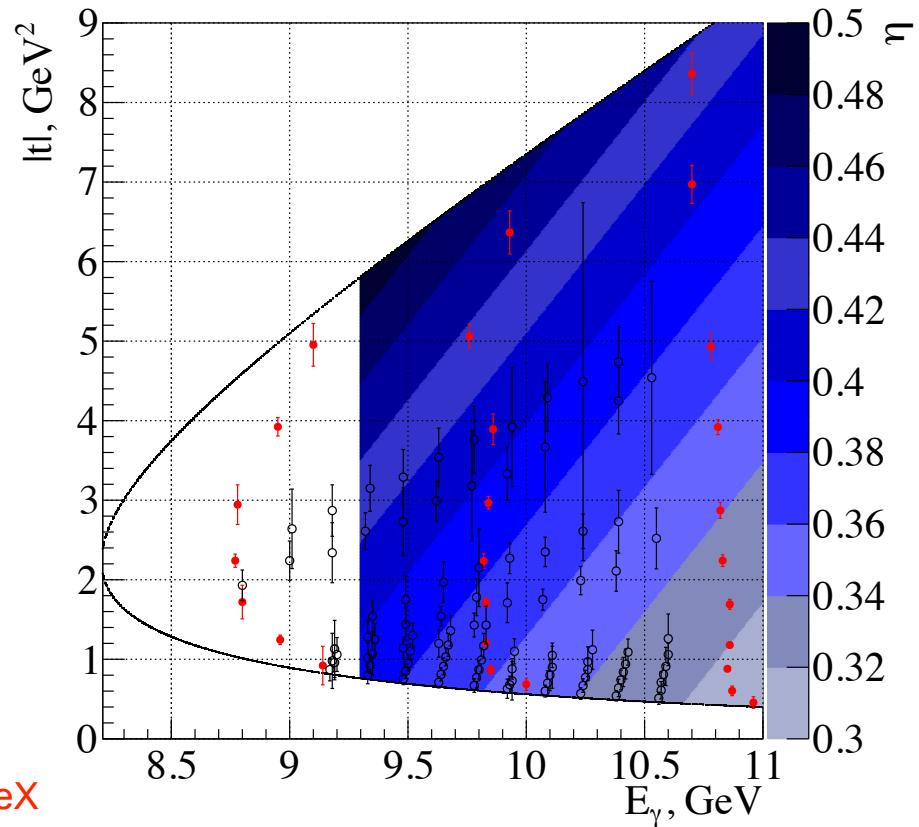
GPD



$\xi > 0.4$

$E_\gamma > 9.3 \text{ GeV}$ (away from
 $\bar{D}\Lambda_C$ and $\bar{D}^*\Lambda_C$ thresholds)

Holography



no constraints on η

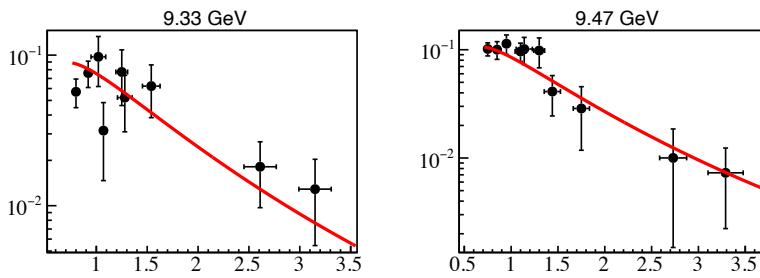
$E_\gamma > 9.3 \text{ GeV}$ (away from
 $\bar{D}\Lambda_C$ and $\bar{D}^*\Lambda_C$ thresholds)

Error bars: relative errors in
A.U. (not related to y-axis)

Using η -scaling to describe data (Holographic approach)

$d\sigma/dt \cdot \eta^{-2}/H$ vs $|t|$ (GeV^2)

χ^2/ndf	108.7 / 102
$A_g(0)$	0.3174 ± 0.06962
m_A	1.164 ± 0.0623
$C_g(0)$	-3.64 ± 1.613
m_C	0.9066 ± 0.08587
const _{G2}	0.0003233 ± 0.000288



Fit to all data ($E_\gamma > 9.3$ GeV, **all t data points included**)

$$\frac{A_g^2(0)}{(1-t/m_A^2)^4}$$

$$\frac{A_g(0)C_g(0)}{(1-t/m_{AC}^2)^4} + \text{const.}$$

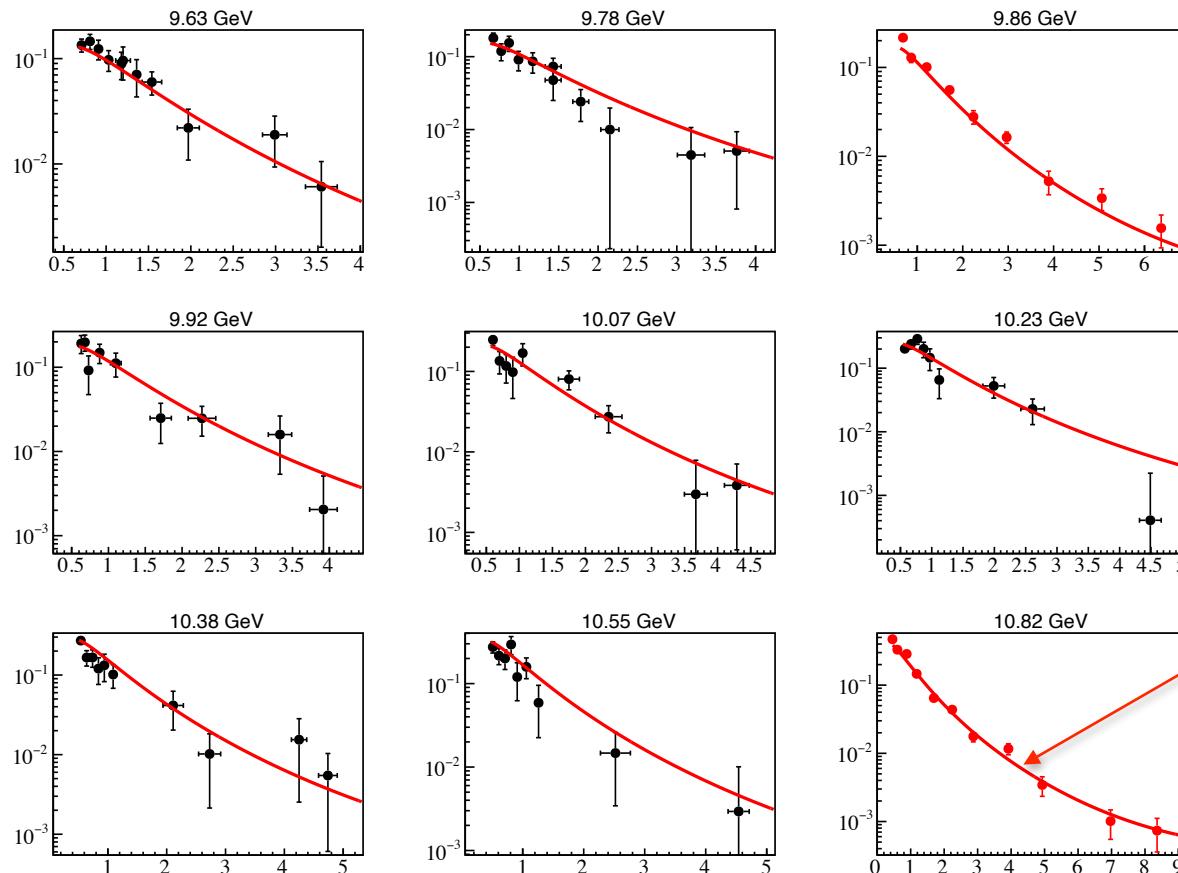
$$\left(\frac{d\sigma}{dt}\right) \frac{\eta^{-2}}{H(E_\gamma)} = A_g^2(t)\eta^{-2} + 8A_g(t)C_g(t) + \dots$$

$$\eta = \frac{M_{J/\psi}^2}{2(s - m_p^2) - M_{J/\psi}^2 + t}$$

- $\chi^2/ndf \sim 1$ in the chosen kinematic region

Only this fitted function used in the analysis

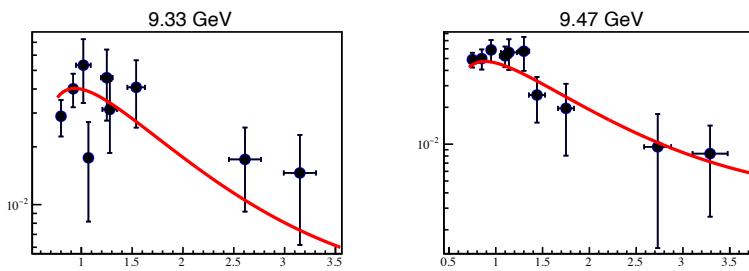
- GlueX
- J/ψ -007



Using ξ -scaling to describe data (GPD approach)

$d\sigma/dt \cdot \xi^2/F$ vs $|t|$ (GeV 2)

χ^2/ndf	53.96 / 63
$G_0(0)$	0.1417 ± 0.07354
m_{G0}	1.386 ± 0.1518
$G_2(0)$	-2.102 ± 1.37
m_{G2}	0.9458 ± 0.1331
const_{G2}	0.003472 ± 0.0007662

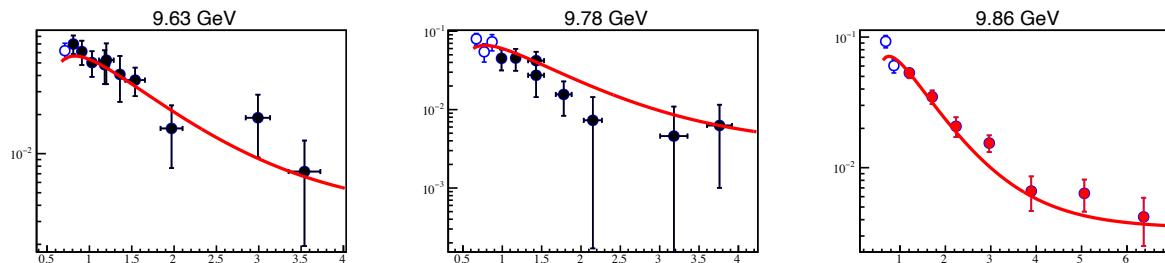


Fit to all data ($E_\gamma > 9.3$ GeV, $\xi > 0.4$)

$$\frac{G_0(0)}{(1-t/m_{G0}^2)^4} + \frac{G_2(0)}{(1-t/m_{G2}^2)^4} + \text{const.}$$

$$\left(\frac{d\sigma}{dt} \right) \frac{\xi^2}{F(E_\gamma)} = G_0(t)\xi^{-2} + G_2(t) + \dots$$

$$\xi = \frac{M_{J/\psi}^2 - t}{2(s - m_p^2) - M_{J/\psi}^2 + t}$$

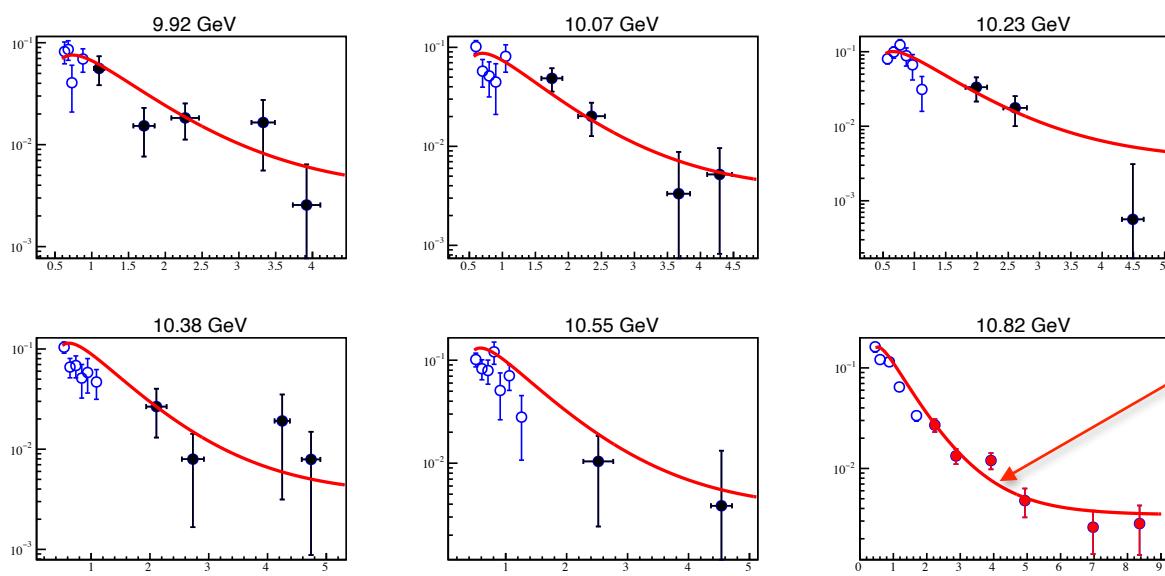


$\chi^2/\text{ndf} \sim 1$ in the chosen kinematic region

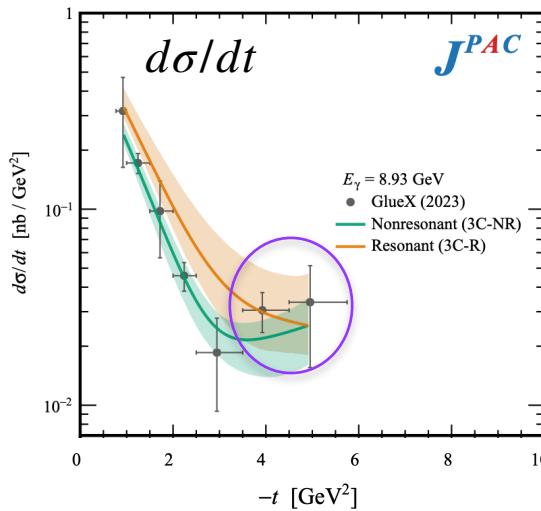
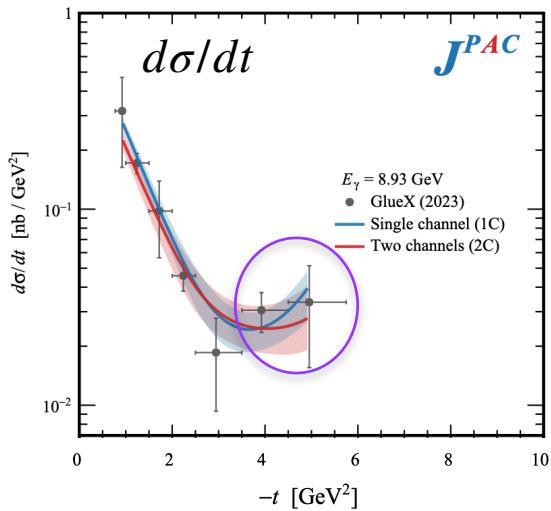
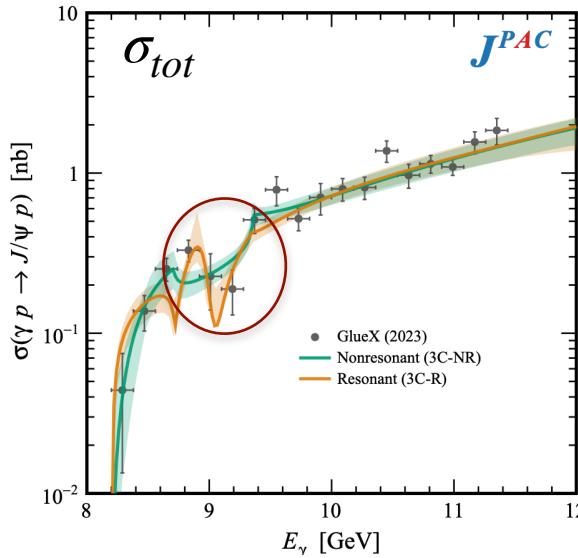
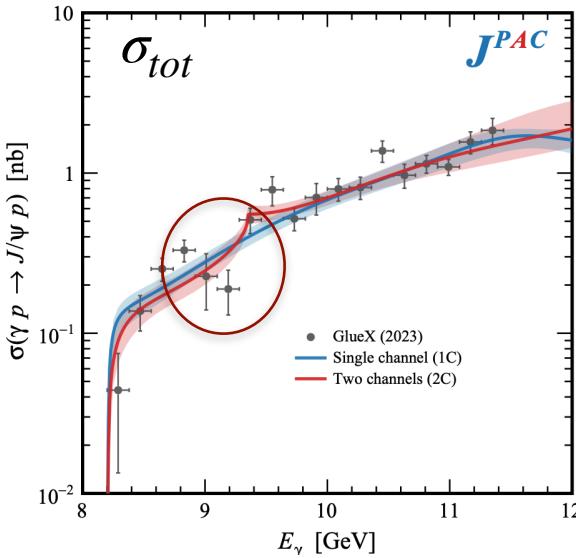
- $\xi < 0.4$ data points (blue) deviate from ξ -scaling

Only this fitted function used in the analysis

- GlueX
- J/ ψ -007



Phenomenological approach: JPAC results



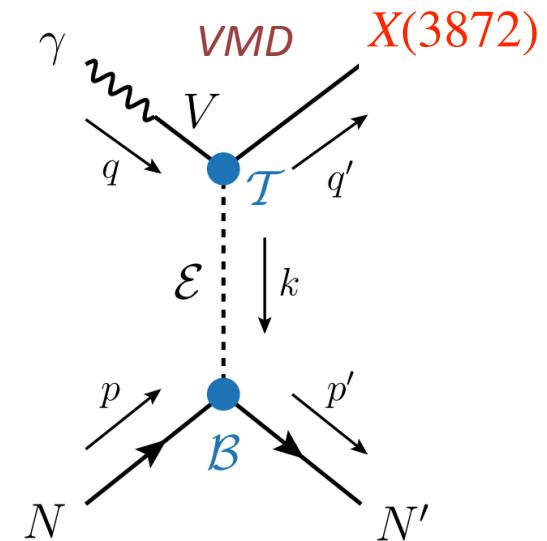
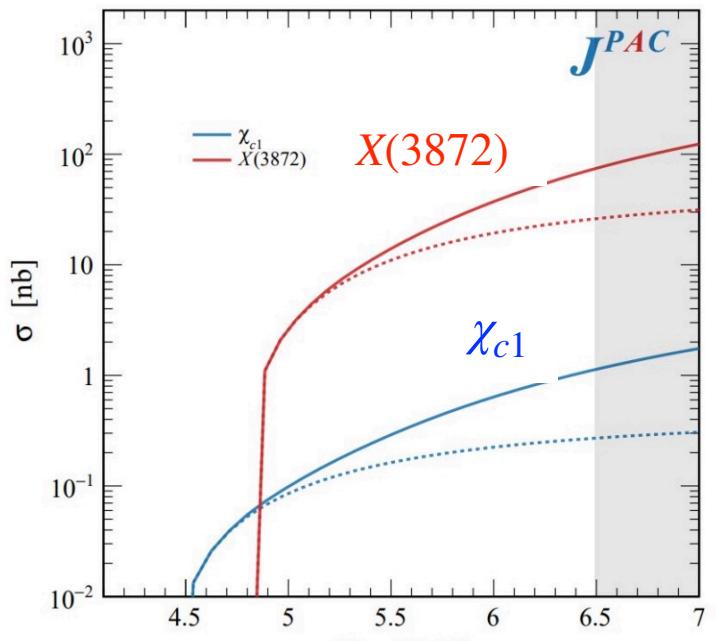
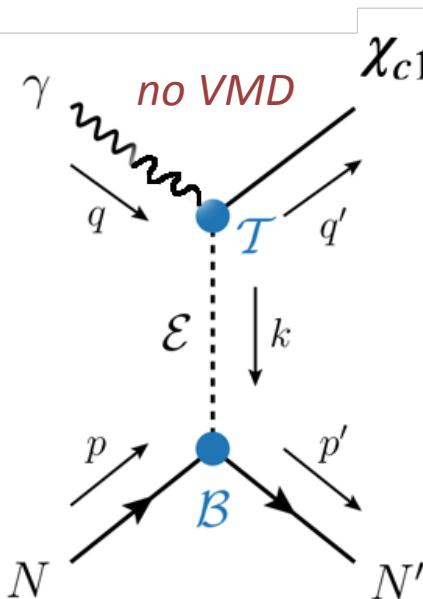
Phenomenological model based on s-channel PW expansion ($l \leq 3$):

- (1C) $J/\psi p$ interaction
- (2C) $J/\psi p$ and $\bar{D}^* \Lambda_C$
- (3C-NR) $J/\psi p$, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (non-resonant solution)
- (3C-NR) $J/\psi p$, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (resonant solution)

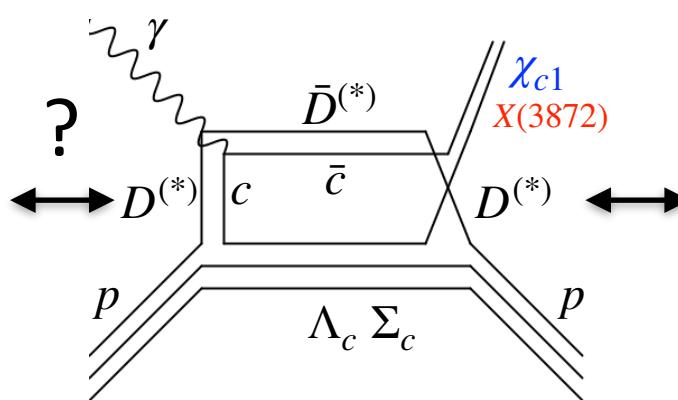
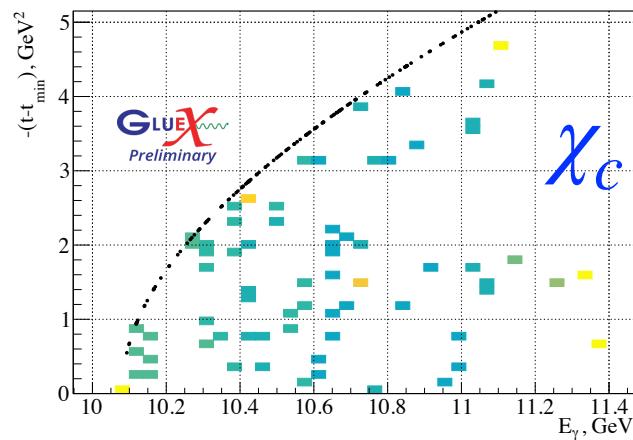
No stat. significant preference:

- 9 GeV structure requires sizable contribution from open charm
- Severe violation of VMD and factorization not excluded
- s-channel resonance not excluded
- t-enhancement indicates s-channel contribution: due to proximity to threshold or open-charm exchange

χ_c vs $X(3872)$ production



JPAC, PRD 102 (2020). t -channel production

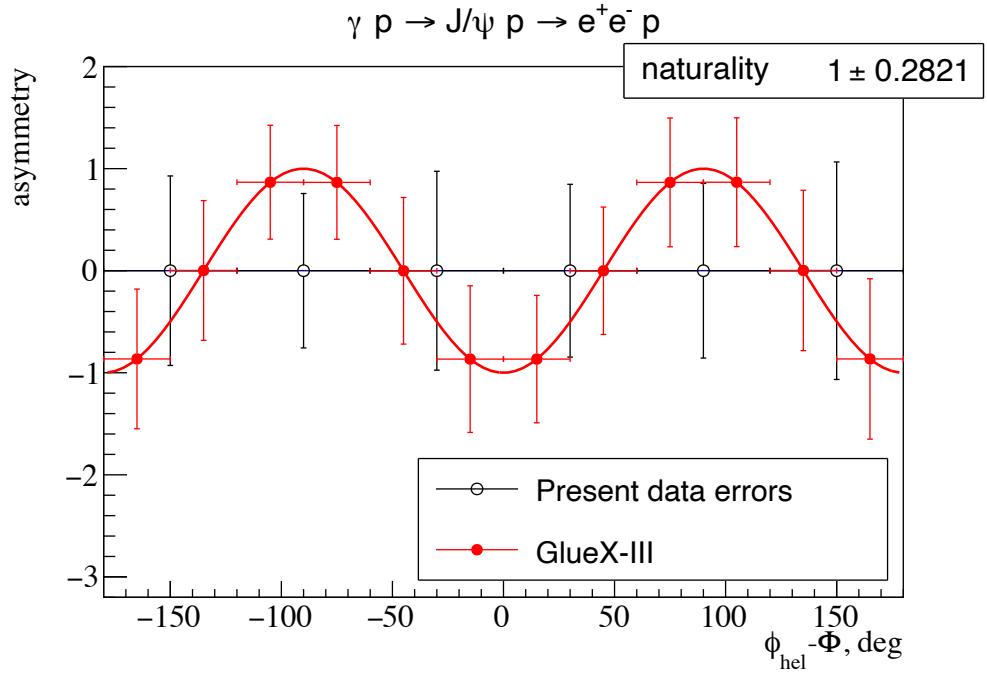
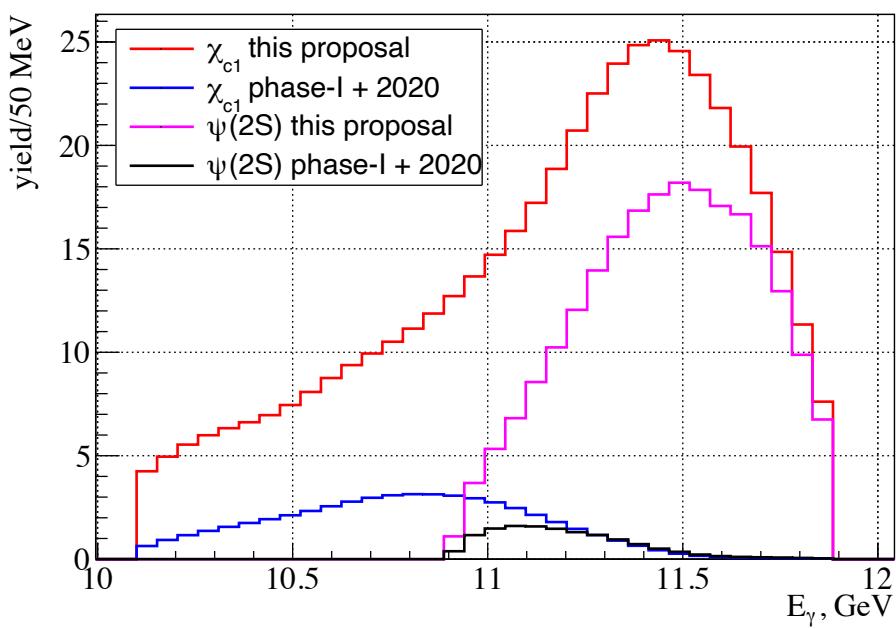


$X(3872)$
?

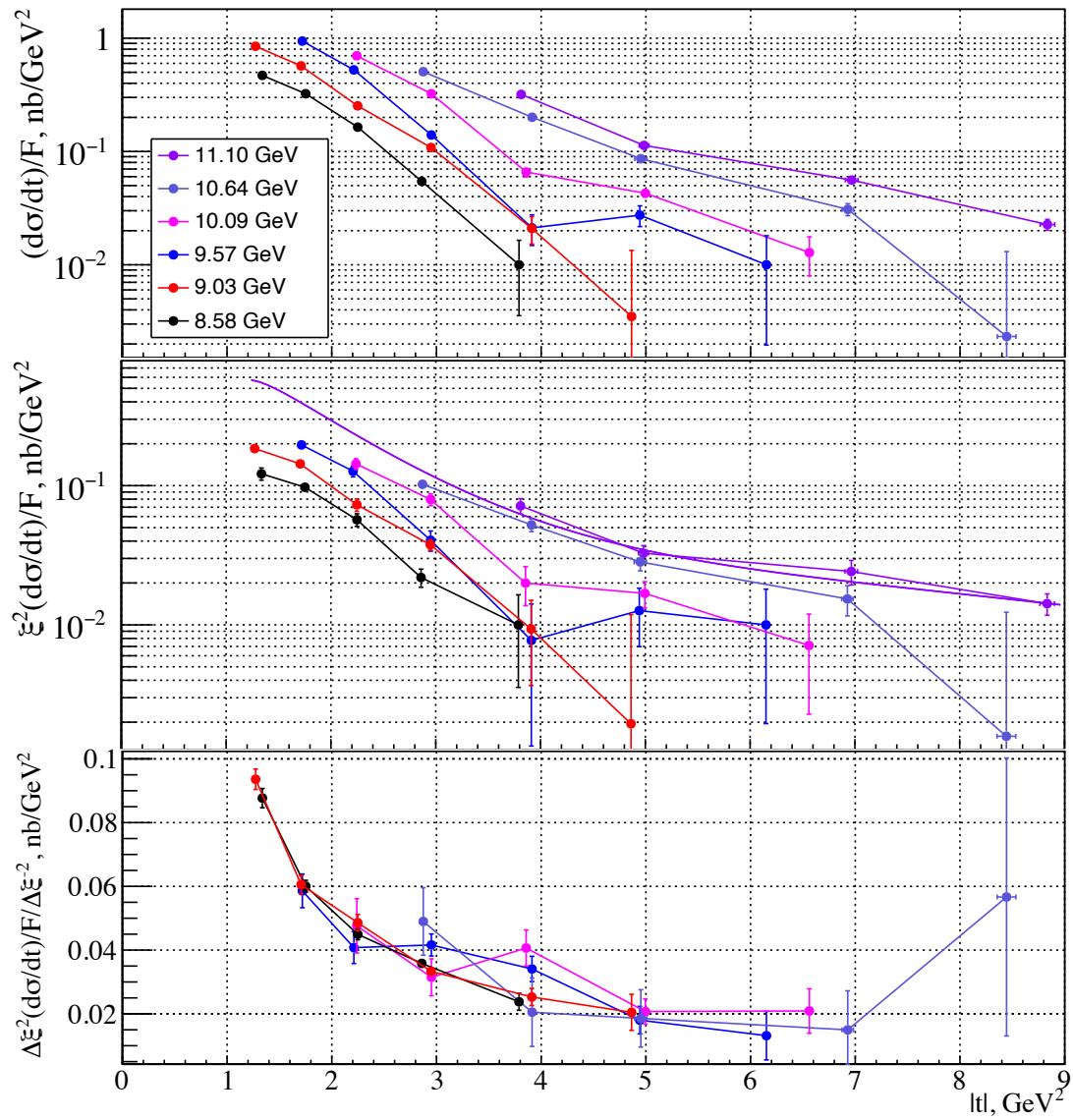
Studying χ_c can help to understand $X(3872)$ production mechanism

Prospect for charmonium threshold production with GlueX-III

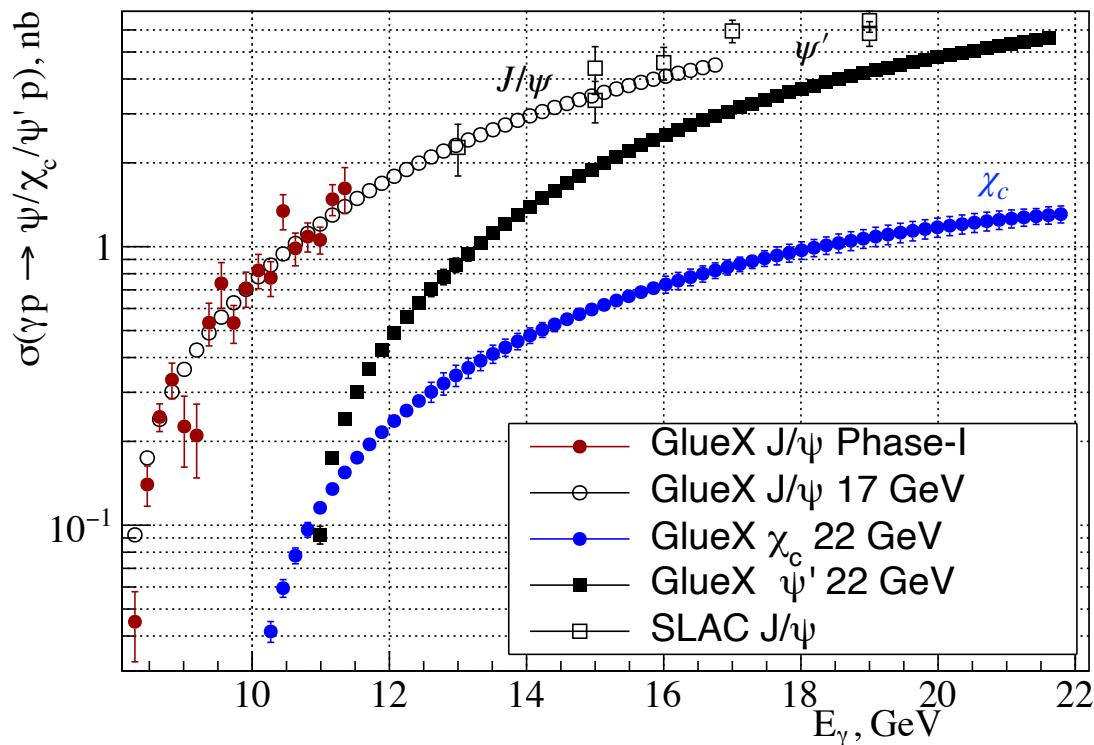
Run Period	J/ψ	χ_{c1}	$\psi(2S)$
2016-2020 Phase I-II	3,960	55	12
2023-2025 Phase II (planned)	3,615	48	11
Phase III (proposal)	11,271	364	178
Projected Total	18,846	467	201



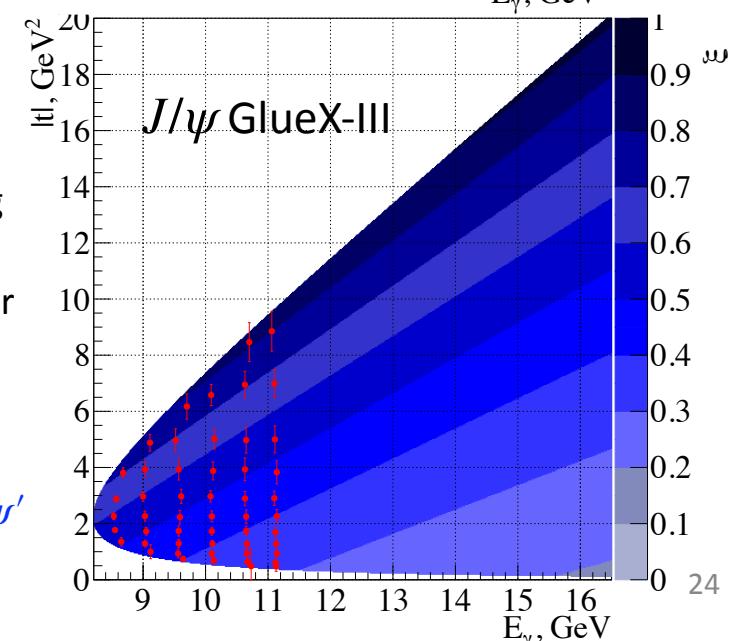
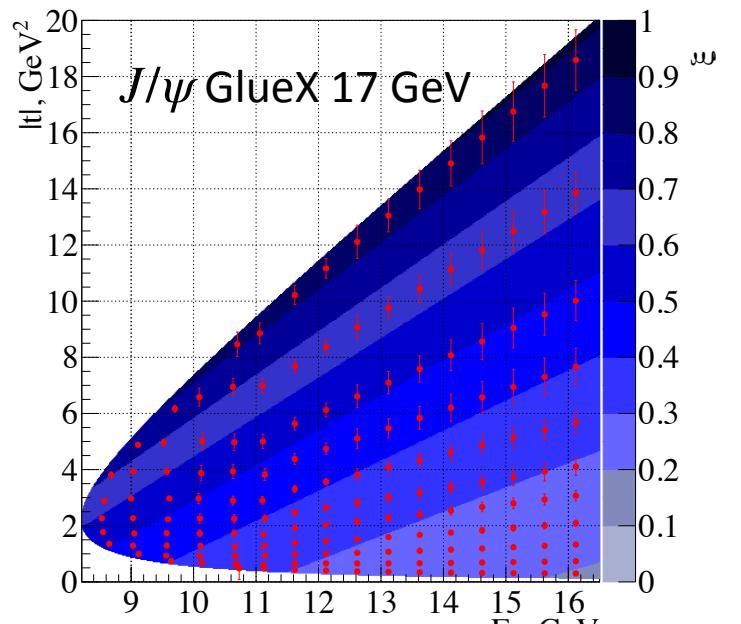
Prospect for charmonium threshold production with GlueX-III



Threshold charmonium photoproduction at JLab22 with GlueX

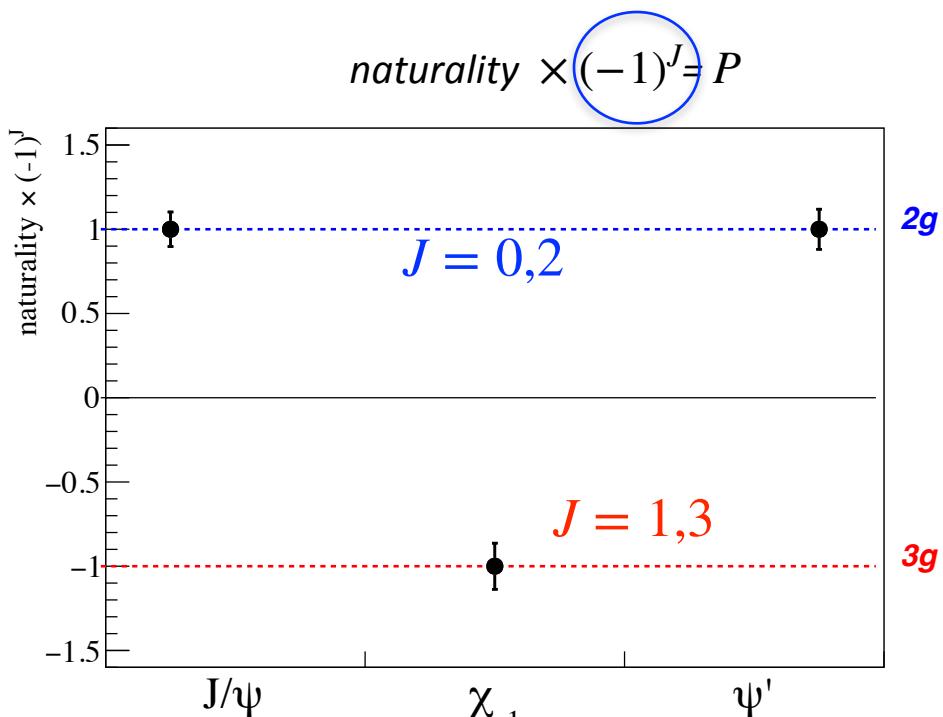
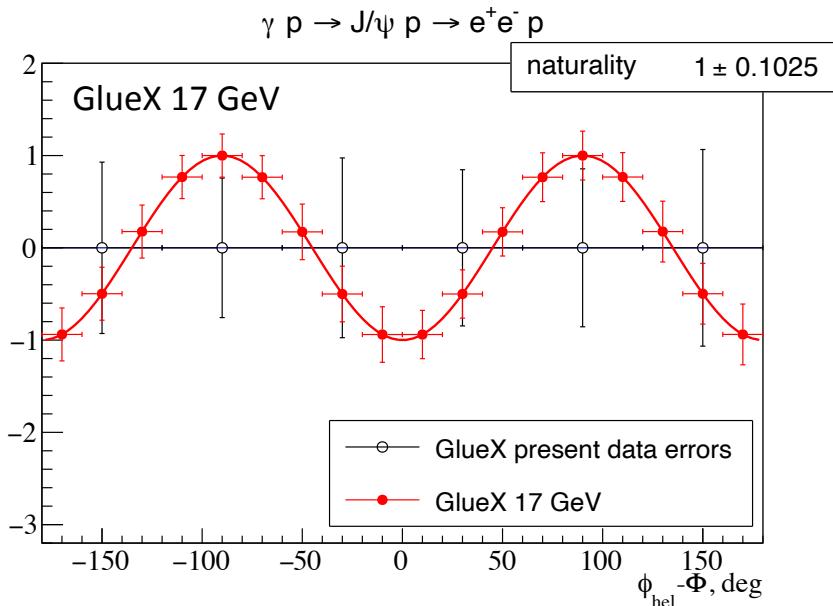


- Absolute normalization and error estimation: scaling existing measurements (<12 GeV) by anticipated flux increase
- Energy extrapolation: using fits of J/ψ data for J/ψ and ψ' , for χ_c based on VM exchange model (*JPAC, Phys.Rev.D 102 (2020)*)
- Efficiencies estimated with MC
- All JLab22 projections for 100 PAC days: 80k J/ψ , 8k χ_c , 18k ψ'

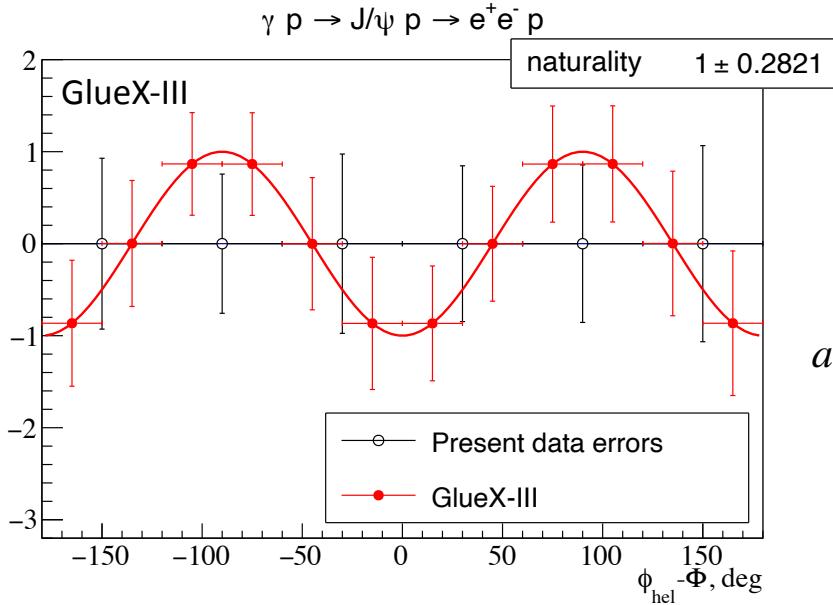


Threshold charmonium photoproduction at JLab22 with GlueX

asymmetry



asymmetry

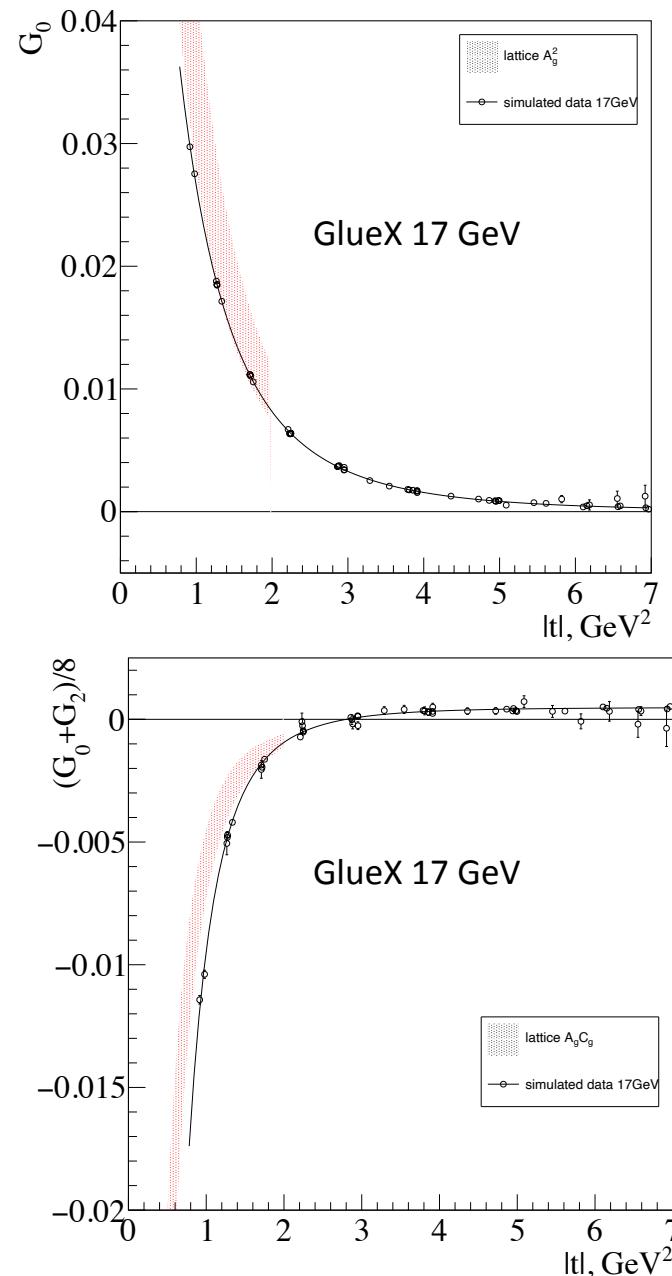
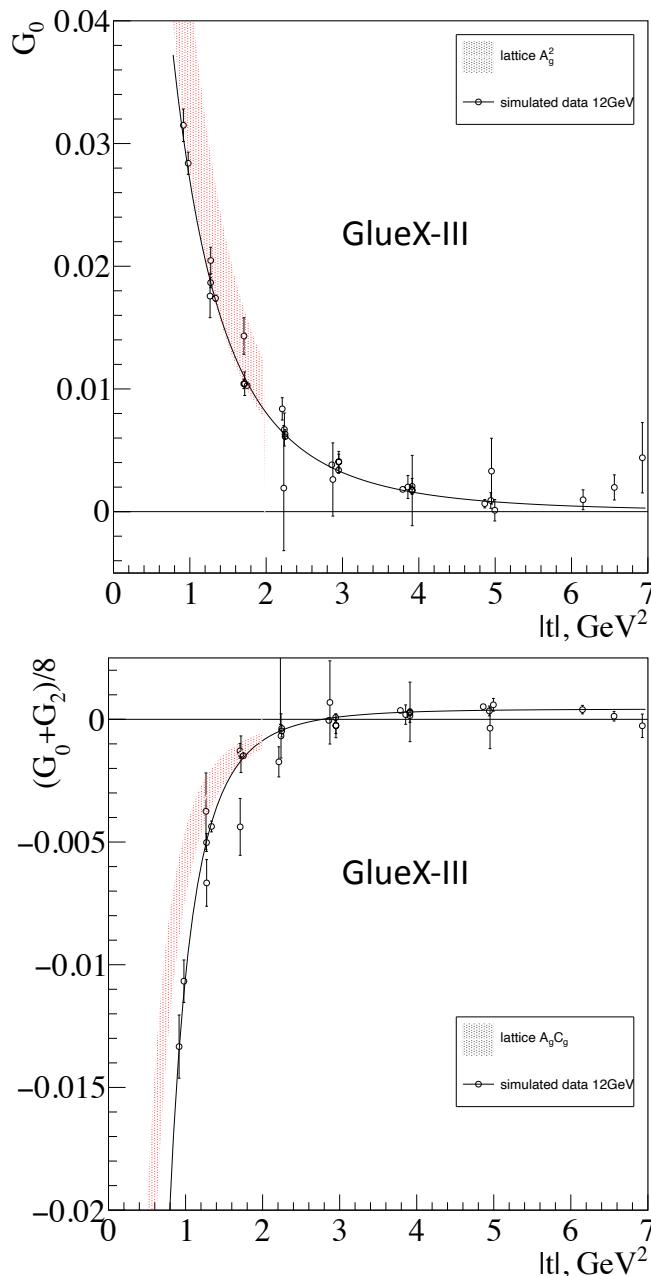


$$\text{asymmetry} = \frac{2}{P_\gamma} \frac{Y_{J/\psi}(0) - Y_{J/\psi}(90)}{Y_{J/\psi}(0) + Y_{J/\psi}(90)} =$$

$$= - (\rho_{1-1}^1 - \text{Im} \rho_{1-1}^2) \cos[2(\phi_{hel} - \Phi)]$$

$$\text{naturality} = (-1)^J P$$

Threshold charmonium photoproduction at JLab22 with GlueX



- Anticipated results for the extracted gluon Form Factors, assuming skewness scaling is valid
- Data randomized around a fit of the current data