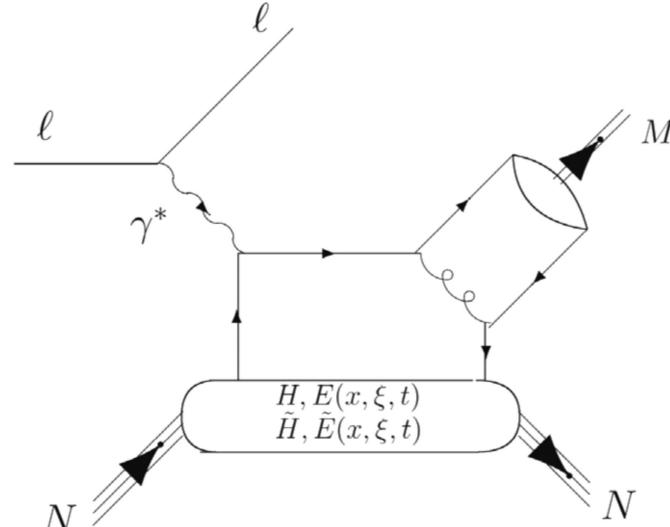
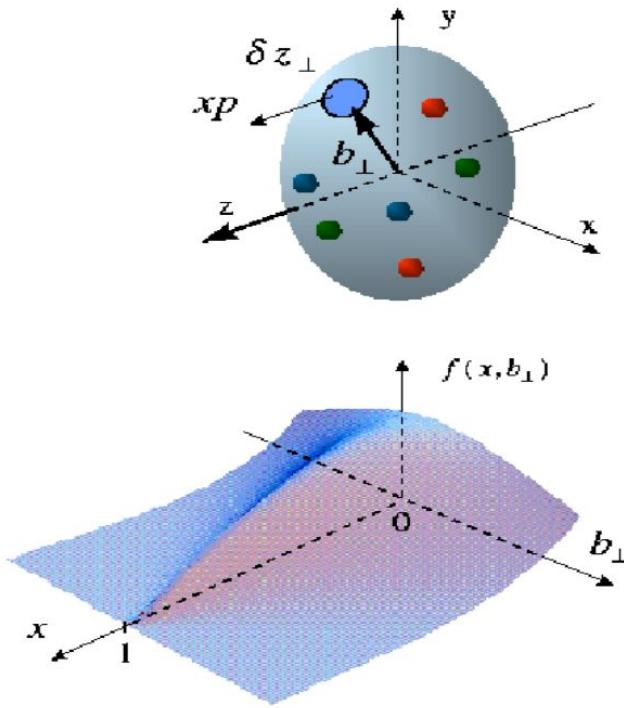

Overview of exclusive meson electroproduction program with CLAS12

Towards improved hardron femtography with hard exclusive reactions
August 10th, 2023

Andrey Kim
(University of Connecticut)

GPDs and exclusive meson electroproduction

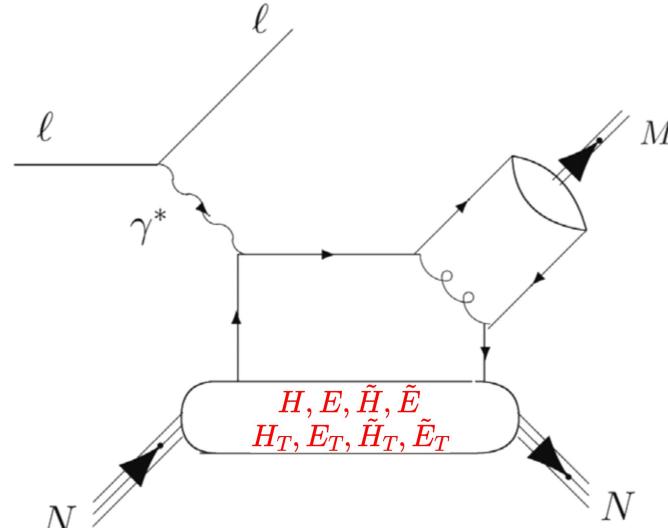
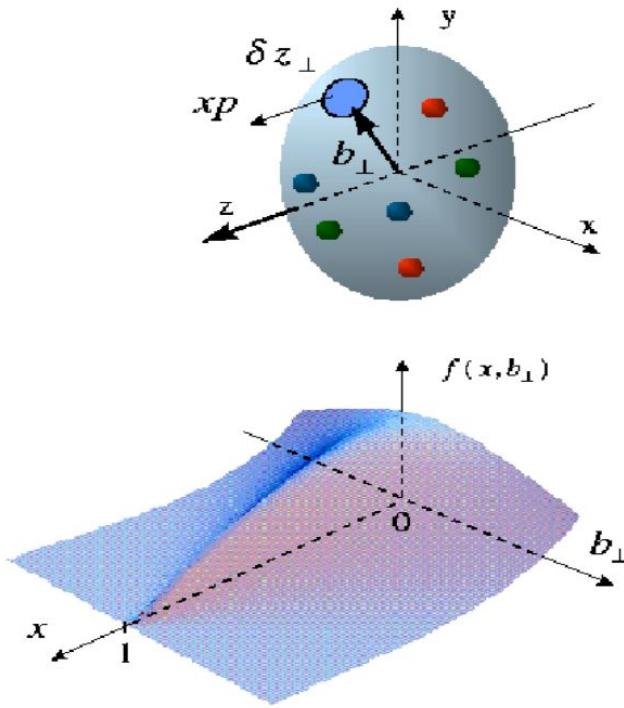
- 4 chiral-even GPDs: H, E, H, E
- 4 chiral-odd GPDs: $H_T, E_T, \tilde{H}_T, \tilde{E}_T$



$$\langle F \rangle = \sum_{\lambda} \int_{-1}^1 dx \mathcal{H}_{0\lambda,\mu\lambda}(x, \xi, Q^2, t) F(x, \xi, t)$$

GPDs and exclusive meson electroproduction

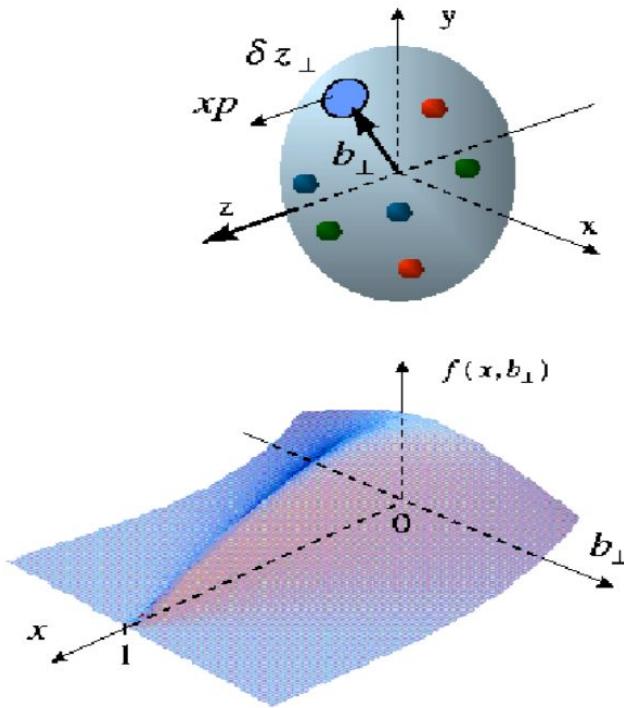
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$$\langle F \rangle = \sum_{\lambda} \int_{-1}^1 dx \mathcal{H}_{0\lambda,\mu\lambda}(x, \xi, Q^2, t) F(x, \xi, t)$$

GPDs and exclusive meson electroproduction

- 4 chiral-even GPDs: H, E, H, E
- 4 chiral-odd GPDs: $H_T, E_T, \tilde{H}_T, \tilde{E}_T$



Meson	GPD flavor composition
π^+	$\Delta u - \Delta d$
π^0	$2\Delta u + \Delta d$
η	$2\Delta u - \Delta d$
<hr/>	
H, E	ρ^0
H_T, \bar{E}_T	ρ^+
	ω

Link GPDs to PDF and Form Factors

PDFs:

in the forward limit

$$\xi = t = 0:$$

$$H^q(x, 0, 0) = q(x)$$
$$\tilde{H}^q(x, 0, 0) = \Delta q(x)$$

Form Factors:

$$\int dx H^q(x, \xi, t) = F_1(t)$$

$$\int dx E^q(x, \xi, t) = F_2(t)$$

$$\int dx \tilde{H}^q(x, \xi, t) = G_A(t)$$

$$\int dx \tilde{E}^q(x, \xi, t) = G_P(t)$$

Link GPDs to Tensor charge and Transversity

PDFs:

- Proton tensor charge

$$\delta_T^u = \int dx H_T^u(x, \xi, t=0)$$

$$\delta_T^d = \int dx H_T^d(x, \xi, t=0)$$

- Density of transversely polarized quarks
in an unpolarized proton in the transverse plane

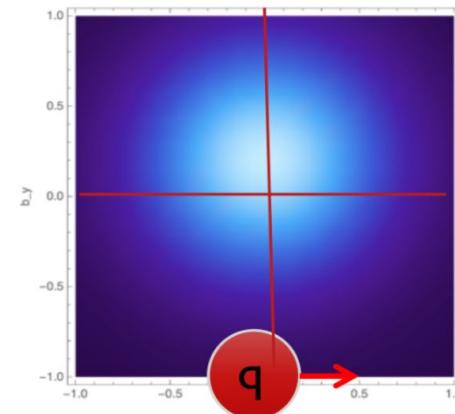
$$\delta(x, \vec{b}) = \frac{1}{2} [H(x, \vec{b}) - \frac{b_y}{m} \frac{\partial}{\partial b^2} \bar{E}_T(x, \vec{b})]$$

Form Factors:

- Proton anomalous tensor magnetic moment

$$\kappa_T^u = \int dx \bar{E}_T^u(x, \xi, t=0)$$

$$\kappa_T^d = \int dx \bar{E}_T^d(x, \xi, t=0)$$



CLAS12 to map exclusive reactions in a wide kinematic space

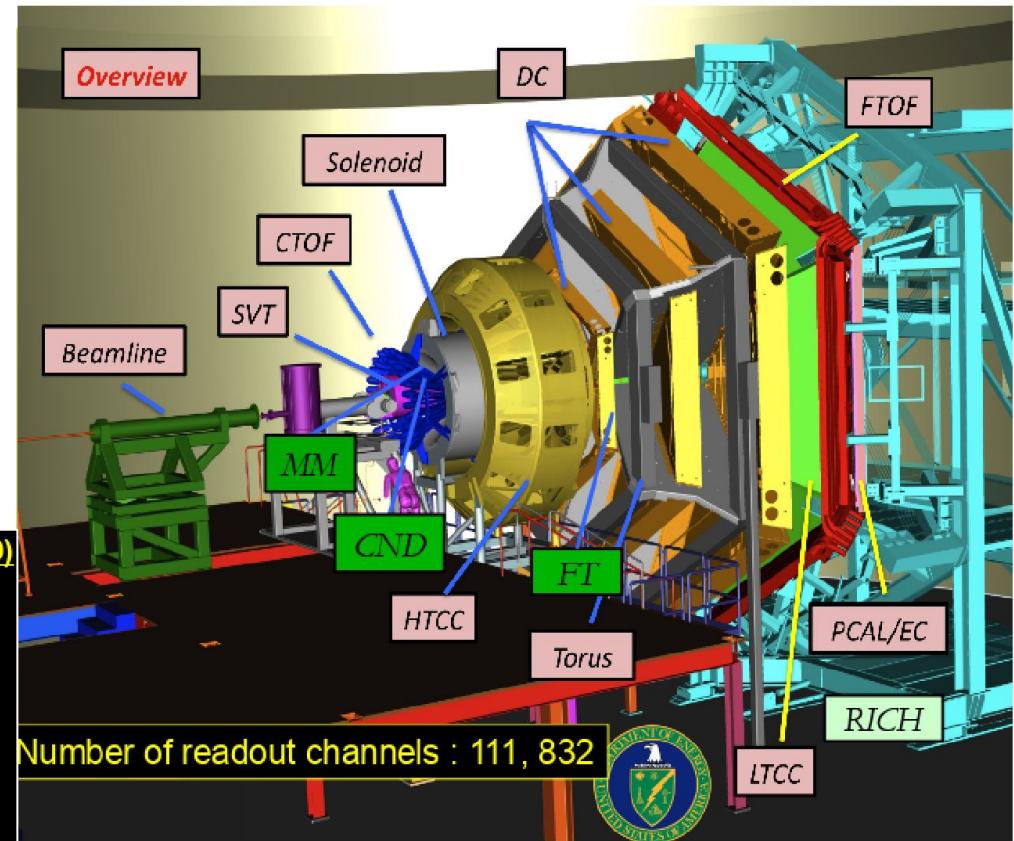
- CEBAF Large Acceptance Spectrometer
- High luminosity frontier
- 86 % electron beam longitudinal polarization
- Unpolarized and polarized fixed targets
- Comprehensive detection system
- Access to the Q^2 up to 10 GeV^2

Forward Detector (FD)

- TORUS magnet
- HT Cherenkov Counter
- Drift chamber system
- LT Cherenkov Counter
- Forward ToF System
- Pre-shower calorimeter
- E.M. calorimeter
- Forward Tagger
- RICH detector

Central Detector (CD)

- Solenoid magnet
- Silicon Vertex Tracker
- Central Time-of-Flight
- Central Neutron Det.
- MicroMegas

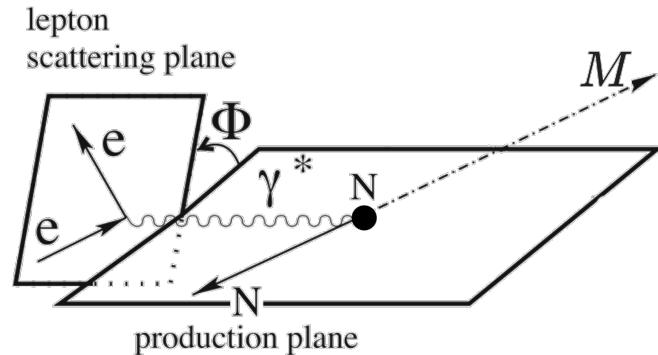


Experimental observables for exclusive meson electroproduction

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\Phi} = \Gamma(Q^2, x_B, E)$$

$$\begin{aligned} & \frac{1}{2\pi} \left\{ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right. \\ & + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\Phi) + \sqrt{\epsilon(2\epsilon+1)} \frac{d\sigma_{LT}}{dt} \cos(\Phi) \\ & \left. + \lambda \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{LT'}}{dt} \sin(\Phi) \right\} \end{aligned}$$

where λ is the helicity state of the incident electron beam



Access to GPDs

PHYSICAL REVIEW D **84**, 034007 (2011)

Flexible parametrization of generalized parton distributions from deeply virtual Compton scattering observables

Gary R. Goldstein,^{1,*} J. Osvaldo Gonzalez Hernandez,^{2,†} and Simonetta Liuti^{2,‡}

¹Department of Physics and Astronomy, Tufts University, Medford, Massachusetts 02155, USA

²Department of Physics, University of Virginia, Charlottesville, Virginia 22901, USA

(Received 16 February 2011; published 5 August 2011)

Eur. Phys. J. A (2011) **47**: 112
DOI 10.1140/epja/i2011-11112-6

THE EUROPEAN
PHYSICAL JOURNAL A

Regular Article – Theoretical Physics

Transversity in hard exclusive electroproduction of pseudoscalar mesons

S.V. Goloskokov^{1,a} and P. Kroll^{2,3,b}

Unpolarized beam and target within GK formalism

$$\sigma_L \sim \left\{ (1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

$$\sigma_T \sim \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\sigma_{LT} \sim \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\sigma_{TT} \sim \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

Access to GPDs

PHYSICAL REVIEW D 84, 034007 (2011)

Flexible parametrization of generalized parton distributions from deeply virtual Compton scattering observables

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THE EUROPEAN
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Regular Article – Theoretical Physics

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Unpolarized beam and target within GK formalism

Longitudinally polarized beam and longitudinally polarized target

$$A_{LU}^{\sin(\phi)} \sigma_0 \sim -\sqrt{\epsilon(1-\epsilon)} \sqrt{-t'} \operatorname{Im} \left[\langle \bar{E}_T \rangle^* \langle \tilde{H}_{\text{eff}} \rangle + \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

$$A_{LL}^{\text{const}} \sigma_0 \sim \sqrt{1-\epsilon^2} | \langle H_T \rangle |^2$$

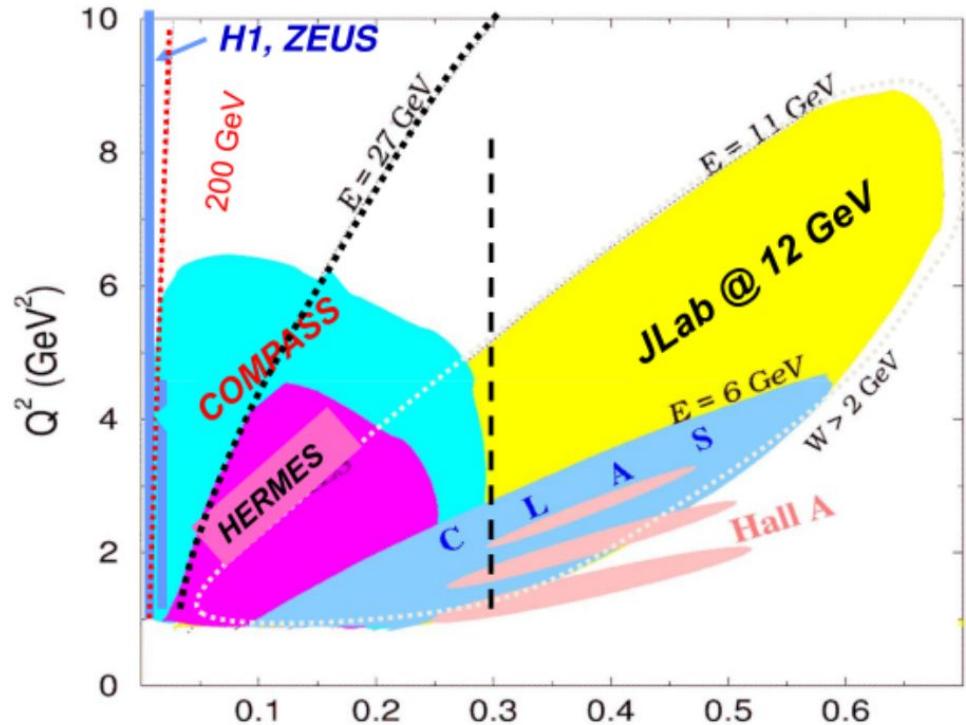
$$A_{LL}^{\cos(\phi)} \sigma_0 \sim -\sqrt{\epsilon(1-\epsilon)} \sqrt{-t'} \operatorname{Re} \left[2 \langle \bar{E}_T \rangle^* \langle \tilde{H}_{\text{eff}} \rangle + \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

within GK formalism

Measurements of exclusive meson production with CLAS

CLAS data:

- Phys. Rev. C63: 065205, 2001 (ϕ)
- Phys. Lett. B605: 256-264, 2005 (ρ^0)
- Eur. Phys. J. A24: 445-458, 2005 (ω)
- Phys. Rev. C78: 025210, 2008 (ϕ)
- Eur. Phys. J. A39: 5-31, 2009 (ρ^0)
- Phys. Rev. Lett. 109, 112001 (2012) (π^0)
- Phys. Rev. C 95, 035207 (2017) (η)
- Phys. Rev. C 95, 035206 (2017) (π^0)
- Phys. Rev. C 95, 035202 (2017) (π^+)
- Phys. Lett. B 768, 168 (2017) (π^0)
- Phys. Lett B. 789, 426 (2019) (η)



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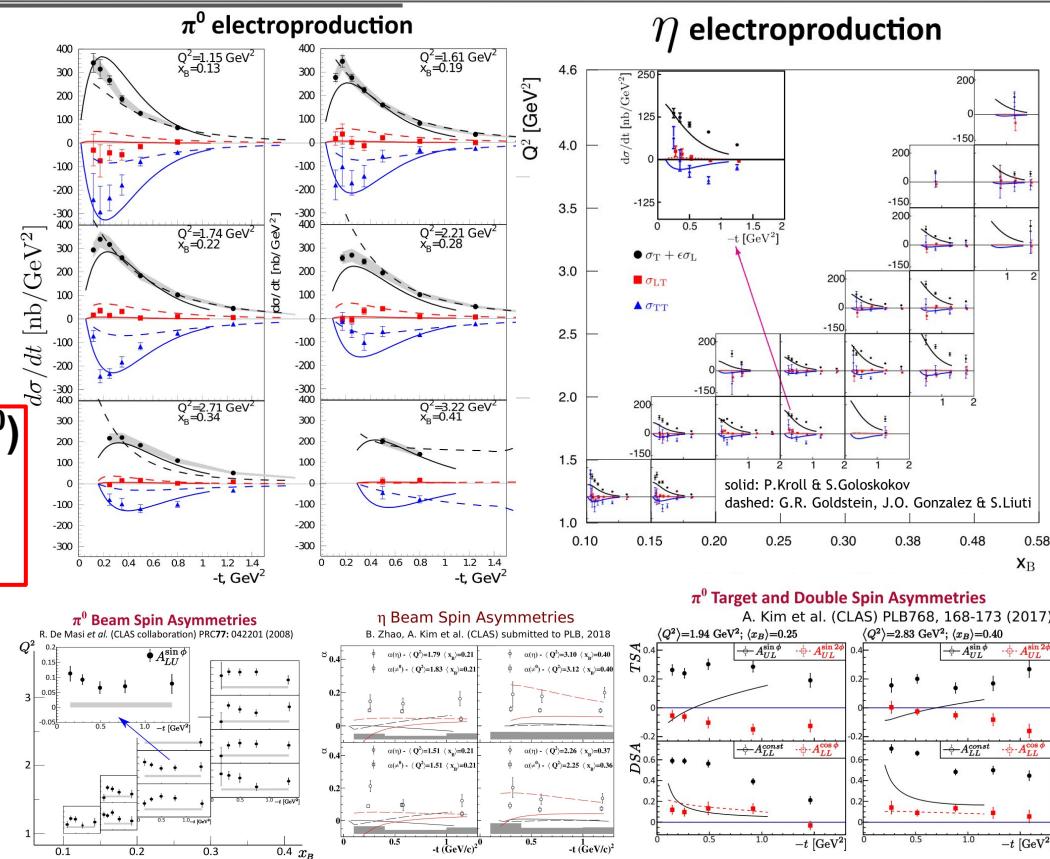
Phys. Rev. C 95, 035207 (2017) (η)

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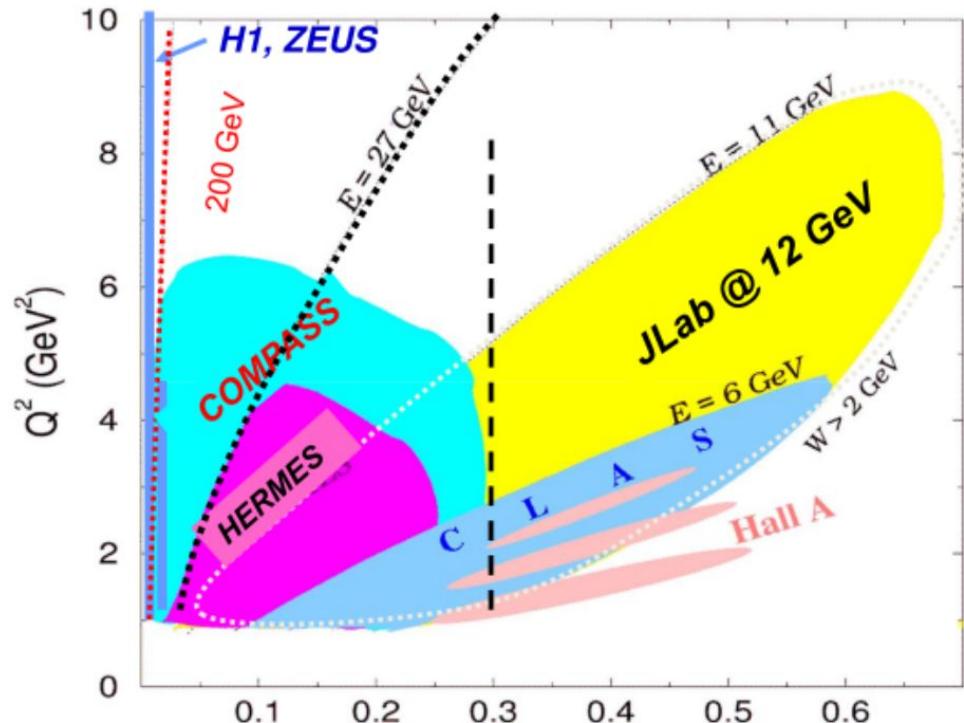
Phys. Lett B. 789, 426 (2019) (η)



Measurements of exclusive meson production with CLAS

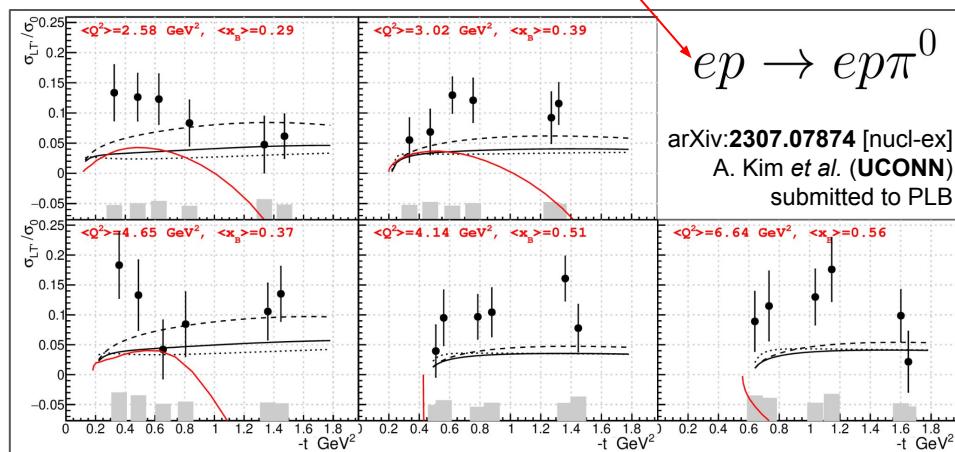
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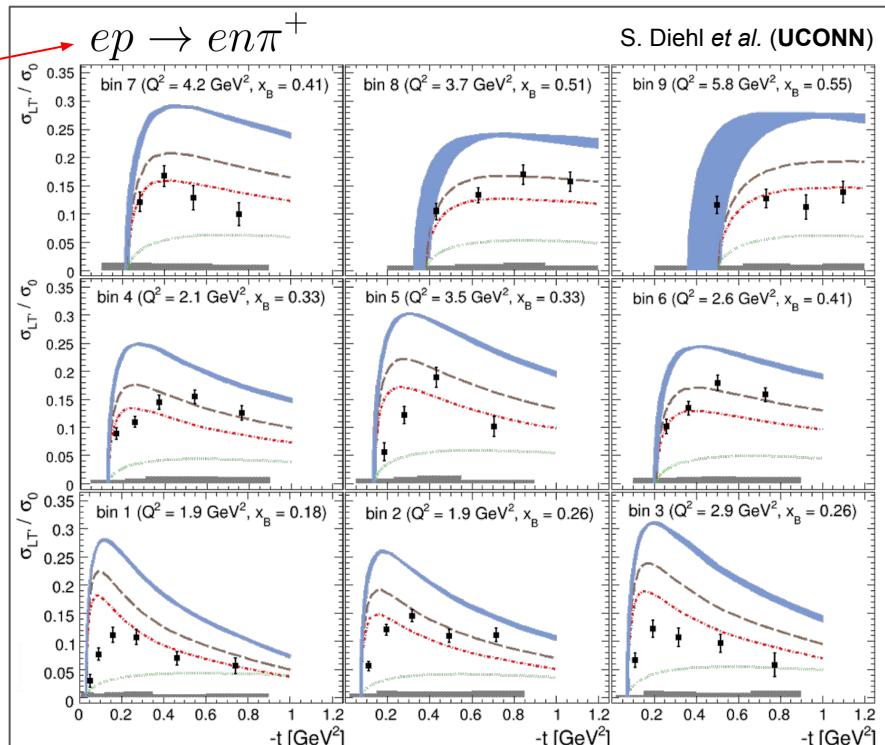


Pseudoscalar meson electroproduction with CLAS12

$$\sigma_{LT'} = \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \times \\ \times \text{Im} [\langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle]$$

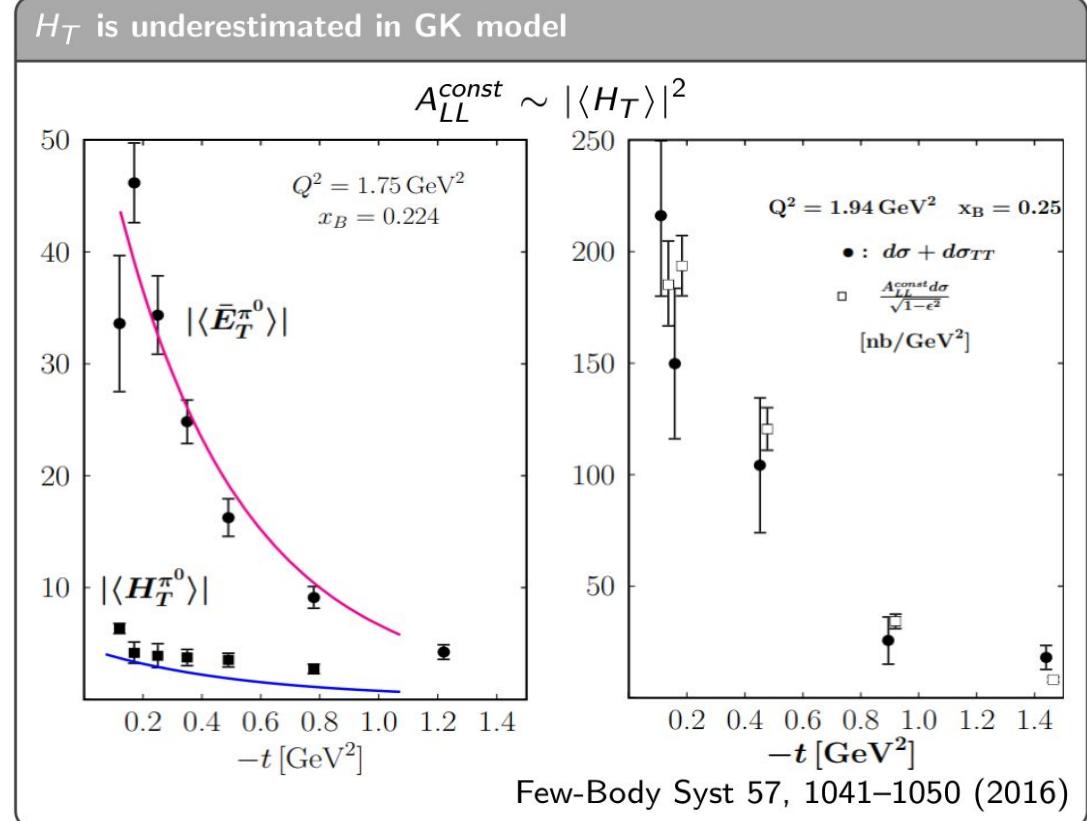
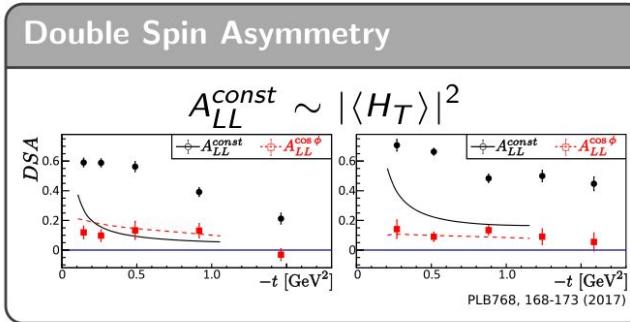
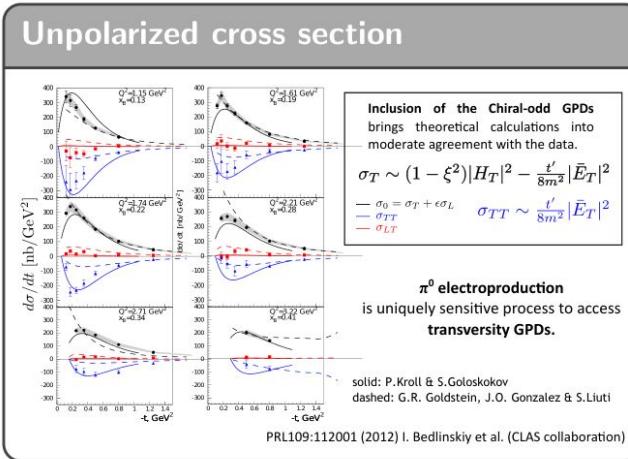


- active work on cross-section extraction by **R. Johnston (MIT)**
- active work on η beam spin asymmetry and cross-section extraction **A. Kim (UCONN)**



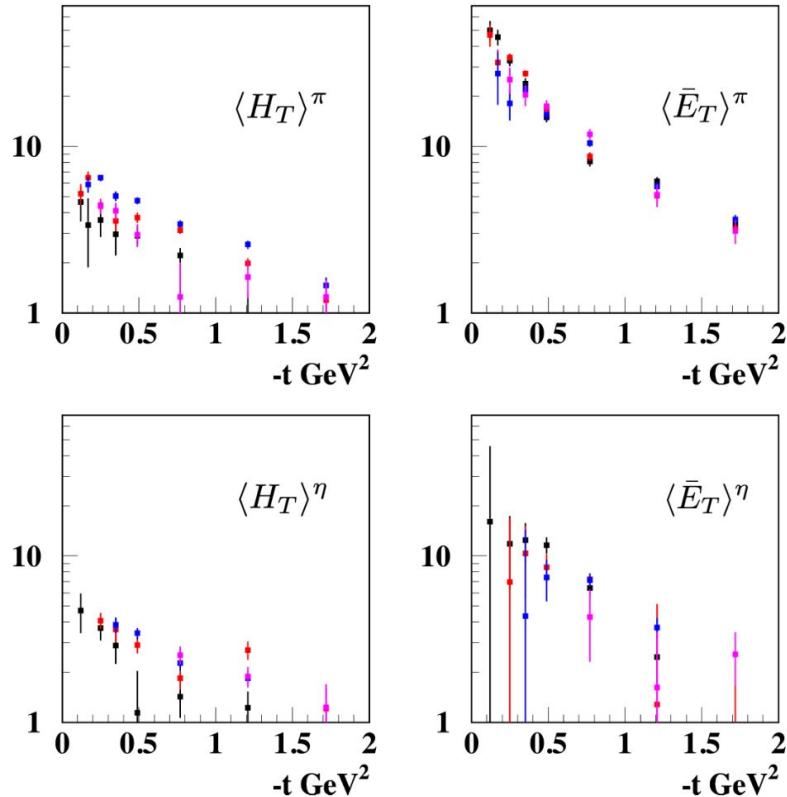
Phys. Rev. Lett. 125, 182001 – Published 28 October 2020

GPD insight (in progress)

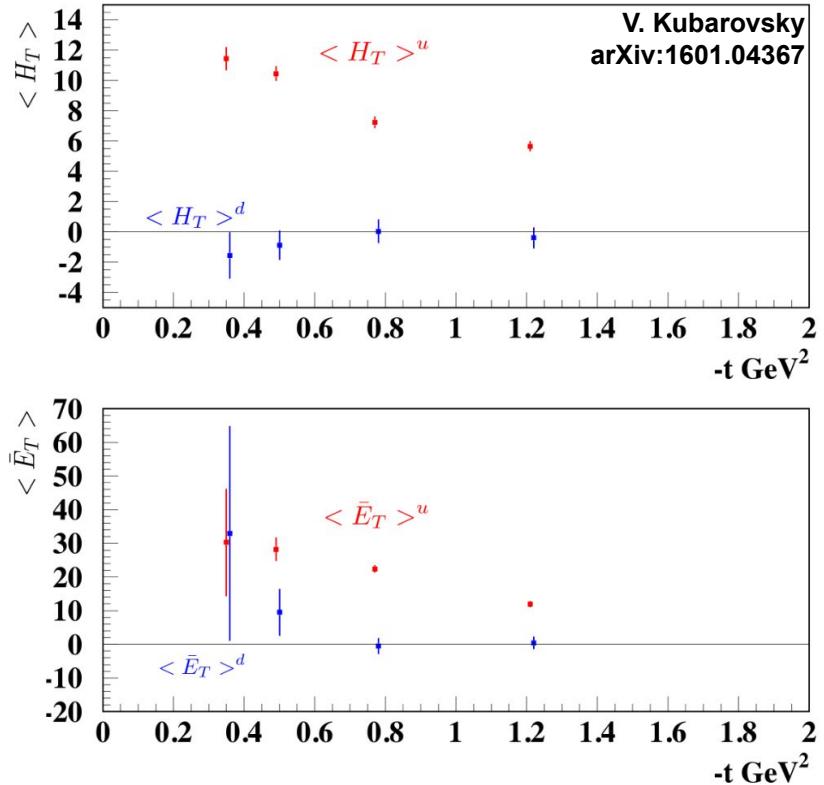


GPD insight (in progress)

Generalized Form Factors



Quark flavor decomposition



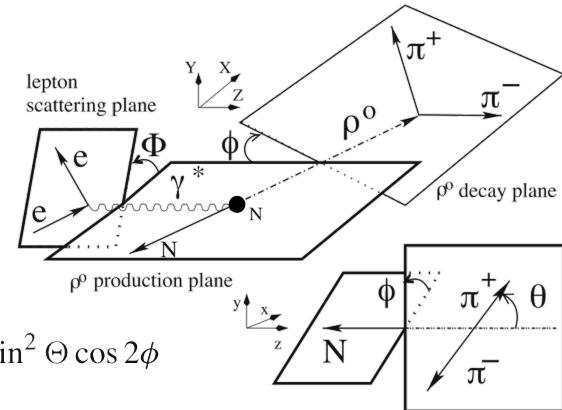
Vector meson production to access Spin Density Matrix Elements (SDME)

$$\frac{d\sigma}{d\phi \, d\Phi \, d\Theta \, dQ^2 \, dx_B \, dt} = \Gamma(Q^2, x_B, E) \frac{1}{2\pi} \left\{ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\} \mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta)$$

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta),$$

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi (r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi) \\ & - \epsilon \sin 2\Phi (\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 + \epsilon) \cos \Phi (r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi) \\ & \left. + \sqrt{2}\epsilon(1 + \epsilon) \sin \Phi (\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi) \right], \end{aligned}$$

$$\begin{aligned} \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} [\sqrt{1 - \epsilon^2} (\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 - \epsilon) \cos \Phi (\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 - \epsilon) \sin \Phi (r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi)] \end{aligned}$$



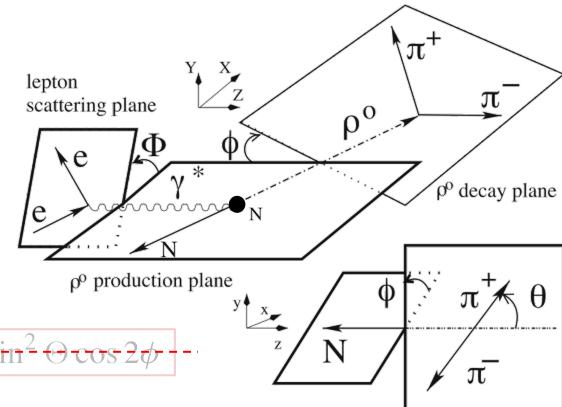
Vector meson production to access Spin Density Matrix Elements (SDME)

- Terms that are cancelled after integration over decay plane ϕ
- Terms that become constant after integration over decay plane Θ

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta),$$

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi [r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta] - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi) \\ & - \epsilon \sin 2\Phi (\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 + \epsilon) \cos \Phi [r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta] - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi) \\ & \left. + \sqrt{2}\epsilon(1 + \epsilon) \sin \Phi (\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi) \right], \end{aligned}$$

$$\begin{aligned} \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} [\sqrt{1 - \epsilon^2} (\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 - \epsilon) \cos \Phi (\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2}\epsilon(1 - \epsilon) \sin \Phi [r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta] - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi)] \end{aligned}$$



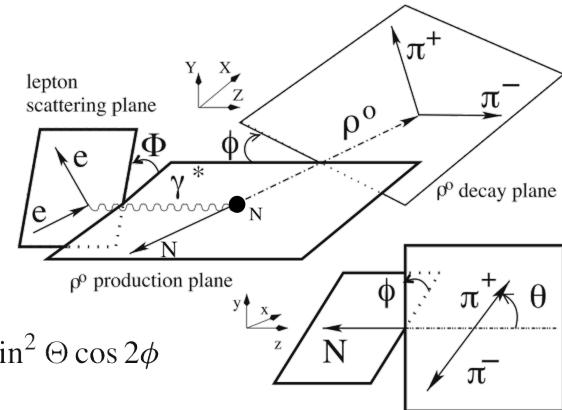
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$$\begin{aligned} \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} [\sqrt{1-\epsilon^2} (\sqrt{2} \text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi (\sqrt{2} \text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi) \\ & + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi (r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2} \text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi)] \end{aligned}$$



From SDME to GPDs

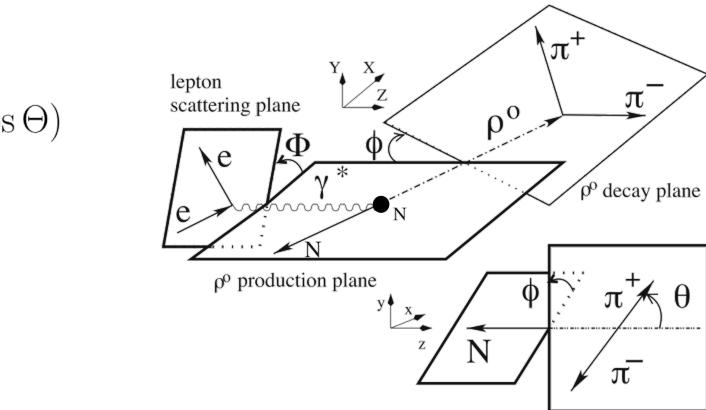
$$\frac{d\sigma}{d\phi \, d\Phi \, d\Theta \, dQ^2 \, dx_B \, dt} = \Gamma(Q^2, x_B, E) \frac{1}{2\pi} \left\{ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\} \mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta)$$

After simplifications from Eur. Phys. J. C (2014):

$$r_{00}^1 \sigma_0 \sim |\bar{E}_T|^2$$

$$r_{00}^5 \sigma_0 \sim \text{Re} [\langle \bar{E}_T \rangle \langle H \rangle + \langle H_T \rangle \langle E \rangle]$$

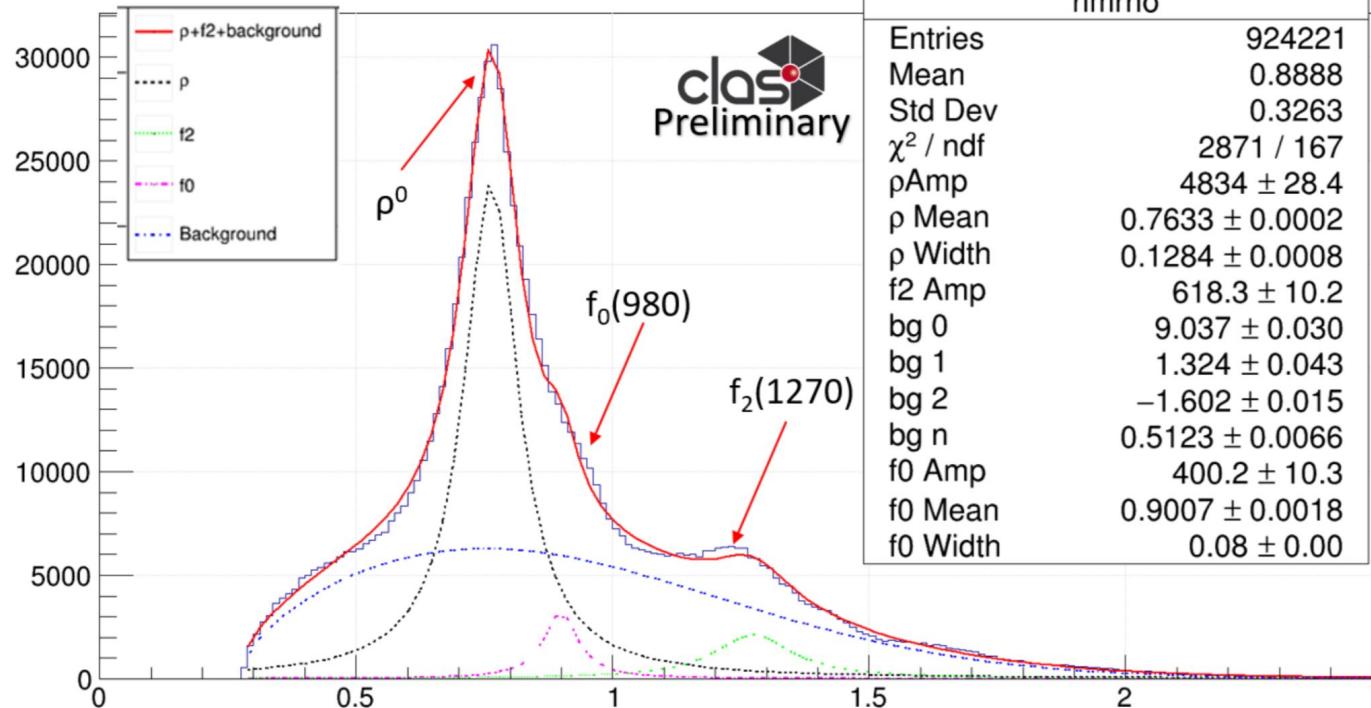
$$r_{00}^8 \sigma_0 \sim \text{Im} [\langle \bar{E}_T \rangle \langle H \rangle + \langle H_T \rangle \langle E \rangle]$$



$$R = \frac{\sigma_L}{\sigma_T} = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$$

Vector meson electroproduction with CLAS12: $ep \rightarrow epp\rho^0 \rightarrow e\pi^+\pi^-(P)$

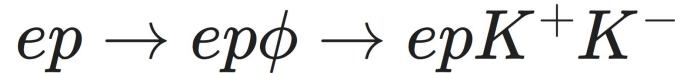
$$\sigma_{LT'} \sim r_{00}^8 \sim \text{Im} [\langle H_T \rangle^* \langle E \rangle + \langle \bar{E}_T \rangle^* \langle H \rangle]$$



parameter	PDG value
ρ^0 mass M_ρ (MeV)	≈ 770
ρ^0 width Γ_ρ (MeV)	≈ 150
f_0 mass M_{f_0} (MeV)	≈ 980
f_0 width Γ_{f_0} (MeV)	40-100
f_2 mass M_{f_2} (MeV)	≈ 1275
f_2 width Γ_{f_2} (MeV)	≈ 185

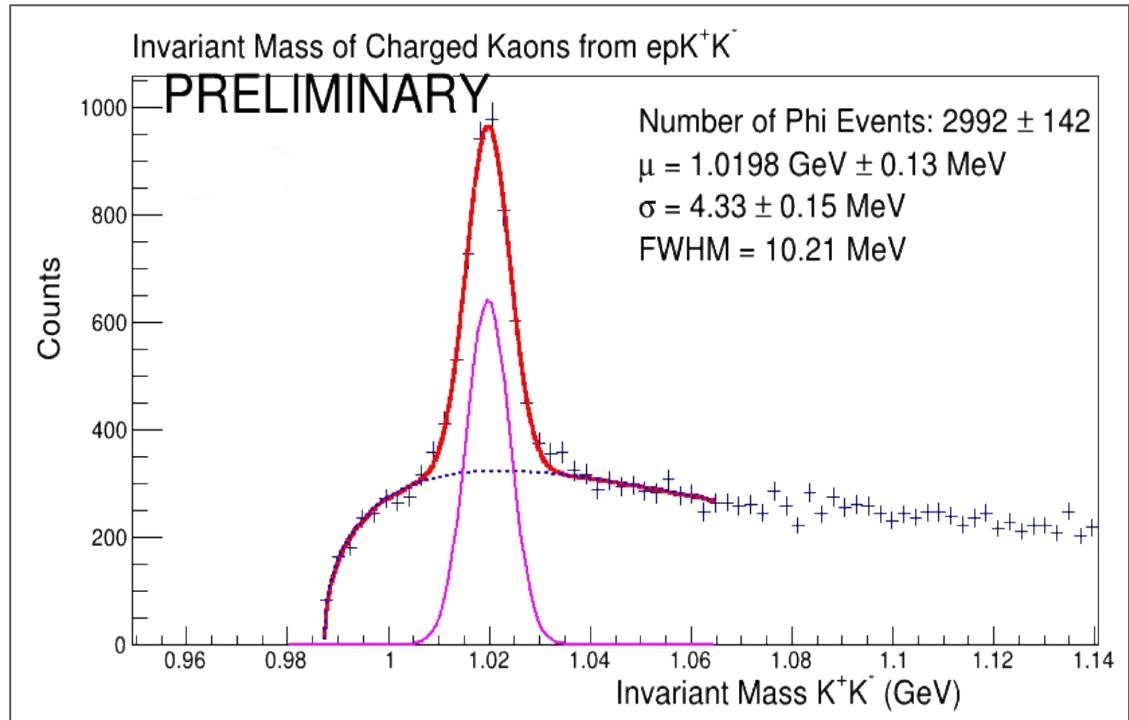
$$\begin{aligned} \frac{dN}{dM_{\pi\pi}} = & BW_{\rho^0}(M_{\pi\pi}) \\ & + BW_{f_0}(M_{\pi\pi}) \\ & + BW_{f_2}(M_{\pi\pi}) \\ & + BG_{ARGUS}(M_{\pi\pi}) \end{aligned}$$

Vector meson electroproduction with CLAS12



Access gluonic GPDs:

- ϕ ($s\bar{s}$) - low $|t|$ measurements where GPDs are relevant
- J/ψ ($c\bar{c}$) - measure the t -dependence of the differential cross section of J/ψ photoproduction*
- active work on φ cross-section measurements by P. Moran (MIT)
- active work on J/ψ cross-section measurements by J. Newton (JLab)



analysis by B. Clary (**UCONN**)

- The exclusive meson electroproduction program at CLAS and CLAS12 has produced important measurements which are already improving our knowledge of Generalized Parton Distributions.
- Stay tuned for the new measurements of **pseudoscalar** and **vector** exclusive meson production channels from CLAS12 data.

THANK YOU