

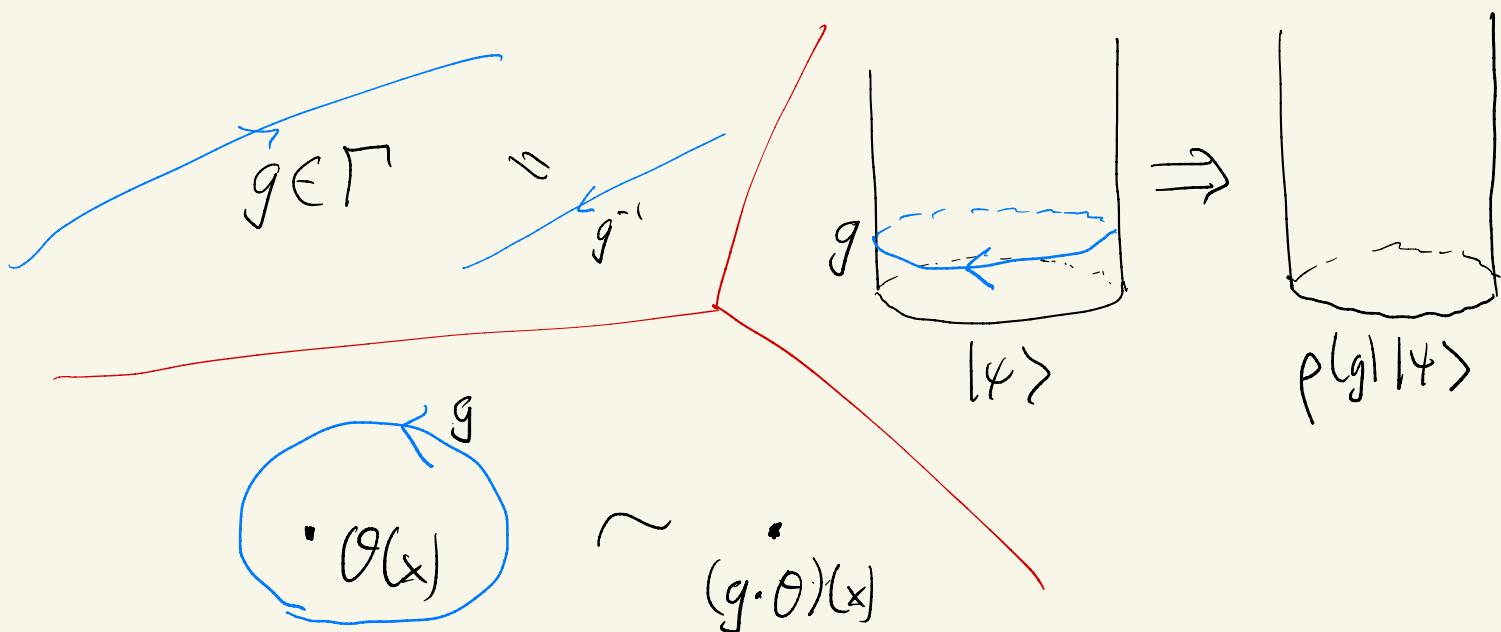
Decomposition in orbifolds with discrete torsion

Decomposition workshop, May 22, 2021

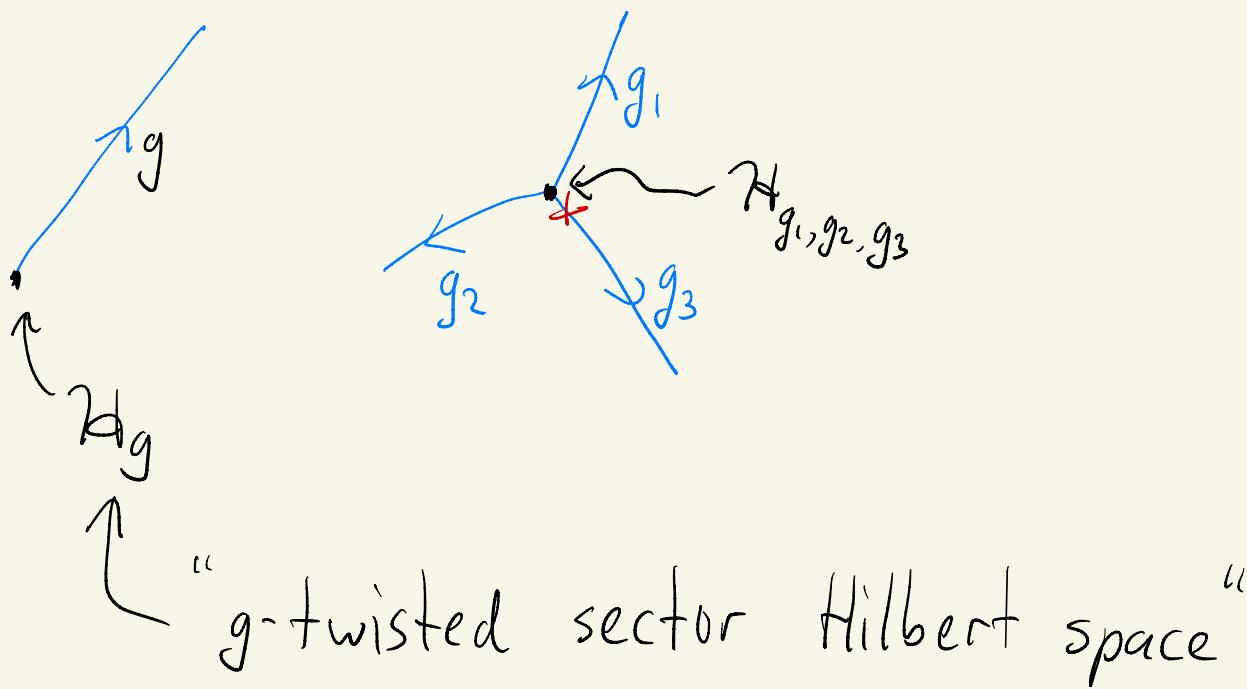
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based on: 2101.11619 w/ E. Sharpe
T. Vandermeulen

- Start with a 2D CFT with a (finite) global symmetry group Γ .
- We can think about the action of Γ in terms of Topological Defect Lines (TDLs)

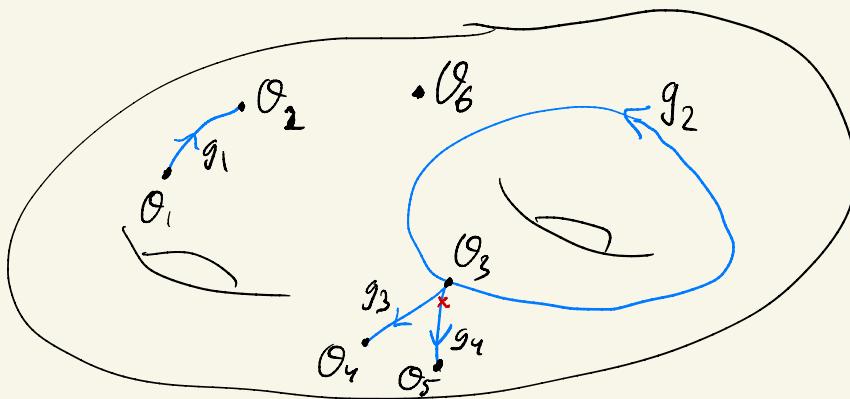


- TDLs can meet at junctions, and for each possible junction there will be a space of operators which can be inserted.

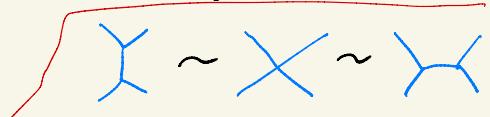


(see 1802.04445, Chang, Lin, Shao, Wang, + Yin)

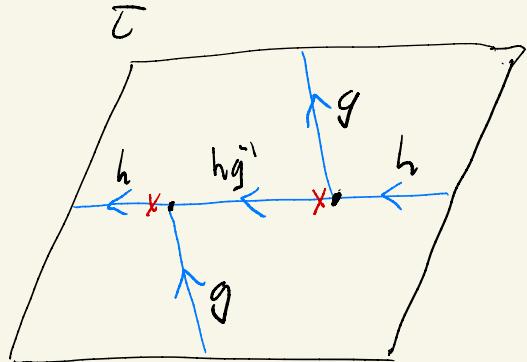
Correlation functions with TDLs:



- Lines can be deformed without changing corr. fn., but vertices can't unless inserted operator is weight zero.
- For effective symmetries, H_{g_1, \dots, g_n} contains weight zero states iff $g_1 g_2 \dots g_n = 1$.
- For anomaly-free Γ , can also freely resolve junctions



- “Partial Traces” $Z_{g,h}(\tau, \bar{\tau})$, $gh = hg$.



- Modular transformations: $Z_{g,h}(\tau+1) = Z_{g,hg^{-1}}(\tau)$, $Z_{g,h}(-\frac{1}{\tau}) = Z_{h,g^{-1}}(\tau)$

$$Z_{g,h}(\tau+1) = Z_{g,hg^{-1}}(\tau) = Z_{h,g^{-1}}(\tau)$$

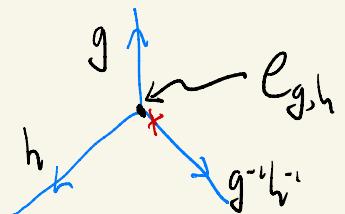
The equation shows three diagrams connected by equals signs, illustrating the modular transformation properties of the partition function.

- Left Diagram:** A genus-1 surface with boundary components τ (top), h (bottom left), and hg^{-1} (bottom right). Blue arrows point from the interior to the boundary at the midpoints of the h and hg^{-1} segments.
- Middle Diagram:** A genus-1 surface with boundary components τ (top), h (bottom left), and hg^{-1} (bottom right). The h segment is now labeled g above it. Blue arrows point from the interior to the boundary at the midpoints of the h and hg^{-1} segments.
- Right Diagram:** A genus-1 surface with boundary components τ (top), h (bottom left), and hg^{-1} (bottom right). The h segment is now labeled g above it. The hg^{-1} segment is now labeled hg^{-2} above it. Blue arrows point from the interior to the boundary at the midpoints of the h and hg^{-2} segments.

- A network of these TDLs can be thought of as a background Γ gauge field.
- If we want to gauge Γ (construct orbifold), we should sum over all background configurations.
- For instance, one-loop Γ -orbifold partition function,

$$Z^\Gamma(\tau, \bar{\tau}) = \frac{1}{|\Gamma|} \sum_{gh=hg} Z_{g,h}(\tau, \bar{\tau})$$

- Implicitly, we are inserting some weight zero vectors $e_{g,h} \in \mathbb{H}_{g,h, g^{-1}h^{-1}}$. We can make a new choice $\underline{e'_{gh} = w_{g,h} e_{g,h}}$
- Consistency ($\cancel{\text{X}} = \text{X}$) implies $\omega(g_2, g_3) \omega(g_1, g_2 g_3) = \omega(g_1 g_2, g_3) \omega(g_1, g_2)$
- Shifting $\omega \rightarrow \omega'(g, h) = \omega(g, h) \mu(g) \mu(h) \mu(gh)^{-1}$
 $\Rightarrow \omega \in H^2(\Gamma, U(1))$
- $Z_{g,h}(\tau) \rightarrow Z_{g,h}(\tau) \cdot \frac{\omega(g,h)}{\omega(h,g)}$ \rightsquigarrow discrete torsion.
 $\epsilon(g,h) = \frac{\omega(g,h)}{\omega(h,g)}$



- Now consider the case when a subgroup $K \subseteq \Gamma$ acts non-effectively (on local operators)
- K must be normal ($gkg^{-1}|4\rangle = g \cdot k(g^{-1}|4\rangle) = g \cdot (g^{-1}|4\rangle) = |4\rangle$
 $\Rightarrow gkg^{-1}$ acts trivially)

Define $G = \Gamma/K$.

SES: $1 \longrightarrow K \xrightarrow{\imath} \Gamma \xrightarrow{\pi} G \longrightarrow 1$

- Now $\mathcal{H}_{g_1, \dots, g_n}$ contains weight zero states iff $g_1 g_2 \cdots g_n \in K$.

The Γ -partial traces Z_{Y_1, Y_2}^Γ are expected to be related to G -partial traces. We'll take*

$$Z_{Y_1, Y_2}^\Gamma = Z_{\pi(Y_1), \pi(Y_2)}^G$$

$$\text{Then } Z^\Gamma = \frac{1}{|\Gamma|} \sum_{Y_1, Y_2 = Y_2 Y_1} \epsilon(Y_1, Y_2) Z_{Y_1, Y_2}^\Gamma$$

$$= \frac{1}{|G|} \sum_{g, h \in G} \left(\frac{1}{|K|} \sum_{\substack{Y_1 Y_2 = Y_2 Y_1 \\ \pi(Y_1) = g \\ \pi(Y_2) = h}} \epsilon(Y_1, Y_2) \right) Z_{g, h}^G$$

* until end of the talk

- Consider first $K = \Gamma$ (orbifold of a point)
 Each $\mathcal{H}_k \cong \mathcal{H}_1 = \mathcal{H}$. For every $\mathcal{O} \in \mathcal{H}$, maps under isomorphism to $\mathcal{O}^{(k)} \in \mathcal{H}_k$.

- If there is no discrete torsion ($\epsilon(r_1, r_2) = 1$), then

$$\mathbb{Z}^\Gamma = N \mathbb{Z}_{\text{irr}}^{\otimes \Gamma}, \quad \text{where} \quad N := \underbrace{\frac{1}{|\Gamma|} \sum_{r_1 r_2 = r_2 r_1} 1}_{= |\widehat{K}|} = \underbrace{\# \text{ of conj. classes of } K}_{= |\widehat{K}|},$$

where \widehat{K} is the set of isomorphism classes of irreducible representations of K .

- For example, if K is abelian, then $\mathcal{N} = \frac{|K|^2}{|K|} = |K|$.
- For an example of K nonabelian, take $K = S_3$.

Commuting pairs:

$(1, g)$	6×1
$((12), 1)$, $((12), (12))$	2×3
$((123), 1)$ $((123), (123))$, $((123), (132))$	$\frac{3 \times 2}{18}$

$$\Rightarrow Z^{S_3} = 3 Z_{\text{tot}}^{(1)} = 3 Z$$

- How does discrete torsion change the story?

Take $K = \Gamma = \mathbb{Z}_2 \times \mathbb{Z}_2 = \{1, a, b, ab\}$.

$$H^2(\mathbb{Z}_2 \times \mathbb{Z}_2, \omega(1)) \cong \mathbb{Z}_2.$$

Nontrivial element $\rightsquigarrow \epsilon(a, b) = \epsilon(a, ab) = \epsilon(b, a)$
 $= \epsilon(b, ab) = \epsilon(ab, a) = \epsilon(ab, b) = -1$,
 all other $\epsilon(g, h) = 1$.

$$Z_w^{\mathbb{Z}_2 \times \mathbb{Z}_2} = \underbrace{\frac{1}{4} \left(\sum_{y_1, y_2 \in \mathbb{Z}_2^2} \epsilon(y_1, y_2) \right)}_{10 \times (+1) + 6 \times (-1) = 4} Z_{1,1}^{\{1\}} = Z \rightsquigarrow \text{equivalent to the parent theory!}$$

In general, $N_w := \frac{1}{|K|} \sum_{k_1 k_2 \in K, k_1 = k_2} \epsilon(k_1, k_2) = |\hat{K}_w|$,

where \hat{K}_w is the set of isomorphism classes of irreducible w -projective representations of K .

- ω -projective representation of K for $\omega \in Z^2(k, U)$
 is a vector space V and a map $\rho: K \rightarrow GL(V)$
 satisfying $\rho(k_1) \rho(k_2) = \omega(k_1, k_2) \rho(k_1 k_2)$
- e.g. for $K = \mathbb{Z}_2 \times \mathbb{Z}_2$
 $\omega(a, b) = \omega(b, ab) = \omega(ab, a) = i, \quad \omega(b, a) = \omega(ab, b) = \omega(a, ab) = -i$

Up to isomorphism, only projective irrep is

$$\rho(1) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \rho(a) = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \rho(b) = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \rho(ab) = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

i.e. the Pauli matrices.

• What if K is not all of Γ ?

Then \mathbb{Z}^Γ should decompose into orbifolds by subgroups of G .

$$\begin{matrix} \{1, a\} \\ \text{``} \end{matrix} \quad \begin{matrix} \{1, a, b, ab\} \\ \text{``} \end{matrix} \quad \begin{matrix} \{\bar{1}, \bar{b}\} \\ \text{``} \end{matrix}$$

• Ex 1. $1 \rightarrow \mathbb{Z}_2 \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow \mathbb{Z}_2 \rightarrow 1$

i) No discrete torsion:

$$\mathbb{Z}^{\mathbb{Z}_2 \times \mathbb{Z}_2} = \frac{1}{4} \left(4 \mathbb{Z}_{\bar{1}, \bar{1}}^{\mathbb{Z}_2} + 4 \mathbb{Z}_{\bar{1}, \bar{b}}^{\mathbb{Z}_2} + 4 \mathbb{Z}_{\bar{b}, \bar{1}}^{\mathbb{Z}_2} + 4 \mathbb{Z}_{\bar{b}, \bar{b}}^{\mathbb{Z}_2} \right)$$

$$= 4 \mathbb{Z}^{\mathbb{Z}_2}$$

\rightsquigarrow 4 disjoint copies of the \mathbb{Z}_2 orbifold

ii) With d.t.: $\mathbb{Z}_w^{\mathbb{Z}_2 \times \mathbb{Z}_2} = \frac{1}{4} \left(4 \mathbb{Z}_{\bar{1}, \bar{1}}^{\mathbb{Z}_2} + (2-2) \left(\mathbb{Z}_{\bar{1}, \bar{b}}^{\mathbb{Z}_2} + \mathbb{Z}_{\bar{b}, \bar{1}}^{\mathbb{Z}_2} + \mathbb{Z}_{\bar{b}, \bar{b}}^{\mathbb{Z}_2} \right) \right)$

$= \mathbb{Z}$ \rightsquigarrow one copy of the parent theory.

$$\cdot \text{Ex. 2. } 1 \rightarrow \mathbb{Z}_2 \rightarrow D_4 \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow 1.$$

$$H^2(D_4, \mathcal{U}(1)) \cong \mathbb{Z}_2$$

i) without d.t.: $\mathbb{Z}^{D_4} = \mathbb{Z}^{\mathbb{Z}_2 \times \mathbb{Z}_2} + \mathbb{Z}_{\omega}^{\mathbb{Z}_2 \times \mathbb{Z}_2}$ discrete torsion
in $G = \mathbb{Z}_2 \times \mathbb{Z}_2$

ii) with d.t.: $\mathbb{Z}_{\omega}^{D_4} = \mathbb{Z}^{\mathbb{Z}_2} \subset \mathbb{Z}_2$ subgroup of $G = \mathbb{Z}_2 \times \mathbb{Z}_2$

$$\underline{\text{Ex. 3}} \quad 1 \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow S_4 \rightarrow S_3 \rightarrow 1$$

$$H^2(S_4, \mathcal{U}(1)) \cong \mathbb{Z}_2$$

i) without d.t.: $\mathbb{Z}^{S_4} = \mathbb{Z} + \mathbb{Z}^{S_3}$

ii) with d.t.: $\mathbb{Z}_{\omega}^{S_4} = \mathbb{Z}^{S_3}$

- Conjectured algorithm:

- Given $| \rightarrow K \xrightarrow{\gamma} \Gamma \xrightarrow{\pi} G \rightarrow |$

and $\omega \in H^2(\Gamma, U(1))$:

- Construct $\hat{K}_{\gamma^*\omega}$

- Construct an action of G on $\hat{K}_{\gamma^*\omega}$:

$$(L_g \varphi)(k) = \frac{\omega(s^{-1}k, s)}{\omega(s, s^{-1}k)} \varphi(s^{-1}ks) \quad s = s(g)$$

Not a group action on φ 's, but is an isomorphism classes.

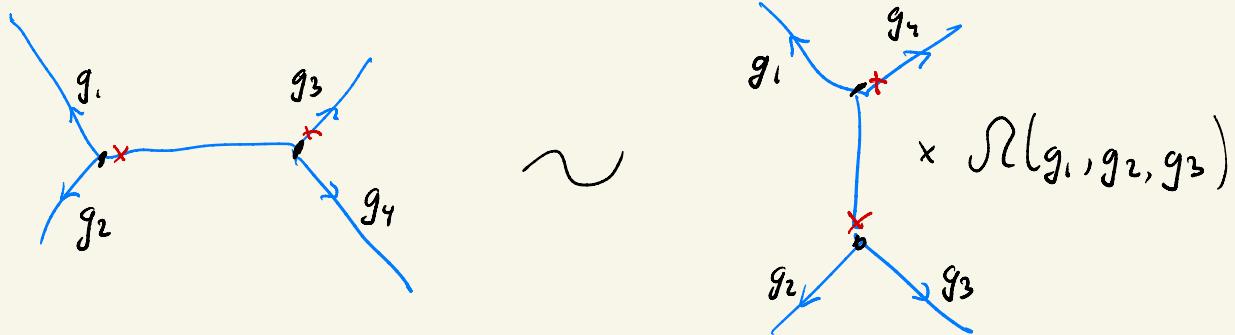
- Decompose $\hat{K}_{\gamma^*\omega}$ into G -orbits with representatives $[\rho_a]$ and $G_a = \text{Stab}([\rho_a]) \subseteq G$.

- Each orbit gives a G_a orbifold in decomposition with discrete torsion $w_a \in \mathbb{Z}^2(G_a, U(1))$
- w_a is determined by constructing from ρ_a a proj. rep $\tilde{\rho}_a$ of Γ with $\tilde{w}_a = \frac{\omega}{\pi^* w_a}$
(see paper for details)

This algorithm works in all examples we checked.

Anomalies. For a CFT with symmetry G , the potential anomaly lives in $H^3(G, \mathbb{U}(1))$.

How do we see its effects?



Modular transformations of $Z_{g,h}$ pick up phases.

$$G = \mathbb{Z}_2, \quad H^3(\mathbb{Z}_2, U(1)) \cong \mathbb{Z}_2$$

$$Z_{1,0}(\tau+2) = - Z_{1,0}(\tau)$$

$$\text{Resolution: } 1 \longrightarrow \mathbb{Z}_2 \xrightarrow{\iota} \mathbb{Z}_4 \xrightarrow{\pi} \mathbb{Z}_2 \longrightarrow 1$$

$$[\pi^* \mathcal{S}] = 0 \in H^3(\mathbb{Z}_4, U(1))$$

$$Z_{r_1, r_2}^\Gamma = \underbrace{\epsilon(r_1, r_2)}_{\uparrow} Z_{\pi(r_1), \pi(r_2)}^G$$

not from discrete torsion here!

$$B(g, k) \in \underline{Z^1(G, H^1(K, U(1)))}$$

$$\epsilon(r_1, r_2) = \frac{B(\pi(r_2), s(\pi(r_1))^{-1}r_1)}{B(\pi(r_1), s(\pi(r_2))^{-1}r_2)}$$

Interpretation as a quantum symmetry.
generalization of

Thanks!