

Extracting the Nucleon Axial Form Factor from Lattice QCD using χ PT

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New Directions in Neutrino-Nucleus Scattering

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- **Nucleon Axial Form Factor,** $G_A(q^2)$

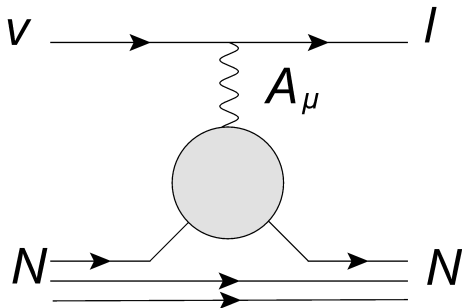
- fundamental property of the nucleon
- Information on the spins distribution
- $A_\mu^i(x) = \bar{q}(x) \gamma_\mu \gamma_5 \frac{\tau^i}{2} q(x)$
- $\langle N(p') | A_\mu^i | N(p) \rangle = \bar{u} \left\{ \gamma_\mu G_A(q^2) + \frac{q_\mu}{2m_N} G_P(q^2) \right\} \gamma_5 \frac{\tau^i}{2} u(p)$

- $G_A(q^2) = g_A \left[1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4) \right]$

- g_A and G_A dependence in q^2 are necessary in

ν oscillations experiments.

- $G_A(q^2)$ enters $\sigma_{\nu N}$. Input for MC and theoretical models.



- $G_A(q^2) = g_A \left[1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4) \right]$
- β decay: $g_A = 1.2723(23)$
(C. Patrignani et al. (PDG), Chin. Phys. C 40 (2016))
- $G_A(q^2)$ more challenging
 - Neutrino-induced charged-current QE scattering on deuteron:
 - Low statistics, unknown E_V
 - Old experiments: low accuracy
 - New experiments in nuclei \rightarrow problem: nuclear physics
 - Also experiments in: π electro-production and weak capture in muonic hydrogen



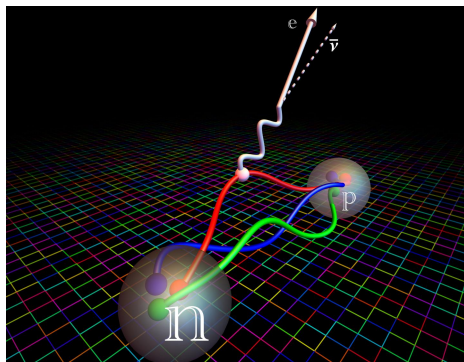
Figure: Fermilab

• Lattice

- several studies on $G_A(q^2)$
 - → some technical difficulties
- Experimental and lattice q^2 parametrisation:
 - dipole ansatz
 - ⇒ different results
 - z-expansion

• Chiral Perturbation Theory (χ PT)

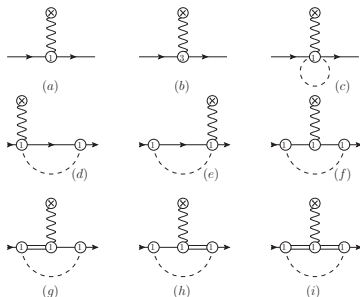
- QCD based parametrisation of q^2 and M_π dependencies
- χ PT is useful to extrapolate Lattice to the physical limit
- Fits to lattice can determine the low energy constants (LECs) values:
 - Lattice- χ PT synergy: M_π is "variable" in both theories



- χ PT is the effective field theory for QCD at low energy
 - Invariance under chiral symmetry: $G \equiv SU(n_f)_L \otimes SU(n_f)_R$
 - quarks, gluons \rightarrow hadronic degrees of freedom
 - Successful description of meson properties
- Baryon χ PT
 - Problem: $\underbrace{m_B \rightarrow 0}_{\chi^{\text{limit}}} \Rightarrow$ Power counting breaking (PCB)
 - \Rightarrow additional finite renormalisation: extended on mass-shell (EOMS) renormalisation
 - PCB terms absorbed by LECs
 - Covariance and analytic properties of loops preserved

G_A calculation in (B χ PT)

- Yao, Alvarez-Ruso, Vicente Vacas, PRD 96 (2017)
- Up to leading one-loop $\mathcal{O}(p^3)$
- Explicit $\Delta(1232)$
 - $\delta = m_\Delta - m_N \sim \mathcal{O}(p)$
- EOMS renormalisation
- Correct power counting and analytical properties
 \Rightarrow appropriate for chiral extrapolations
- G_A fit to combined lattice data



Combined fit

• G_A fit to combined lattice data

- 2017 fit data = PNDME¹ + "Cyprus"² + "Mainz"³
- 2021 fit data = 2017 (updating "Cyprus"⁴) + RQCD⁵ + PACS⁶
 - lattice spacing dependence corrected
- χ PT range:
 - $Q^2 < 0.36 \text{ GeV}^2$, $M_\pi < 400 \text{ MeV}$ ($M_\pi L \geq 3.5$)
- LECs g , d_{16} and d_{22} are determined by the fit
- Δ LECs (h_A , g_1) fixed to large- N_c
- **Good description: improved by Δ**

¹ R. Gupta et al., Phys. Rev. D 96 (2017)

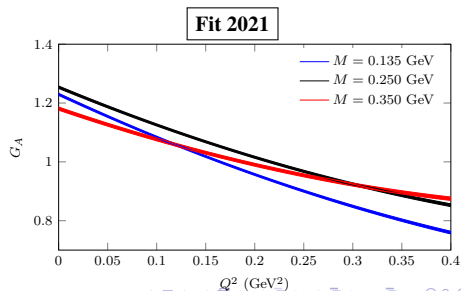
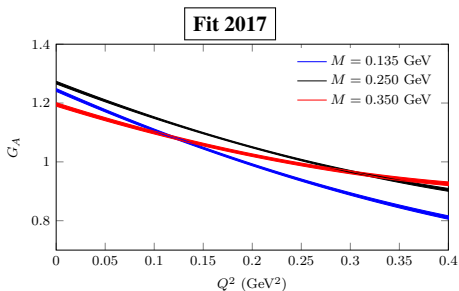
² C. Alexandrou et al., Phys. Rev. D 96 (2017)

³ S. Capitani et al., Int.J.Mod.Phys.A 34 (2019)

⁴ Phys. Rev. D 103 (2021)

⁵ G. S. Bali et al., JHEP 126 (2020)

⁶ Phys. Rev. D 102 (2020)



- $G_A(q^2) = g_A \left[1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4) \right]$

- | | Fit-2017 | Fit-2021 | Experimental* |
|-------|------------|------------|---------------|
| g_A | 1.2444(48) | 1.2295(37) | 1.2723(23) |

- Our fits: statistical errors alone

*(C. Patrignani et al. (PDG), Chin. Phys. C 40 (2016))

Axial radius, $\langle r_A^2 \rangle$

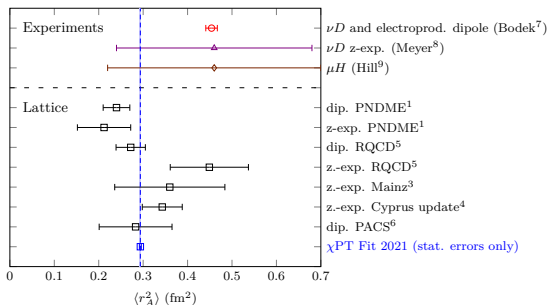
$$G_A(q^2) = g_A \left[1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4) \right]$$

Source	$\langle r_A^2 \rangle$
Fit-2017	0.2732(32)
Fit-2021	0.2937(22)
νD and electroprod. dipole (Bodek ⁷)	0.454(13)
νD z-exp. (Meyer ⁸)	0.46(22)

⁷ A. Bodek et al., Eur. Phys. J. C 53, 349 (2008)

⁸ A. S. Meyer et al., Phys. Rev. D 93, 113015 (2016)

⁹ R. J. Hill et al. Rept. Prog. Phys. 81 (2018)



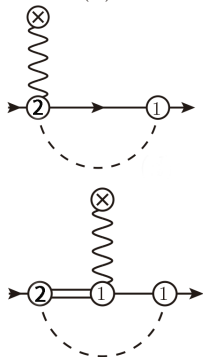
$G_A \mathcal{O}(p^4)$

- A more precise description of $G_A(q^2)$ and a determination of the chiral series truncation error $\implies \mathcal{O}(p^4)$ calculation

- We have calculated the new diagrams

\longrightarrow new $\mathcal{O}(p^2)$ LECs appear

- $\mathcal{L}_{\pi N}^{(2)} = \sum_i c_i \bar{\Psi} O_i \Psi \implies c_1, c_2, c_3, c_4$
- $\mathcal{L}_{\pi\Delta}^{(2)} = a_1 \bar{\Psi}_\mu^j \langle \chi_+ \rangle \Psi^{j\mu} + \text{h. c.}$
- $\mathcal{L}_{\pi N\Delta}^{(2)} = \sum_i \bar{\Psi}_\alpha^j b_i O^{aj} \Psi_N + \text{h. c.} \implies b_1, b_2, b_4, b_5$



+...

+ w. f. renormalisation

+ mass insertions

- g_A $\mathcal{O}(4)$ Δ -less
 - g_A lattice dataset*
 - The $\mathcal{O}(3)$ and $\mathcal{O}(4)$ Δ -less LECs (d_{16} and the c_i) extracted from $\pi N \rightarrow \pi N$ and $\pi N \rightarrow \pi\pi N$ (Siemens et al. PRC 96 (2017)) do not describe the g_A in the lattice.

* g_A dataset: "Mainz" (PRD 100 (2019)), RQCD⁵, PNDME (PRD 98 (2018)), CallLat (Nature 558 (2018))

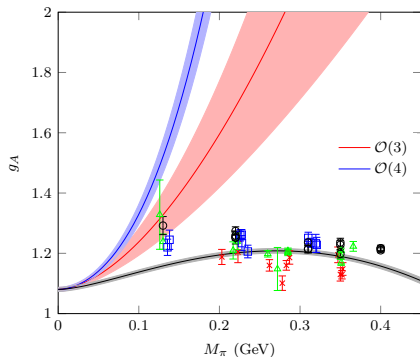


Figure: Gray line: our fit of the g_A $\mathcal{O}(3)$ to the lattice data (correcting a). Red line: g_A $\mathcal{O}(3)$ prediction of PRC 96 (2017) (their d_{16} , but our g). Blue line: g_A $\mathcal{O}(4)$ prediction of PRC 96 (2017) (their d_{16} and c_i , but our g). The bands show the errors (only statistical errors are considered for our fit). The points represent the lattice data.

- $g_A \mathcal{O}(4) \Delta$ -full fit:

- $g_A = 1.199 \pm 0.012$

- d_{16}

- key source of uncertainty in χ^2 PT
extrap. of nuclear observables

- $\pi N \rightarrow \pi\pi N \Rightarrow$

$$d_{16} = 2.98 \pm 1.00 \text{ GeV}^{-2\dagger}$$

- g_A is more sensitive to d_{16} than
 $\pi N \rightarrow \pi\pi N$

- our fit $\Rightarrow d_{16} = 0.24 \pm 0.18 \text{ GeV}^{-2}$
(stat. error only)

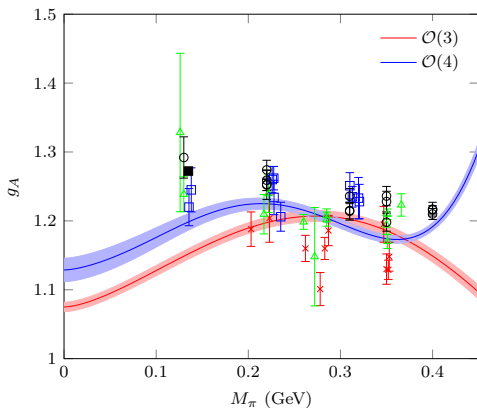


Figure: Red curve: $g_A \mathcal{O}(3) \Delta$ -full fit; blue curve: $g_A \mathcal{O}(4) \Delta$ -full fit. The points are the lattice data, except for the filled black square, which is the experimental determination.

[†]Siemens et al. PRC 96 (2017) (value converted to standard EOMS)

- $G_A(q^2)$ essential in ν oscillations.
- χ PT and Lattice complement one another
- The extraction of $\langle r_A^2 \rangle$ is challenging both from experiment and lattice
- Our combined fit $\mathcal{O}(p^3)$ Δ -full successfully describes the lattice data
 - $g_A = 1.2295(37)$, $\langle r_A^2 \rangle = 0.2937(22)$ (fm²)
 - No ad hoc parametrisation
- $\mathcal{O}(p^4)$ Δ -full fit $\Rightarrow d_{16} = 0.24 \pm 0.18$ GeV⁻²

Thanks for watching!

Any questions?

- Dipole ansatz: $G_A(q^2) = g_A \left(1 - \frac{q^2}{M_A^2}\right)^{-2}$
- z-exp.: $G_A(q^2) = \sum_k a_k z^k(q^2)$, with $z(q^2, t_{\text{cut}}, t_0)$