

# Neutrino-Nucleon Form Factors from Lattice QCD

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New Directions in Neutrino Nucleus Scattering

# Outline

- ▶ Introduction
  - Motivation
  - Interaction Modes
  - $z$  expansion
- ▶ (High-Level) LQCD Results
  - Fermilab Lattice  $g_A$
  - Callat  $g_A$
  - Quick survey of  $g_A$ ,  $r_A^2$
- ▶ Future Prospects
  - QE
  - RES
  - NN
- ▶ Conclusions

# Introduction

# Challenges with Oscillation Experiments

Deep Underground Neutrino Experiment – Flagship \$1b experiment at LBNF

Measure osc. prob. over first two oscillation peaks

⇒  $E_\nu$  btw. [0.5, 10] GeV

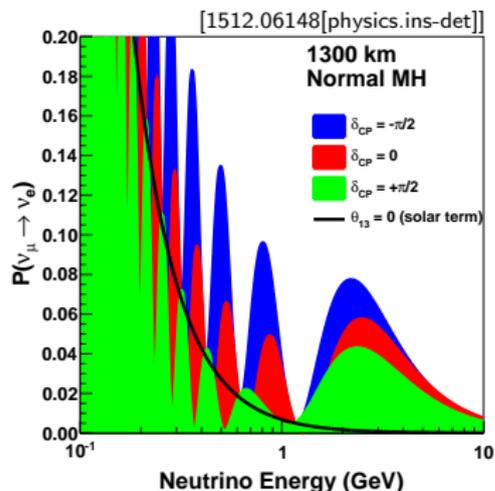
Osc. prob. is fn of  $L/E_\nu$

⇒ must classify events by  $E_\nu$

Neutrinos from secondary beam

⇒  $E_\nu$  not known event by event

⇒  $E_\nu$  inferred from distribution



Many requisite inputs difficult/impractical to measure from expt

⇒ Need ambitious theory support to supplement the experimental effort!

# Final State Interactions

Liquid Argon target

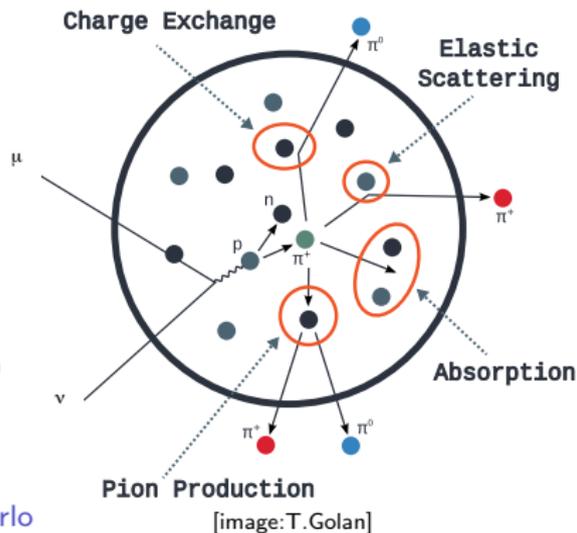
Nuclear xsec  $\implies$  nuclear modeling

Nuclear amplitudes constructed from  
one/few nucleon response

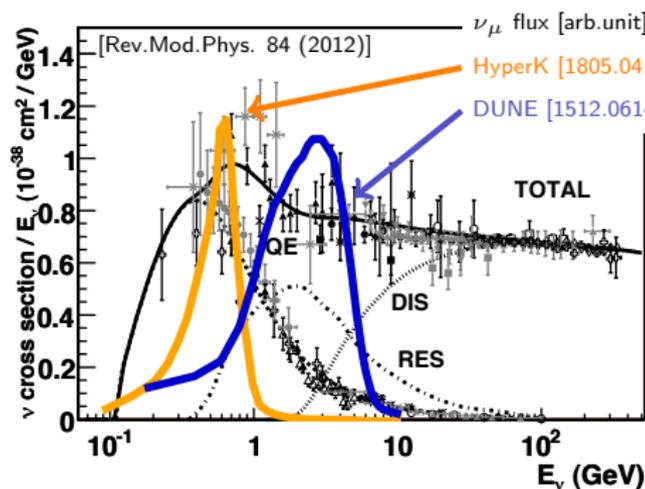
Event characteristics change in nuclear medium

In general:  $E_\nu \neq E_{\text{vis.}}$

$\implies$  Must infer  $E_\nu$  statistically from Monte Carlo



# Neutrino-Nucleon Interaction Modes

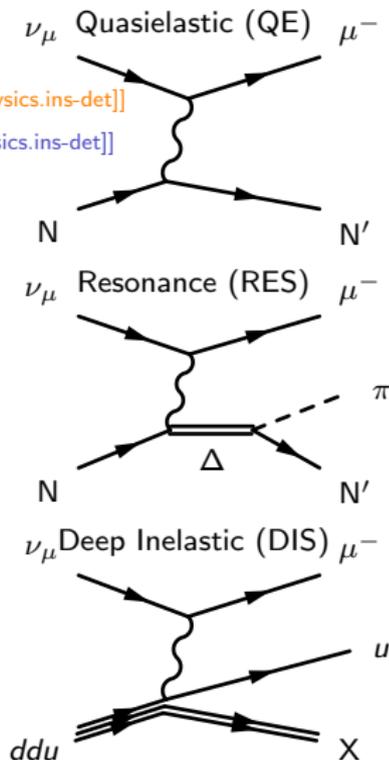


Event rate/ $E_\nu$  bin  $\sim \int_{\text{bin}} dE_\nu$  (Flux)  $\times$  (Xsec)

DUNE appx. 1:1:1 events for QE:RES:DIS

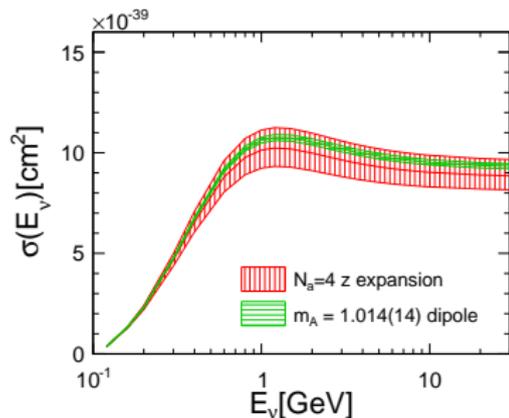
$\Rightarrow$  all interaction topologies are important

Focus on QE & RES, but LQCD comp. strategies exist for DIS (& SIS) too

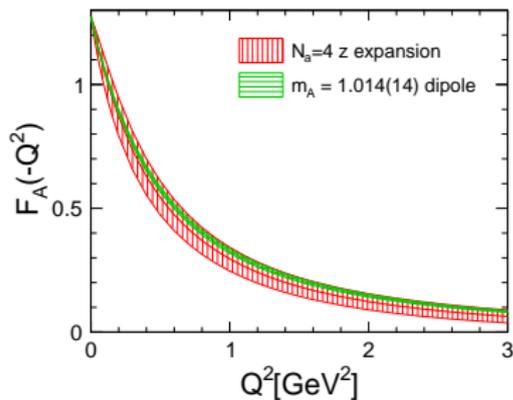


## z Expansion Fit to QE $F_A$

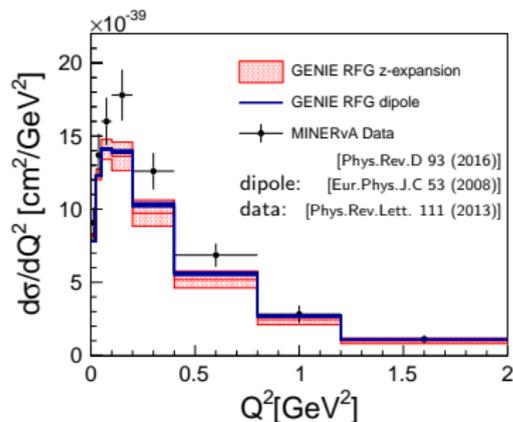
- ▶ Dipole has strict  $Q^2$  shape, inconsistent w/ QCD
- ▶ Dipole FF ansatz significantly underestimates FF uncertainty



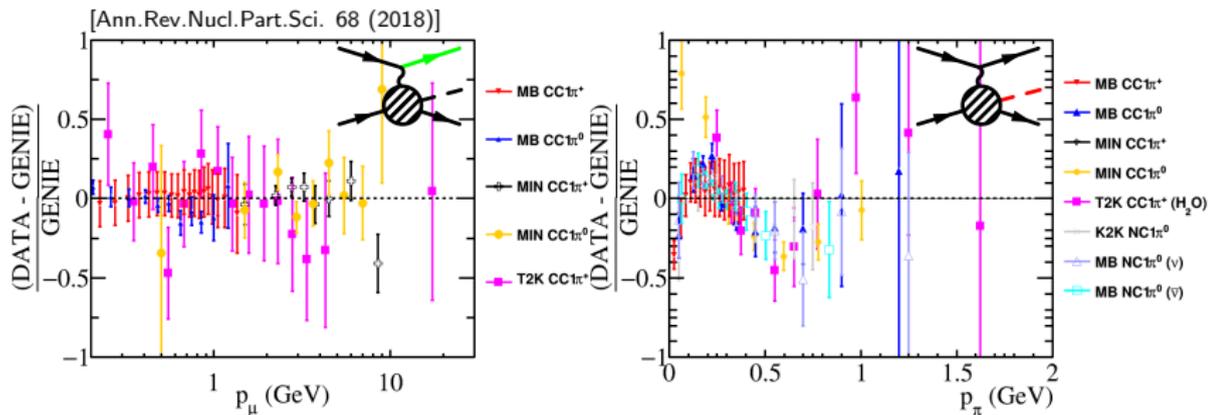
- ▶ Nucl. xsec uncertainty from FF same size as data-MC tensions
- ▶ Source of tensions unclear btw. nucleon/nuclear



- ▶ Model-independent parameterization  
Order of mag. increase in  $\delta\sigma$



# Leptonic vs Hadronic



Correlated differences between data & MC in leptonic, hadronic models

Balancing act to reconcile two variables

Insufficient model to describe interactions?

# Moving Forward

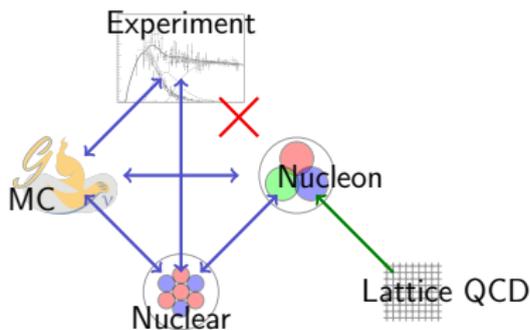
Room for improvement, but what is needed?

Ideal: Modern high stats  $\nu$ -D<sub>2</sub> scattering bubble chamber expt

Some community push, safety concerns

⇒ LQCD as a alternative/complement to expt

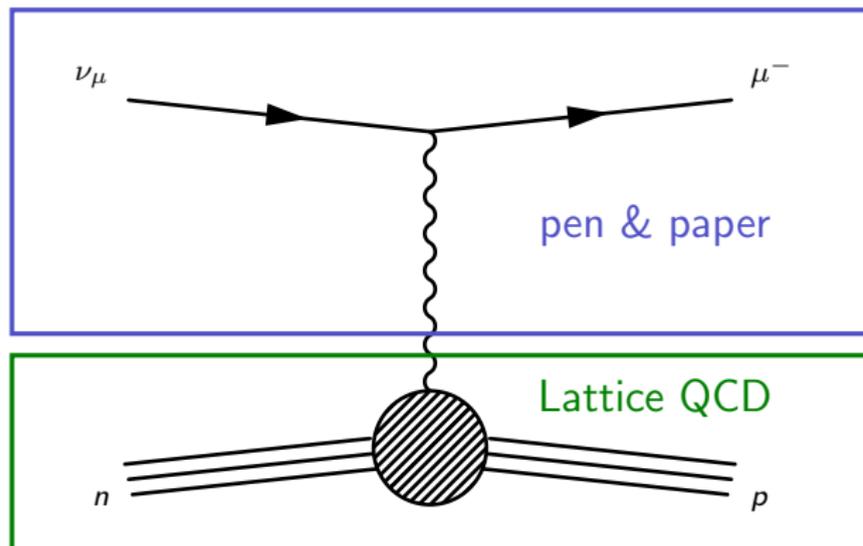
- ✓ No nuclear effects
- ✓ Realistic uncertainty estimates
- ✓ Systematically improvable
- ✓ Computers are (relatively) inexpensive



# How can Lattice Help?

Lattice is well suited to compute matrix elements:

$$\mathcal{M}_{\nu_{\mu} n \rightarrow \mu p} = \langle \mu | (V - A)_{\mu} | \nu \rangle \langle p | (V - A)_{\mu} | n \rangle$$



# Lattice QCD

# Lattice QCD: Formalism 1/2

Numerical eval of path integral

Quark, gluon DOFs —

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}U \exp(-S) \mathcal{O}_\psi [U]$$

Few inputs —

Computational:  $am_{(u,d),\text{bare}}$   
 $am_{s,\text{bare}}$   
 $\beta = 6/g_{\text{bare}}^2$

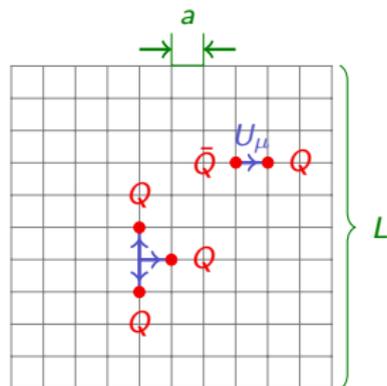
Scale setting: e.g.  $\frac{M_\pi}{M_\Omega}$ ,  $\frac{M_K}{M_\Omega}$ ,  $M_\Omega$   
one per computational input

Results — first principles predictions of QCD

$$M_{\text{hadron}}, \langle F | \mathcal{O} | I \rangle$$

Euclidean time  $\implies C(t) \sim e^{-M \cdot t}$

“Complete” error budget  $\implies$  extrapolation in  $a$ ,  $L$ ,  $(M_\pi)$



# Lattice QCD: Formalism 2/2

Correlation functions computed in euclidean time:

2-point function

$$\langle \mathcal{O}_1(t) \mathcal{O}_2(0) \rangle = \sum_n \langle 0 | \mathcal{O}_1 | n \rangle \langle n | \mathcal{O}_2 | 0 \rangle e^{-E_n t}$$

3-point function

$$\langle \mathcal{O}_1(t) \mathcal{O}_2(\tau) \mathcal{O}_3(0) \rangle = \sum_{mn} \langle 0 | \mathcal{O}_1 | n \rangle \langle n | \mathcal{O}_2 | m \rangle \langle m | \mathcal{O}_3 | 0 \rangle e^{-E_n(t-\tau) - E_m \tau}$$

Large  $t$ : excited states decay away, signal-to-noise degrades

Computations performed on “gauge ensembles” with fixed physics:

$$\{a, L, M_\pi, \dots\}$$

Extrapolate in ensemble parameters to arrive at physical point

- ▶  $a \rightarrow 0$  (continuum limit)
- ▶  $L \rightarrow \infty$  (infinite volume limit)
- ▶  $M_\pi \rightarrow M_\pi^{\text{phys}}$  (chiral limit)

## Word of caution:

Many collaborations will compare experiment to unextrapolated results...

These values will have uncontrolled systematics.

Make sure you know what you are looking at!

# Fermilab Lattice — $M_N$

Nucleon mass —  $C_{2pt}(t) \sim \sum_n z_n z_n^\dagger e^{-E_n t}$ ,  $z_n = \langle 0 | \mathcal{O} | n \rangle$

Demonstration of method:

Baryons w/ “Highly-Improved Staggered Quarks” (HISQ)

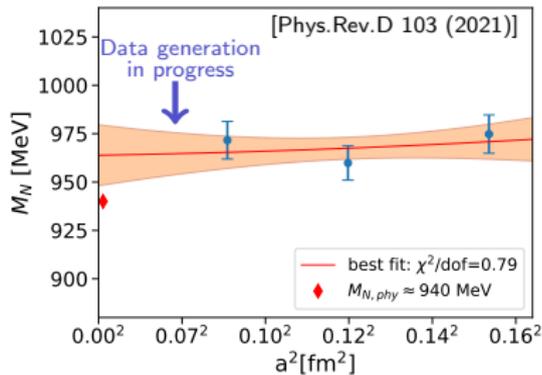
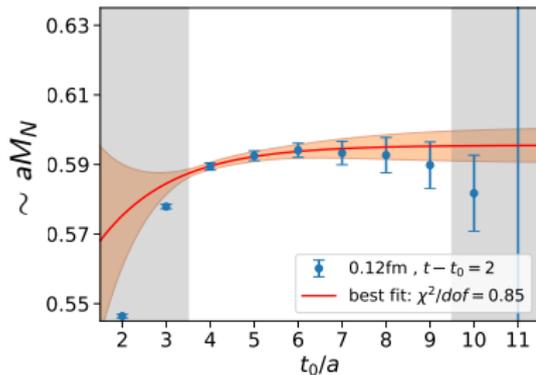
Additional SU(4) “taste” symmetry

Baryon octet  $\rightarrow$  baryon 572-plet

Complicated group thy, cheaper computation (1-comp spinors)

3 ensembles: all  $M_\pi^{\text{phys}}$ ; various  $a$ ,  $L$

Analysis credit: Yin Lin



# Fermilab Lattice — $g_A, g_V$

Axial charge —

$$C_{3pt}(t, \tau) \sim \sum_{mn} z_n z_m^\dagger \langle n | \mathcal{A}_\mu | m \rangle e^{-E_n(t-\tau) - E_m \tau}$$

$$\langle N | \mathcal{A}_\mu | N \rangle \sim g_A$$

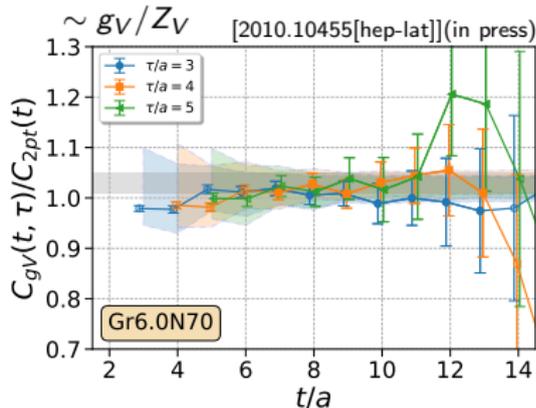
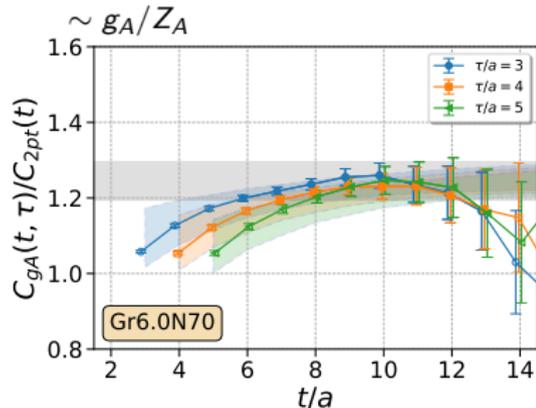
Single ensemble:  $M_\pi \approx 305$  MeV

World-first 3-point fn w/ HISQ

Interpreting results requires

$$SU(2)_{\text{flavor}} \times SU(4)_{\text{taste}} \subset SU(8) \text{ CG coefs/Wigner-Eckart}$$

Analysis credit: Yin Lin



Left:  $O(1\%)$  on  $g_A$  using Feynman-Hellman inspired technique:

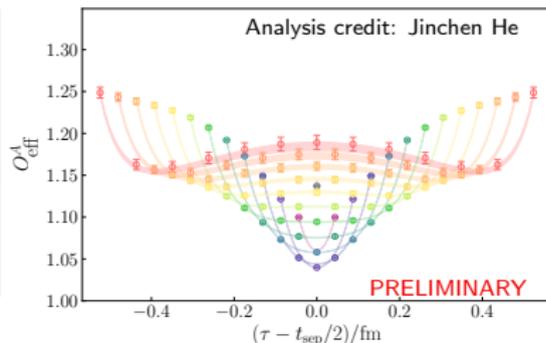
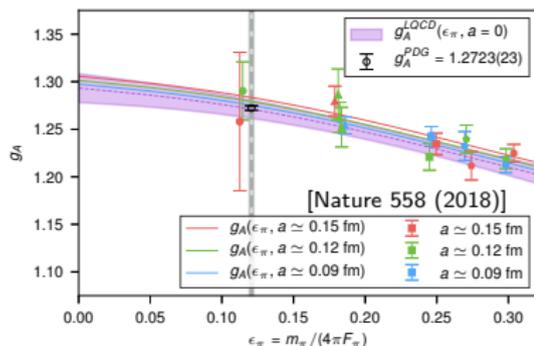
$$\underbrace{\frac{dE_\lambda}{d\lambda}}_{2\text{-pt}} = \underbrace{\langle \psi_\lambda | \frac{dH_\lambda}{d\lambda} | \psi_\lambda \rangle}_{3\text{-pt}} \quad \text{w/ source term} \quad \lambda \int d^4x \mathcal{A}_\mu(x)$$

Right: Work in progress —

One ensemble:  $a \approx 0.09$  fm,  $M_\pi \approx 310$  MeV

5-state fit, detailed analysis of excited state contamination

Compare w/ traditional three-point method



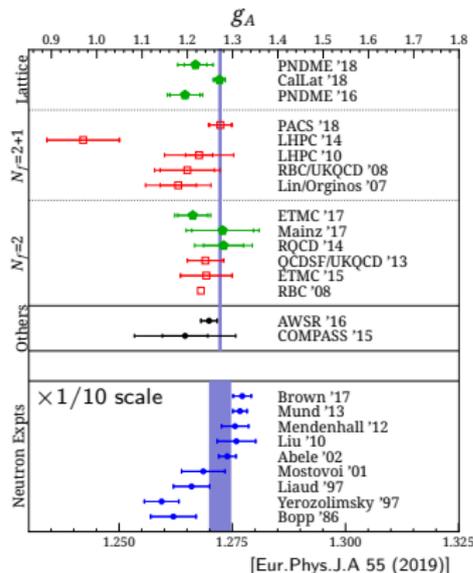
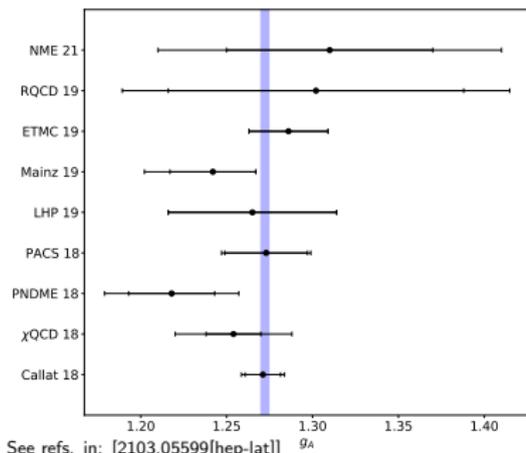
# Survey — $g_A$

State of field circa 2019 summarized in white paper,  
written for nonpractitioners: [Eur.Phys.J.A 55 (2019)]

Historically low  $g_A$  attributed to  $N$  excitations ( $N\pi$ , RES)  
led to apparent violations of PCAC relation

Now:

- Agreement w/ PDG seen consistently
- Details about excitations still unclear



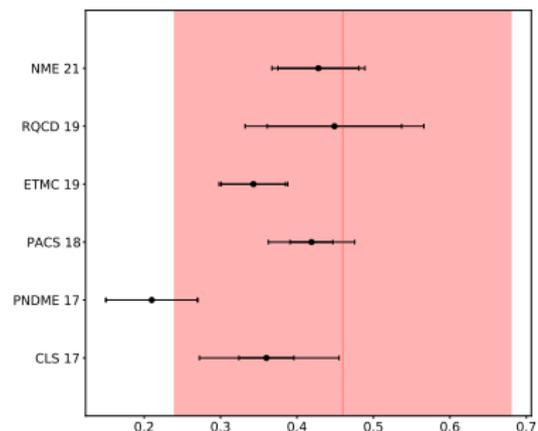
# Survey — $r_A^2$

Avoid  $M_A$ : only makes sense in dipole  $\implies r_A^2 = -(6/g_A)dF_A/dQ^2|_{Q^2=0}$   
 $\implies r_{A,\text{dipole}}^2 = (1 \text{ GeV}^2/m_A^2) \times 0.466\text{fm}^2$

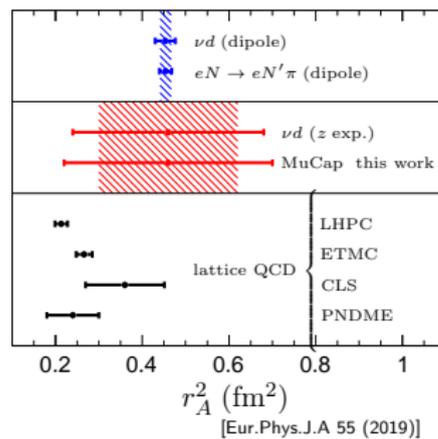
Most collaborations have adopted  $z$  expansion parameterization

Uncertainties will continue to decrease w/ time

Opinion: Still too early by few years — wait for resolution of excited state issues



See refs. in: [2103.05599[hep-lat]]  $r_A^2/\text{fm}^2$

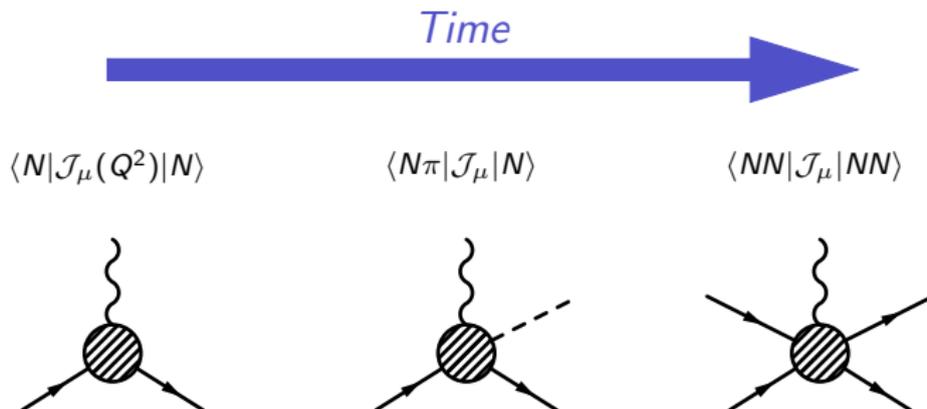


$r_A^2$  (fm<sup>2</sup>)  
[Eur.Phys.J.A 55 (2019)]

# Prospects

# Timeline

Very rough sketch of my interpretation of timeline for LQCD computations



# Forward-Looking: QE FFs

High-precision meas.  $\langle N | \mathcal{J}(\vec{q}, \nu) | N \rangle \sim F_A(Q^2), F_V(Q^2),$

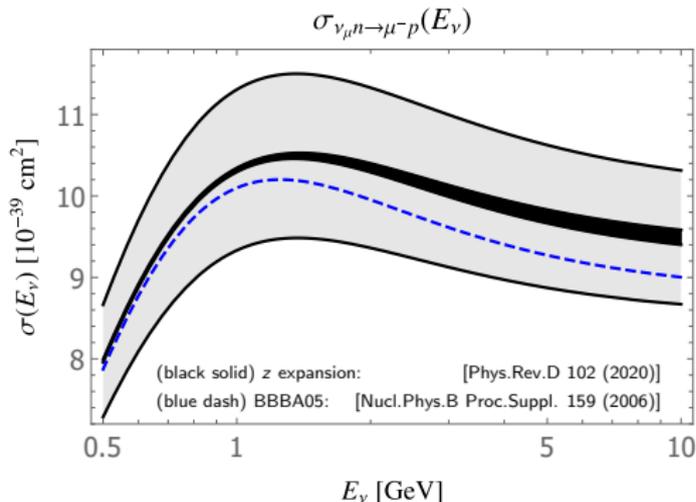
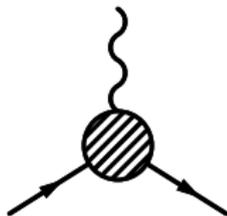
Realistically:  $Q^2 \sim 0 - 1 \text{ GeV}^2$

Reduce uncertainties of axial/pseudoscalar FFs

Fill in where difficult/impractical to get expt data

Resolve tensions btw vector FF parameterizations

Discrepancy of proton mag FF  $\implies$  uncertainty floor for axial FF



# Forward-Looking: Resonant Transitions

Constrain amplitudes/FFs that are **inaccessible to expt**

$$\langle N\pi | \mathcal{J}_\mu(\vec{q}, W) | N \rangle [\text{res.} + \text{nonres.}] \sim C_{5A}(Q^2)$$

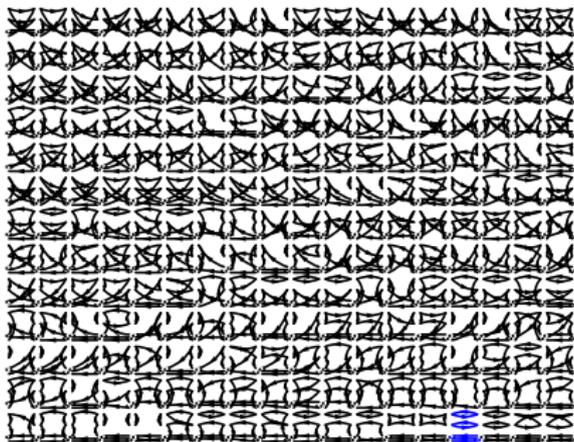
Theoretically & computationally challenging:

⇒ hard cutoff at  $N\pi\pi$  threshold (for now)

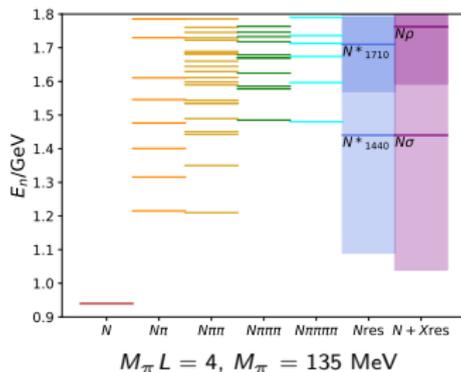
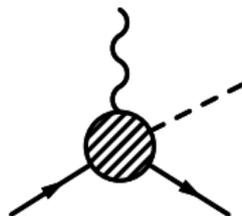
⇒ raise  $M_\pi$  to circumvent,  $\Delta \rightarrow$  stable

Dense spectrum of states

⇒ In practice requires  $N\pi$ ,  $(N\pi\pi)$  operators



[1912.04917[hep-lat]]  
[github.com/lehner/Wick]



# Forward-Looking: $NN$ $g_A$

$$\langle NN | \mathcal{J}_\mu | NN \rangle \sim F_V^{D_2}, F_A^{D_2}$$

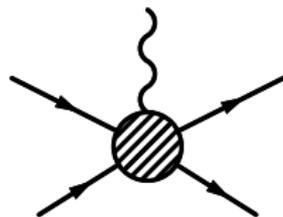
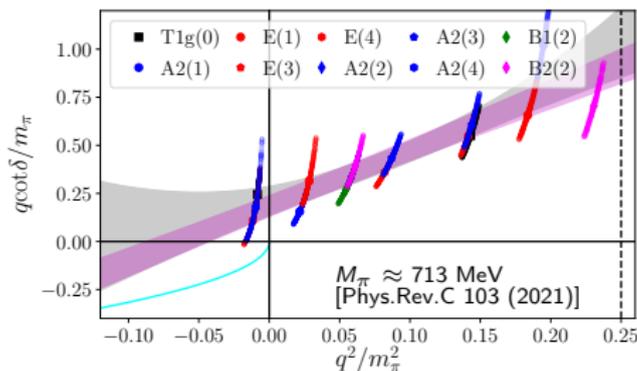
Direct comparisons to  $D_2$  scattering, nucl. model

$\Rightarrow$  Test nucl. corrections to QE assumption on  $D_2$

Signal to noise exponentially degrades  $\sim e^{-(2M_N - 3M_\pi)t}$

$\Rightarrow$  heavy  $M_\pi$  required (for now)

Nuclear models w/ large  $M_\pi$  could offer direct comparisons in near future



# Closing Remarks

Nuclear xsecs necessary for  $\nu$  oscillation experiments,  
but nucleon amplitude uncertainties still leave something to be desired —

- ▶ Nuclear corrections make extraction of nucleon amplitudes difficult
- ▶ Nucleon amplitudes from  $D_2$  have large uncertainties
- ▶ New  $D_2$  data not forthcoming (yet)

In absence of modern  $D_2$  expt, LQCD can fill missing pieces to puzzle —

- ▶ Improved stats on QE form factors
- ▶ Resolve tensions in vector form factor parameterizations
- ▶ Compute amplitudes that are difficult/impractical to measure in expt
- ▶ Match directly to (small) nuclear target data/models

Lots of work to be done, but path forward is clear!