

# GLSMs, localisation, and the stringy Kähler moduli space

[with M. Romo, E. Scheidegger: [arXiv:2003.00182\[hep-th\]](https://arxiv.org/abs/2003.00182)] [with D. Erkinger: [arXiv:2008.03089\[hep-th\]](https://arxiv.org/abs/2008.03089)]

Johanna Knapp

School of Mathematics and Statistics, University of Melbourne

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# Outline

Overview

GLSMs and localisation

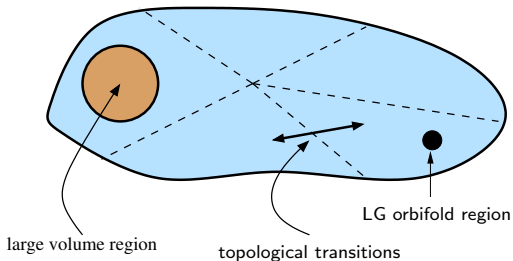
Hemisphere

Sphere

Conclusions

## Stringy Kähler moduli space $\mathcal{M}_K$

- **This talk:** Calabi-Yau threefolds in a type II setting (+ branes)
- $\mathcal{M}_K$  is divided into **chambers**



- **Quantum corrections** due to worldsheet instantons.
- One can cross the chamber boundaries.
- Categories associated to chambers (conjecturally) equivalent.
- Study  $\mathcal{M}_K$  in limiting regions that are **not large volume**.

$\mathcal{M}_K$  via GLSMs

- The 2D (2, 2) **gauged linear sigma model (GLSM)** provides a common UV description of the CFTs in  $\mathcal{M}_K$ . [Witten 93]
- **Phases:** different low-energy (IR) configurations depending on the values of the FI-theta parameters:

$$t = \zeta - i\theta \quad \leftrightarrow \quad t \in \mathcal{M}_K \quad \leftrightarrow \quad \text{chambers}$$

- Use the GLSM to map out  $\mathcal{M}_K$ .
- **SUSY localisation** computes (quantum) exact expressions.
  - Can be evaluated in **any phase**.
  - $t$ -dependence  $\leftrightarrow$  worldsheet instanton corrections.
  - Compute quantum corrections directly in the GLSM.

## Beyond geometry

- **Geometric phases are well-understood** and we know how to do computations.
  - Toric geometry, topological strings, mirror symmetry, enumerative invariants, etc.
- The structures we know in geometry should also be visible in **non-geometric settings**. Why?
  - The **worldsheet CFT** of the phases does not care whether there is a geometry.
  - Whatever is computed in non-geometric phases using **localisation** has the same UV origin as the geometric results.
- If we evaluate localisation results in phases of the GLSM, **how do we interpret the result in the IR theory?**

## Universal structures in phases of GLSMs

- **Claim:** When evaluating the GLSM partition functions in different types of phases the (exact) result has the same structure in every phase.
- To see this, we need **structures that are available beyond geometric settings**, e.g.
  - Worldsheet: CFT, topological gravity,  $tt^*$ -geometry
  - FJRW theory [Witten 92][Fan-Jarvis-Ruan 07]
  - Givental's mirror construction [Givental 96-03]
  - Categories
- **Ingredients** that appear in every phase:
  - State space and pairing
  - I/J-function
  - Gamma class [Hosono 00][Iritani 07][Katzarkov-Kontsevich-Pantev 08]
  - D-brane data

## Results

- Hemisphere partition function:

[Romo-Scheidegger-JK 20]

$$Z_{D^2}^{\text{phases}}(\mathcal{B}) = \langle \text{ch}(\mathcal{B}), \Gamma \cdot I \rangle$$

- Works in geometry and Landau-Ginzburg orbifold phases.
- Sphere partition function:

[Erkinger-JK 20]

$$Z_{S^2}^{\text{phases}} = \langle \bar{I}, I \rangle, \quad \langle \bar{I} | = (-1)^{\text{Gr}} \frac{\Gamma}{\Gamma_*} \bar{I}$$

- Works in hybrid phases that are LG orbifolds fibered over some base manifold.

## GLSM data

- $G$  ... a compact Lie group (gauge group)
- $V$  ... space of chiral fields  $\phi_i \in V$
- $\rho_V : G \rightarrow GL(V)$  ... faithful complex representation
  - **CY condition:**  $G \rightarrow SL(V)$
- $R : U(1)_V \rightarrow GL(V)$  ... vector R-symmetry
  - $R_i$  ... R-charges
- $T \subset G$  ... maximal torus
  - Lie algebras:  $\mathfrak{g} = Lie(G)$ ,  $\mathfrak{t} = Lie(T)$
  - $Q_i^a \in \mathfrak{t}_{\mathbb{C}}^*$  ... gauge charges of chiral fields



## GLSM Data (ctd.)

- $t^a \in \mathfrak{g}_{\mathbb{C}}^*$  ... FI-theta parameters
  - $t^a = \zeta^a - i\theta^a$       $\zeta^a$  : real,     $\theta^a$  :  $2\pi$ -periodic
  - $t^a \leftrightarrow$  Kähler moduli of the CY
- $\sigma_a \in \mathfrak{g}_{\mathbb{C}}$  ... scalar components of the vector multiplet
- $W \in \text{Sym} V^*$  ... superpotential
  - $G$ -invariant
  - $R$ -charge 2
  - non-zero for **compact** CYs

## D-branes in GLSMs

- D-branes (B-type) in the GLSM are  $G$ -invariant **matrix factorisations** of the GLSM potential with  $R$ -charge 1

[Herbst-Hori-Page 08][Honda-Okuda,Hori-Romo 13]

- Data:

- **Sym  $V^*$ -module** (Chan-Paton space):  $M = M^0 \oplus M^1$
- **Matrix Factorisation**:  $Q \in \text{End}^1(M)$  with

$$Q^2 = W \cdot \text{id}_M$$

- **$G$ -action**:  $\rho : G \rightarrow GL(M)$  with

$$\rho(g)^{-1} Q(g\phi) \rho(g) = Q(\phi) \quad g \in G$$

- **$R$ -action**:  $r_* : u(1)_V \rightarrow gl(M)$  with

$$\lambda^{r_*} Q(\lambda^R \phi) \lambda^{-r_*} = \lambda Q(\phi) \quad \lambda \in U(1)_V$$

## Sphere partition function $Z_{S^2}$

- Sphere partition function

[Benini-Cremonesi 12][Doroud-Gomis-LeFloch-Lee 12]

$$Z_{S^2} = C \sum_{m \in \mathbb{Z}^{\text{rk}G}} \int_{\gamma} d^{\text{rk}G} \sigma \prod_{\alpha > 0} \left( \frac{\alpha(m)^2}{4} + \alpha(\sigma)^2 \right) \cdot \prod_{j=1}^{\dim V} \frac{\Gamma \left( iQ_j(\sigma) - \frac{Q_j(m)}{2} + \frac{R_j}{2} \right)}{\Gamma \left( 1 - iQ_j(\sigma) + \frac{Q_j(m)}{2} - \frac{R_j}{2} \right)} e^{-4\pi i \zeta(\sigma) - i\theta(m)}$$

- $\alpha > 0$  positive roots
- $\gamma \dots$  integration contour (s.t. integral is convergent)
- $Z_{S^2}$  computes the **exact Kähler potential** on  $\mathcal{M}_K$ .

[Jockers-Kumar-Lapan-Morrison-Romo 12][Gomis-Lee 12]

[Gerchkovitz-Gomis-Komargodski 14][Gomis-Hsin-Komargodski-Schwimmer-Seiberg-Theisen 15]

## Hemisphere partition function

- **Hemisphere partition function:** [Sugishita-Terashima][Honda-Okuda][Hori-Romo 13]

$$Z_{D^2}(\mathcal{B}) = C \int_{\gamma} d^{\text{rk}_G} \sigma \prod_{\alpha > 0} \alpha(\sigma) \sinh(\pi \alpha(\sigma)) \\ \cdot \prod_{j=1}^{\dim V} \Gamma \left( iQ_j(\sigma) + \frac{R_j}{2} \right) e^{it(\sigma)} f_{\mathcal{B}}(\sigma)$$

- **Brane factor**

$$f_{\mathcal{B}}(\sigma) = \text{tr}_M \left( e^{i\pi \mathbf{r} \cdot \sigma} e^{2\pi \rho(\sigma)} \right)$$

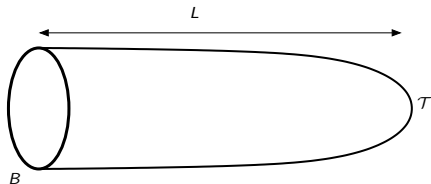
- $Z_{D^2}$  computes the **exact D-brane central charge**.

## D-brane central charge – worldsheet perspective

- The central charge connects A/B-branes with  $(c, c)/(a, c)$ -operators:

[Ooguri-Oz-Yin 96][Hori-Iqbal-Vafa 00]

$$Z(B) = \langle B | \mathbf{0} \rangle$$



- Here we consider B-branes  $\mathcal{B}$  and  $(a, c)$ -operators.
  - Note:** “B-branes in the A-model”

## $Z_{D^2}$ in LG orbifold phases

- We want to collect evidence that

[Romo-Scheidegger-JK 20]

$$Z_{D^2}^{\text{phases}}(\mathcal{B}) = \langle \text{ch}(\mathcal{B}), \Gamma \cdot I \rangle.$$

- $\langle \cdot, \cdot \rangle$  ... pairing (state space)
- $\Gamma$  ... Gamma class
- $I$  ...  $I$ -function
- $\text{ch}(\mathcal{B})$  ... Chern character
- Focus on **Landau-Ginzburg orbifold** phases.
- **Plan:**
  - Define the objects on the right-hand side for LG orbifolds.
  - Show that the results match with  $Z_{D^2}$  and FJRW theory.

## LG data

- **Landau-Ginzburg orbifold:**  $(W, G, \bar{\rho}_m, U(1)_{L/R})$   
[Vafa 89][Intriligator-Vafa 90]
  - $W(x_i)$  ... superpotential
  - $G$  ... orbifold group (assume  $\mathbb{Z}_{d_1} \times \mathbb{Z}_{d_2} \times \dots$ )
  - $\bar{\rho}_m$  ... matter representation
  - $U(1)_{L/R}$  ... left/right R-symmetry
- **State space**
  - $(c, c), (a, c), \dots$ - chiral rings  $\leftrightarrow \gamma$ -twisted sectors ( $\gamma \in G$ ) with basis elements  $\phi_\gamma$
  - Restriction to one-dimensional (“**narrow**”) sectors ( $\delta \in G$ )
  - (Topological) **pairing** on the  $(a, c)$ -ring:

$$\langle \phi_\delta, \phi_{\delta'} \rangle = \frac{1}{|G|} \delta_{\delta, \delta'^{-1}}$$

- **Note:** Spectral flow required to map  $(c, c)$  to  $(a, c)$ .

## LG branes

- **LG orbifold B-brane:**  $\bar{\mathcal{B}} = (\bar{M}, \bar{Q}, \bar{\rho}_\gamma, \bar{r}_*)$

[Kapustin-Li 02][Brunner-Herbst-Lerche-Scheuner 03][Walcher 04]

- $\bar{M}$ ... Chan-Paton space
- $\bar{Q}$ ... matrix factorisation of  $W$
- $\bar{\rho}_\gamma, \bar{r}_*$ ... representation of  $G$  and  $u(1)_V$  on the boundary
- **Chern character:**

$$\text{ch}(\bar{\mathcal{B}})_\gamma = \frac{1}{n_\gamma!} \text{Res}_{W_\gamma} (\Phi_\gamma \cdot \text{Tr}_{\bar{M}} [\bar{\rho}_\gamma (\partial \bar{Q}_\gamma)^{\wedge n_\gamma}])$$

- $\Phi_\gamma$  ... state(s) in the  $\gamma$ -twisted sector  $\mathcal{H}_\gamma$  (to be precise, this is the  $(c, c)$ -ring)
- $n_\gamma$  ...  $\dim \text{Fix}_\gamma$
- $W_\gamma = W|_{\text{Fix}_\gamma}$



## Gamma class

- Taking into account **deformations** away from the LG point, one can define a  $h \times (N + h)$  matrix  $q$ .
  - $N$  ... number of chirals  $x_i$ .
  - $h$  ... number of marginal deformations (narrow sectors).
- **Gamma class:**

$$\Gamma\phi_\delta = \Gamma_\delta\phi_\delta \quad \Gamma_\delta = \prod_{j=1}^N \Gamma \left( 1 - \left\langle \sum_{a=1}^h k_a q_{a,h+j} \right\rangle \right).$$

- $k_a \in \mathbb{Z}_{\geq 0}^h$ : each gets associated to a sector  $\delta$ .
- $\langle x \rangle = x - [x]$

# I-function

- I-function

$$I_{LG}(u) = \sum_{\delta \in G} I_{\delta}(u) \phi_{\delta}$$

$$I_{\delta}(u) = - \sum_{\substack{k_1, \dots, k_h \geq 0 \\ k'_i = \delta_i \pmod{d_i}}} \frac{u^k}{\prod_{a=1}^h \Gamma(k_a + 1)} \\ \times \prod_{j=1}^N \frac{(-1)^{\langle -\sum_{a=1}^h k_a q_{a,h+j} + q_j \rangle} \Gamma(\langle \sum_{a=1}^h k_a q_{a,h+j} - q_j \rangle)}{\Gamma(1 + \sum_{a=1}^h k_a q_{a,h+j} - q_j)}$$

- $u$  ... local coordinates/deformation parameters
- $q_j$  ... left R-charges of the  $x_j$

## $J$ -function

- One obtains the  $J$ -function from the  $I$ -function through a transformation to **flat coordinates**.
  - These correspond to the deformation parameters of the marginal deformations in the worldsheet CFT.
- Select components of the  $I_{\delta_a}$  ( $a = 1, \dots, h$ ) of the  $I$ -function associated to the subspace of (narrow) marginal deformations  $\leftrightarrow (a, c)$  states with  $(q, \bar{q}) = (-1, 1)$  and the unique component  $I_0$  with  $(q, \bar{q}) = (0, 0)$ .
- The flat coordinates and the  $J$ -function are

$$t_a(u) = \frac{I_{\delta_a}(u)}{I_0(u)} \quad J(t) = \frac{I_{LG}(u(t))}{I_0(u(t))}$$

- This coincides with the **mirror map**.

## Consistency checks

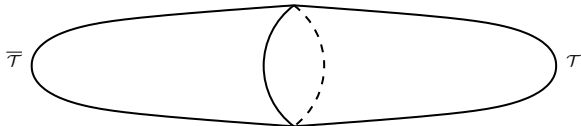
- The LG quantities **consistent with FJRW theory**.  
[Chiodo-Ruan 08][Chiodo-Iritani-Ruan 12]
  - This gives an independent check for our results.
  - Our results **generalise results from FJRW theory**: more moduli, more general  $G$ .
- The **hemisphere partition function** of the GLSM reproduces the results from the central charge formula when evaluated at LG points.
  - GLSM gauge group broken to orbifold group  $G$
  - Matrix  $q$  is related to the gauge charges of the GLSM fields
  - **GLSM branes**: matrix factorisations of the GLSM superpotential reduce to LG matrix factorisations  
[Herbst-Hori-Page 08][Clarke-Guffin 10]
- Tested in **examples** with up to four Kähler moduli, including computation of **FJRW invariants**.

# Kähler potential – worldsheet perspective

- The Kähler potential of  $\mathcal{M}_K$  is

[Cecotti-Vafa '91]

$$e^{-K(t, \bar{t})} = \langle \bar{0} | 0 \rangle$$



- $\langle \bar{0} |$  and  $|0\rangle$  are related by CPT conjugation.

# Hybrids

- Consider a GLSM phase that is a Landau-Ginzburg orbifold with orbifold group  $G$  fibered over some base manifold  $B$ .
- This includes:
  - Calabi-Yau complete intersections in toric ambient spaces
  - Landau-Ginzburg orbifolds
- The **sphere partition function** evaluates to [Erkinger-JK 20]

$$Z_{S^2}^{\text{phases}}(t) = \sum_{\delta \in G} \int_B (-1)^{\text{Gr}} \frac{\Gamma_{\delta}(H)}{\Gamma_{\delta}^*(H)} |I_{\delta}(t, H)|^2 = \langle \bar{I}, I \rangle$$

- $\delta \in G$  ... narrow sectors
- Basis  $H \in H^2(B)$
- Gr ... grading operator on the state space

## $Z_{S^2}$ in phases

- **Landau-Ginzburg phases:**
  - $B$  is a point (no  $H$ )
  - Definitions of  $I/J$ , Gr, pairing, and the Gamma class coincide with definitions above
- **Geometric phases:**
  - $B$  is a complete intersection CY  $X$ ,  $G$  is trivial
  - The pairing is  $\langle \alpha, \beta \rangle = \int_X \alpha^\vee \wedge \beta$ ,  $\alpha, \beta \in H^{even}(X)$ .
  - $I/J$ , Gr, and Gamma class coincide with results in the literature
- Results suggest a definition [Iritani 07][Halverson-Jockers-Lapan-Morrison 13]

$$\langle \bar{I} | = (-1)^{\text{Gr}} \frac{\Gamma}{\Gamma^*} I(\bar{t}).$$

## Consistency checks

- Tested for GLSMs associated to **14 one-parameter complete intersections** in toric ambient phases.
  - The small radius phases are Landau-Ginzburg, hybrids, and pseudo-hybrids. [Aspinwall-Plesser 09]
  - The structure is observed in all phases, even pseudo-hybrids, where we have a sum of terms.
  - $I$ -functions and Gamma class in hybrid phases **match with FJRW theory**. [Clader 13]
- **Two-parameter** CY hypersurface with geometric, LG, and hybrid phase.
  - New conjectural results for the  $I$ -function and the Gamma class in multi-parameter hybrid models.



## Summary

- We conjectured universal expressions for the hemisphere and sphere partition functions for phases of (abelian) GLSMs.
- Evidence that this works for geometric, Landau-Ginzburg and hybrid phases.
- Results match with mathematical results from FJRW theory and mirror symmetry, where available.
- Tested for lots of examples.

## Open Questions

- Gamma class from the worldsheet perspective?
- Non-abelian GLSMs.
- Broad sectors.
- Hybrids, in particular with branes and enumerative invariants.
- More localisation results.
- Mathematical proofs.