

# Benefits of a Near Detector for JUNO



**REBEKAH PESTES**

# Neutrino Oscillation Parameters

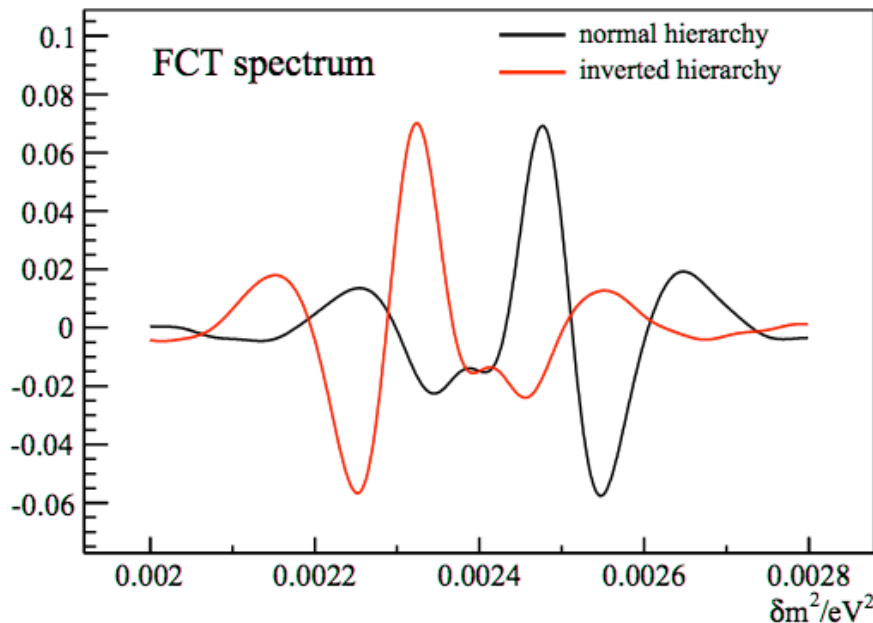
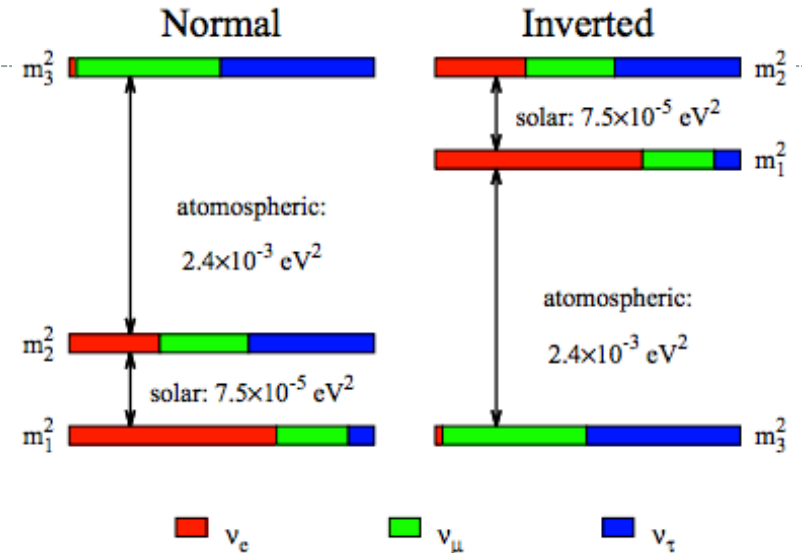
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- 6 independent oscillation parameters in three-flavor oscillation
  - 3 mixing angles:  $\theta_{12}, \theta_{23}, \theta_{13}$
  - 2 mass-differences:  $\Delta m^2_{21}, \Delta m^2_{31}$
  - 1 CP-violating phase:  $\delta$
- Only two remain largely unknown:
  - Sign of  $\Delta m^2_{31}$
  - $\delta$

# Neutrino Mass Hierarchy

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- $\Delta m_{21}^2 \sim 7 \times 10^{-5} \text{ eV}^2$ ,  
 $|\Delta m_{31}^2| \sim 2 \times 10^{-3} \text{ eV}^2$
- Normal (NH):  $\Delta m_{31}^2 > 0$
- Inverted (IH):  $\Delta m_{31}^2 < 0$

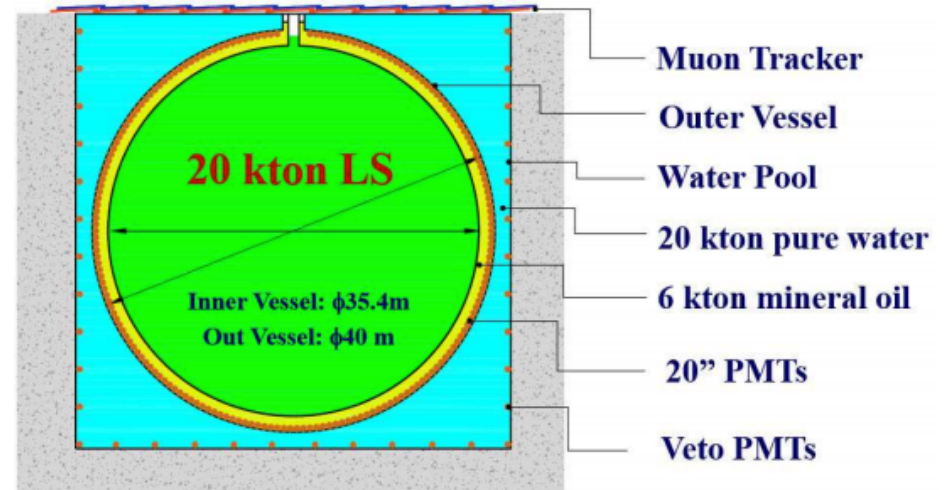


- Determination Method:
  - $L \sim 50 \text{ km}$
  - Observe Beats in Oscillation Frequencies
  - Energy resolution needed:  $\left| \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \right| \sim 0.03$

# JUNO

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- Jiangmen Underground Neutrino Observatory
- Detector needs  $3\%/\sqrt{E}$  energy resolution to determine the neutrino mass hierarchy
- 20 kton of Liquid Scintillator
- Measure  $\bar{\nu}_e$  disappearance from 2 reactors (each at 53 km):

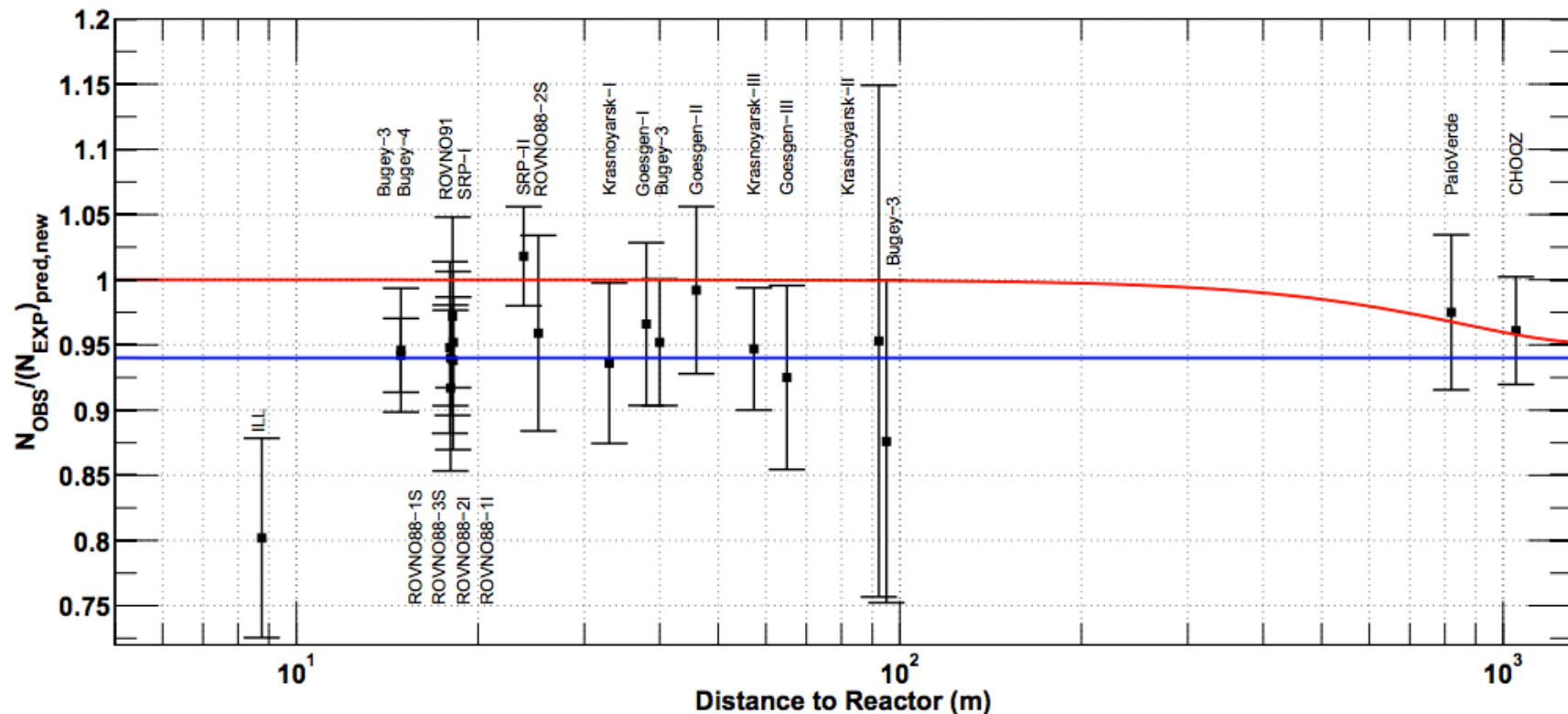


$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2(2\theta_{12})\cos^4\theta_{13}\sin^2\left(\Delta m_{21}^2 \frac{L}{4E}\right) - \sin^2(2\theta_{13})\left[\cos^2\theta_{12}\sin^2\left(\Delta m_{31}^2 \frac{L}{4E}\right) + \sin^2\theta_{12}\sin^2\left(\Delta m_{32}^2 \frac{L}{4E}\right)\right]$$

# Reactor Antineutrino Anomaly

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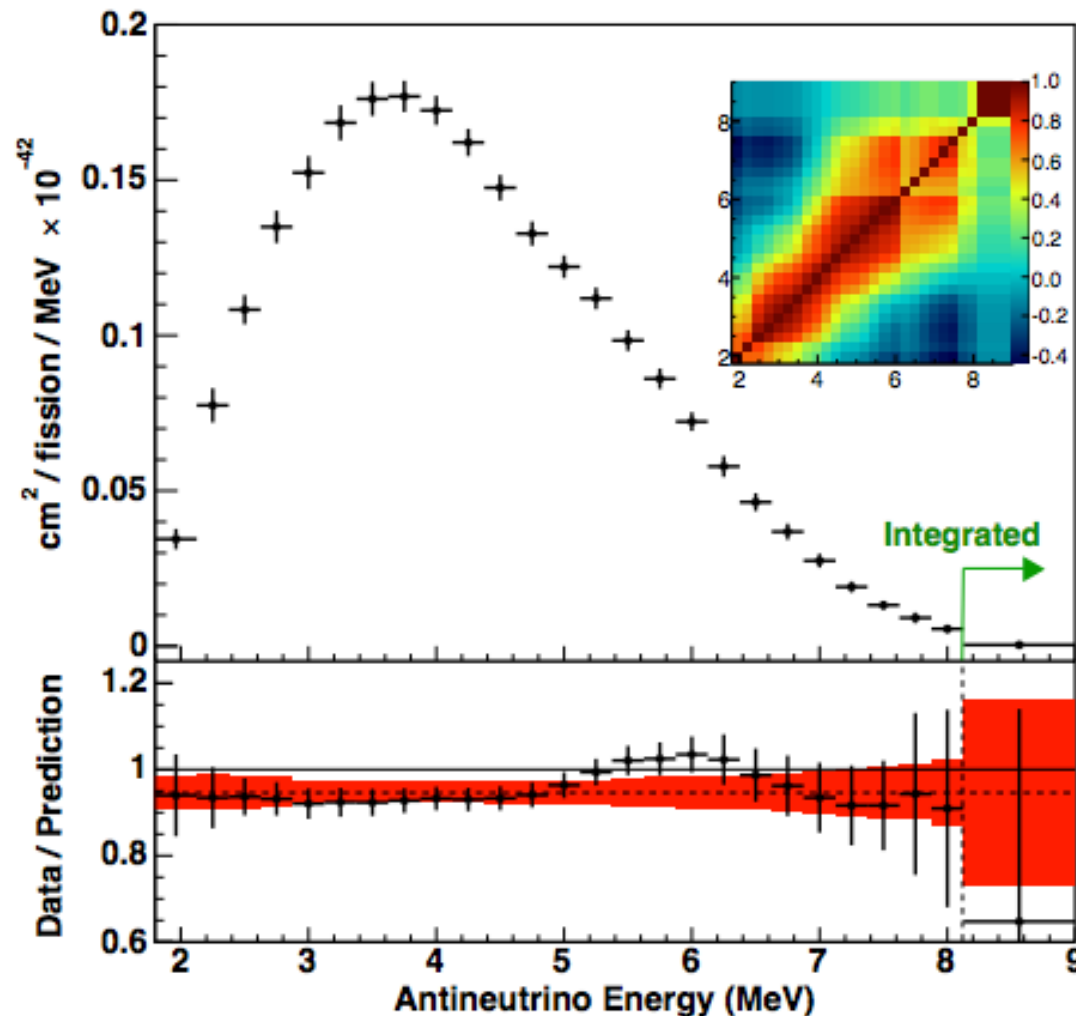
- Disagreement between measured antineutrino energy spectrum from nuclear reactors and updated theoretical calculations



Mention, et al. (arXiv:1101.2755)

# Reactor Antineutrino Anomaly Cont'd

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- Possible Implications:
  - shows ignorance of what is going on in the reactor
    - (i.e. reactor  $\nu_e$  spectrum)
  - something is happening to the neutrinos on short length scales after creation
    - (e.g. sterile  $\nu$ )

## JUNO Cont'd

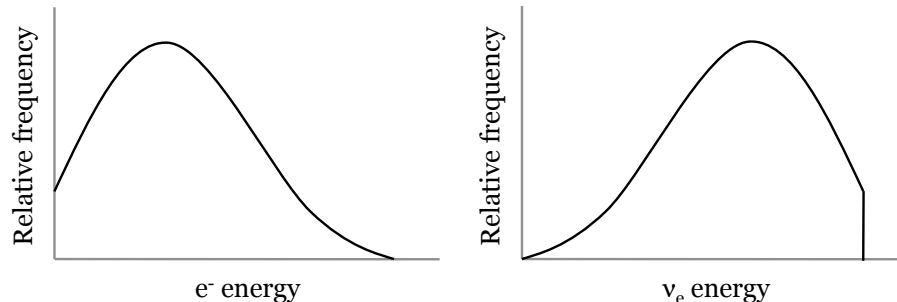
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- Had been proposing to use Daya Bay data to obtain a reference neutrino energy spectrum
  - Daya Bay used virtually identical detectors for near and far detectors
  - Spectrum measured at energy resolution of  $8\%/\sqrt{E}$

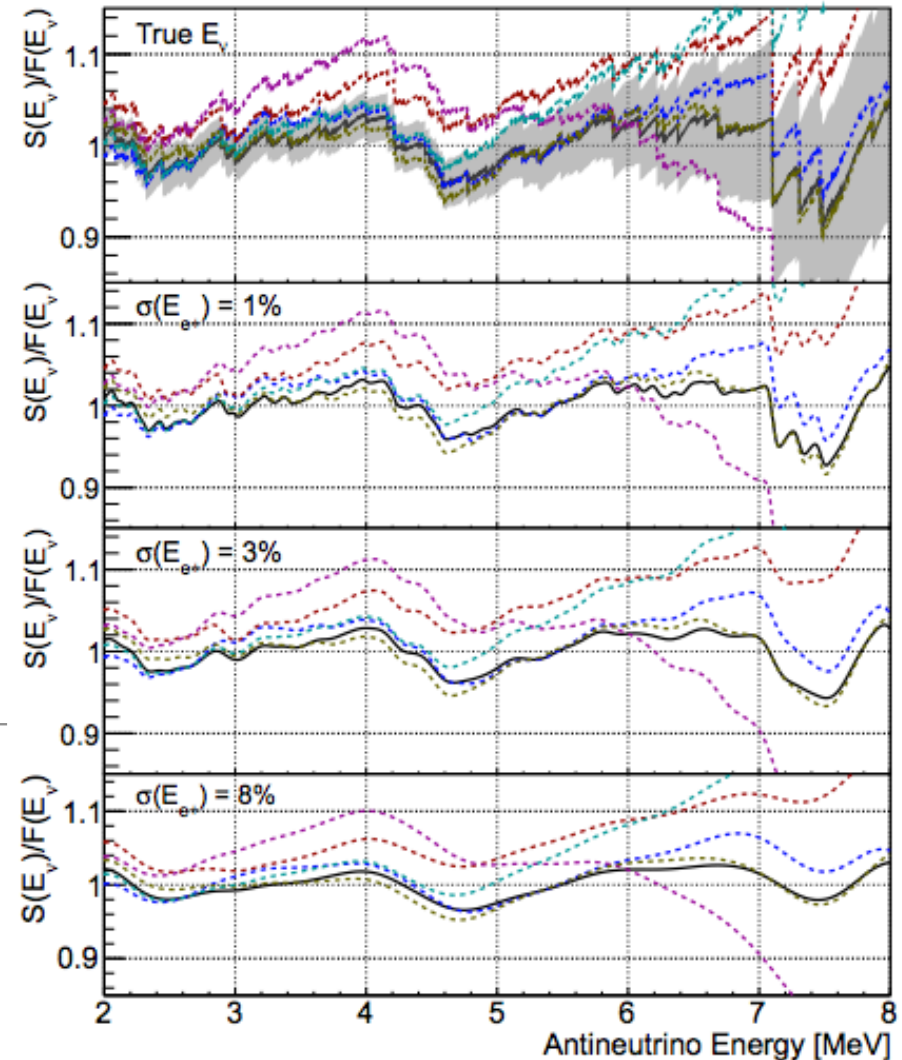
# Possible Problem

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- Fine structure in antineutrino energy spectrum coming from nuclear reactor
  - For single  $\beta$ -decay:



- Obscured at  $8\%/\sqrt{E}$
- Possibly visible at  $3\%/\sqrt{E}$





# JUNO Simulations

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- Used GLoBES
- JUNO's specs for far detector
  - 20 kt liquid scintillator
  - 53 km away from source
  - $3\%/\sqrt{E}$  energy resolution
- Added near detector
  - 5 ton liquid scintillator
  - 0.5 km away from source
  - Varied energy resolution

# JUNO Simulations Cont'd

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- Rate data divided into 100 energy bins

$\phi_i^I \equiv$  detection rate in  $i$ th bin

- Nuisance parameters added for each bin

$$\tilde{\phi}_{fit_i}^I = (1 + \xi_i) \left( \phi_{fit_i}^I \right)^o, \text{ where } \phi_{fit_i}^I = \sum_j M_{ij} \tilde{\phi}_{fit_j}^I$$

- Accounts for uncertainty in energy spectrum
- Accounted for non-linearity in detector energy response

$$\frac{E_{rec}}{E} = 1 + \sum_{k=0}^n \alpha_k E^k$$

- Allowed differences between detectors

# JUNO Simulations Cont'd

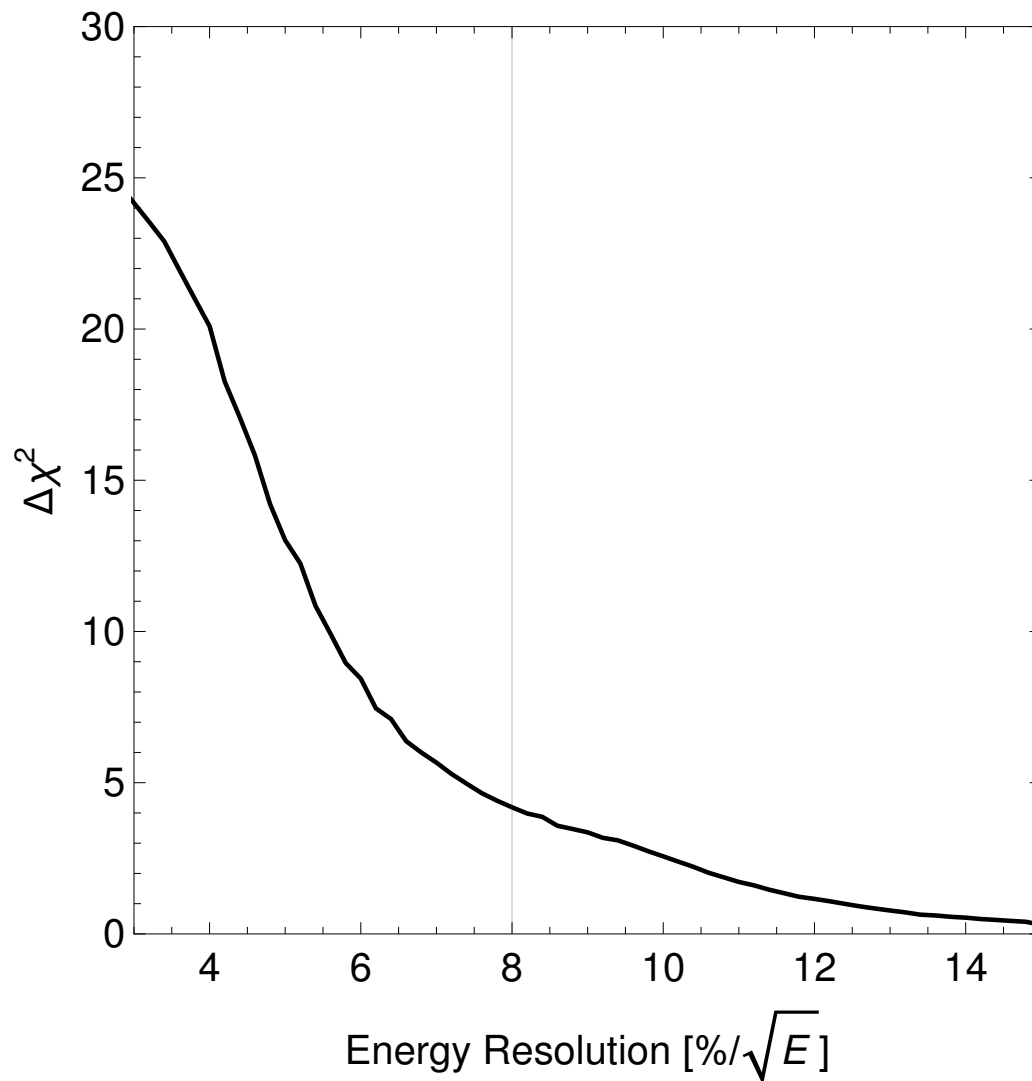
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- Huber-Mueller model used for source antineutrino energy spectrum
- Assumed normal hierarchy to be true for generating data
- Calculated  $\chi^2$  values for fits using the normal and inverted hierarchies

$$\chi^2 = \sum_{i,I} \frac{(\phi_{true_i}^I - \phi_{fit_i}^I)^2}{\phi_{true_i}^I} + \sum_j \left( \frac{s_j}{\sigma_j} \right)^2$$

# Results

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# Consistency Check #1

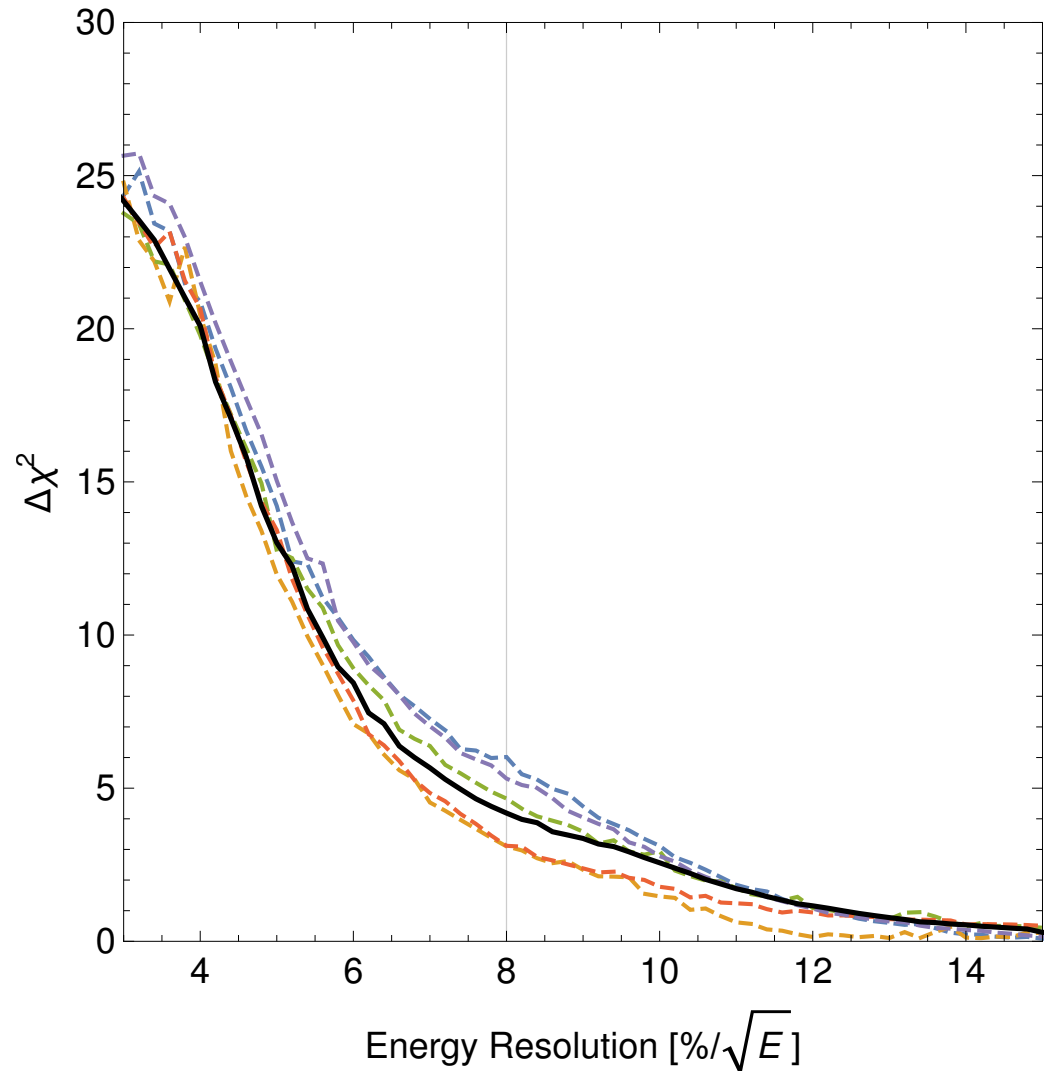
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- Used alternative spectra to generate data
  - Add up antineutrino spectra produced from a randomly chosen (with weights) sample of beta decay branches
  - Renormalize spectrum to match Huber-Mueller model at  $8\%/\sqrt{E}$  energy resolution
- Fit data to Huber-Mueller model
  - Meddled with internal structures of GLoBES

# Consistency Check #1

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- Generated data with alternative spectra
- Fit using Huber-Mueller model



## Consistency Check #2

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- No Near Detector
- Used Daya Bay's Covariance Matrix, instead

$$\chi_A^2 = \sum_i \frac{(\phi_{true_i} - \phi_{fit_i})^2}{\phi_{true_i}} + \sum_j \left( \frac{s_j}{\sigma_j} \right)^2 + \sum_{n,m} \xi_n (V^{-1})_{nm} \xi_m$$

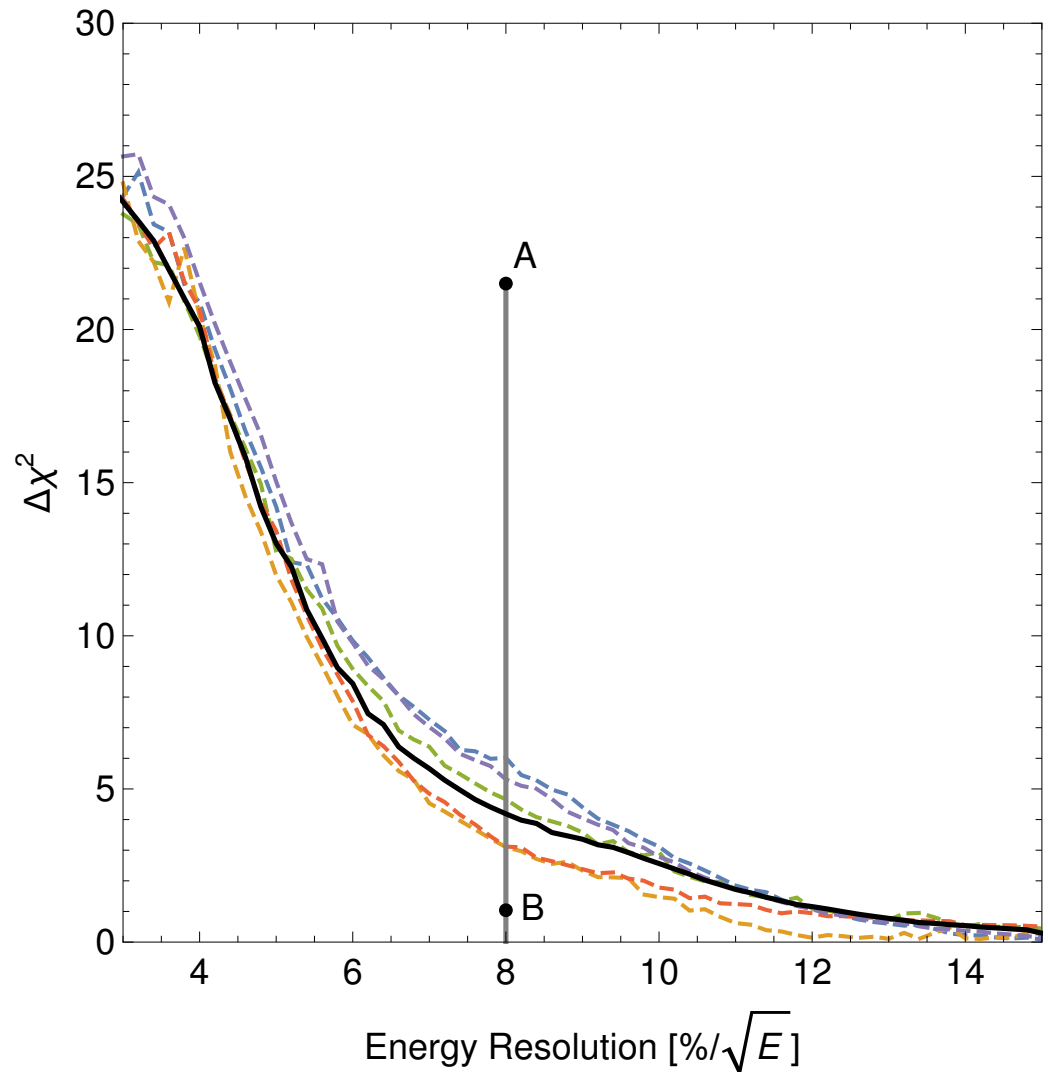
$$\chi_B^2 = \sum_i \frac{(\phi_{true_i} - \phi_{fit_i})^2}{\phi_{true_i}} + \sum_j \left( \frac{s_j}{\sigma_j} \right)^2 + \sum_{n,m} \alpha_n (V^{-1})_{nm} \alpha_m$$

$$\text{where } \alpha_n = \frac{\sum_{i \in N_n} \phi_{fit_i}^o \xi_i}{\sum_{i \in N_n} \phi_{fit_i}^o}$$

# Consistency Check #2

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- No Near Detector
- Used Daya Bay's Covariance Matrix, instead
- A: One nuisance parameter for each of Daya Bay's bins
- B: One nuisance parameter for each of the far detector's bins





## Consistency Check #3

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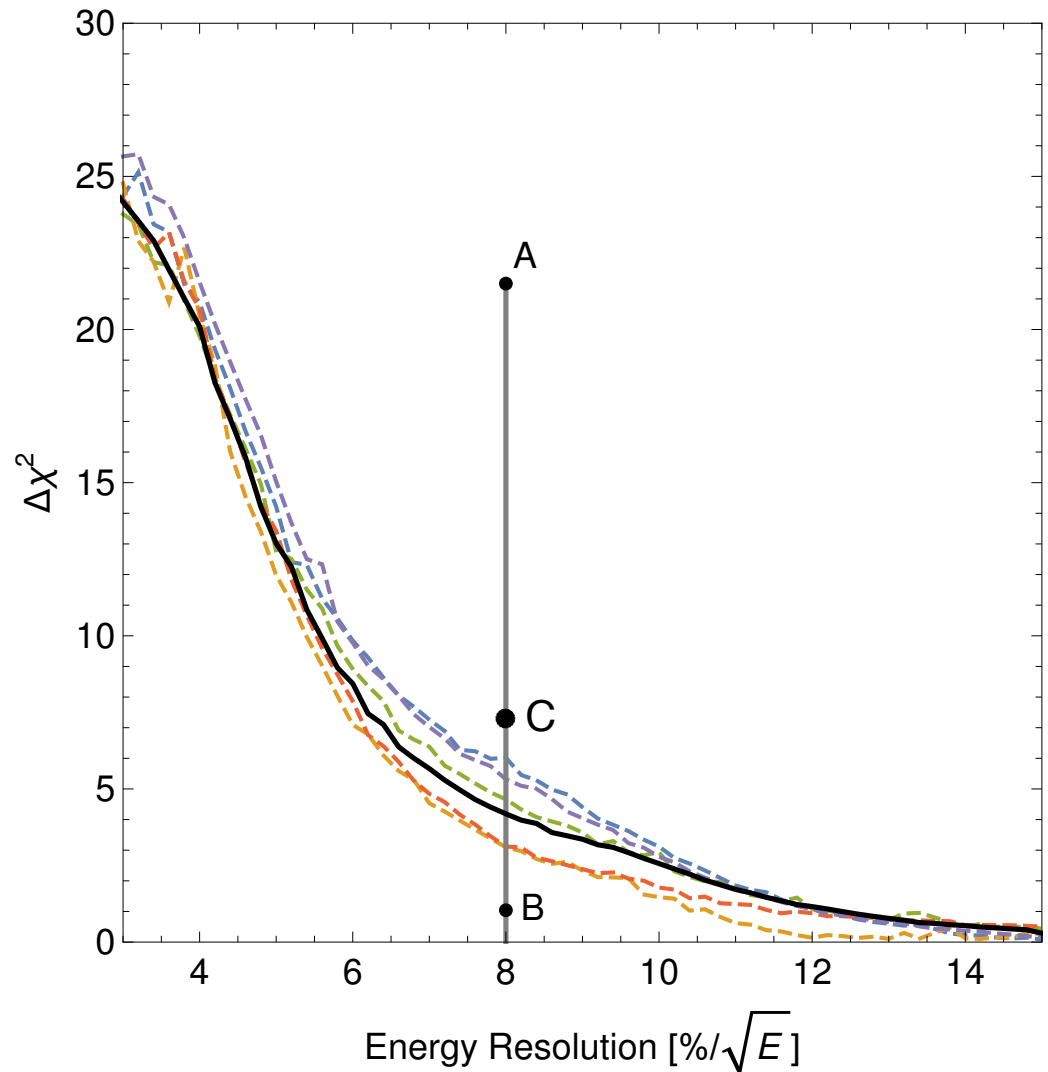
- No Near Detector
- Used the covariance matrix for an idealized detector, instead

$$\chi_C^2 = \sum_i \frac{(\phi_{true_i} - \phi_{fit_i})^2}{\phi_{true_i}} + \sum_j \left( \frac{s_j}{\sigma_j} \right)^2 + \sum_{i,k} \xi_i (V^{-1})_{ik} \xi_k$$

# Consistency Check #3

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- No Near Detector
- Used Idealized Covariance Matrix, instead
- C: One nuisance parameter per Daya Bay energy bin



# Summary

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- Can't trust theoretical reactor antineutrino spectrum
- Antineutrino spectrum fine structure prevents extrapolation to lower energy resolutions
  - JUNO can't use Daya Bay's data
- So, JUNO needs a near detector in order to resolve the mass hierarchy
  - Energy resolution  $\leq 3\%/\sqrt{E}$