

# Systematics in Neutrino Oscillation Experiments

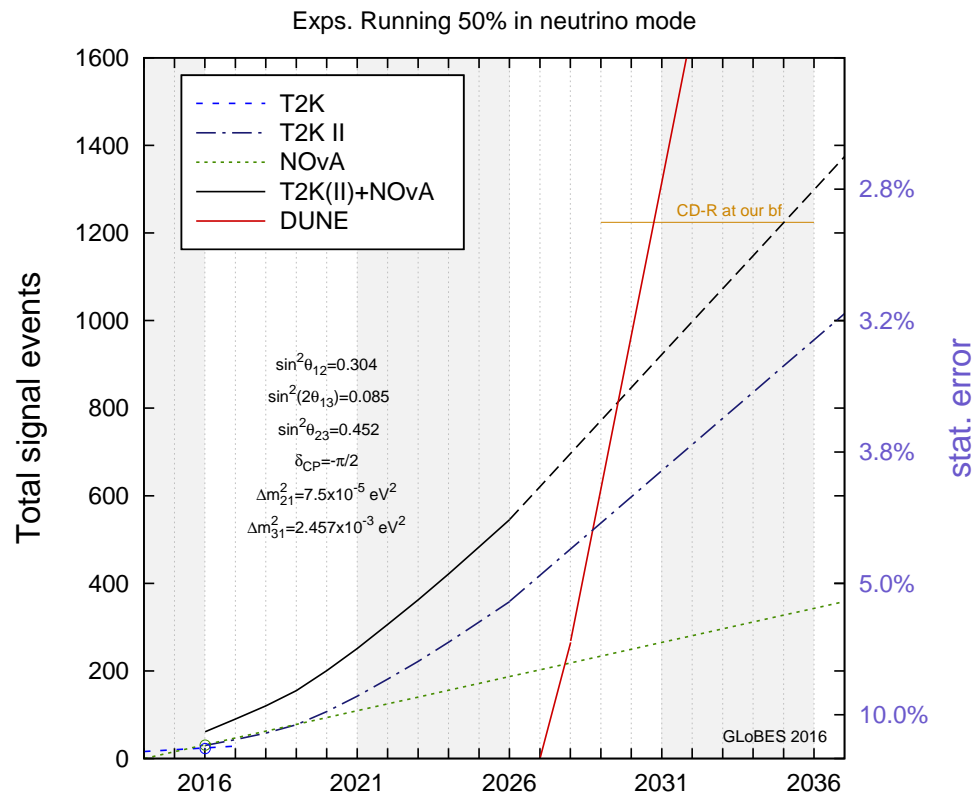
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# Statistical errors



Clearly, we are on the (slow) road towards 3% measurements of the event rates

Translating this into a 3% measurements of the oscillation probability is very difficult

Note, T2HK would reach 1000  $\nu_e$  signal events very quickly.

# The Idea

In order to measure CP violation we need to reconstruct one out of these

$$P(\nu_\mu \rightarrow \nu_e) \text{ or } P(\nu_e \rightarrow \nu_\mu)$$

and one out of these

$$P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \text{ or } P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu)$$

and we'd like to do that at the percent level accuracy

# The Reality

We do not measure probabilities, but event rates!

$$R_{\beta}^{\alpha}(E_{\text{vis}}) = N \int dE \Phi_{\alpha}(E) \sigma_{\beta}(E, E_{\text{vis}}) \epsilon_{\beta}(E) P(\nu_{\alpha} \rightarrow \nu_{\beta}, E)$$

In order to reconstruct  $P$ , we have to know

- $N$  – overall normalization (fiducial mass)
- $\Phi_{\alpha}$  – flux of  $\nu_{\alpha}$
- $\sigma_{\beta}$  – x-section for  $\nu_{\beta}$
- $\epsilon_{\beta}$  – detection efficiency for  $\nu_{\beta}$

Note:  $\sigma_{\beta}\epsilon_{\beta}$  always appears in that combination, hence we can define an effective cross section  $\tilde{\sigma}_{\beta} := \sigma_{\beta}\epsilon_{\beta}$

# The Problem

Even if we ignore all energy dependencies of efficiencies, x-sections *etc.*, we generally can not expect to know any  $\phi$  or any  $\tilde{\sigma}$ . Also, we won't know any kind of ratio

$$\frac{\Phi_{\alpha}}{\Phi_{\bar{\alpha}}} \quad \text{or} \quad \frac{\Phi_{\alpha}}{\Phi_{\beta}}$$

nor

$$\frac{\tilde{\sigma}_{\alpha}}{\tilde{\sigma}_{\bar{\alpha}}} \quad \text{or} \quad \frac{\tilde{\sigma}_{\alpha}}{\tilde{\sigma}_{\beta}}$$

Note: Even if we may be able to know  $\sigma_e/\sigma_{\mu}$  from theory, we won't know the corresponding ratio of efficiencies  $\epsilon_e/\epsilon_{\mu}$

# The Solution

Measure the un-oscillated event rate at a near location and everything is fine, since all uncertainties will cancel, (provided the detectors are identical and have the same acceptance)

$$\frac{R_{\alpha}^{\alpha}(\text{far}) L^2}{R_{\alpha}^{\alpha}(\text{near})} = \frac{N_{\text{far}} \Phi_{\alpha} \tilde{\sigma}_{\alpha} P(\nu_{\alpha} \rightarrow \nu_{\alpha})}{N_{\text{near}} \Phi_{\alpha} \tilde{\sigma}_{\alpha} 1}$$

$$\frac{R_{\alpha}^{\alpha}(\text{far}) L^2}{R_{\alpha}^{\alpha}(\text{near})} = \frac{N_{\text{far}}}{N_{\text{near}}} P(\nu_{\alpha} \rightarrow \nu_{\alpha})$$

And the error on  $\frac{N_{\text{far}}}{N_{\text{near}}}$  will cancel in the  $\nu$  to  $\bar{\nu}$  comparison.

# But ...

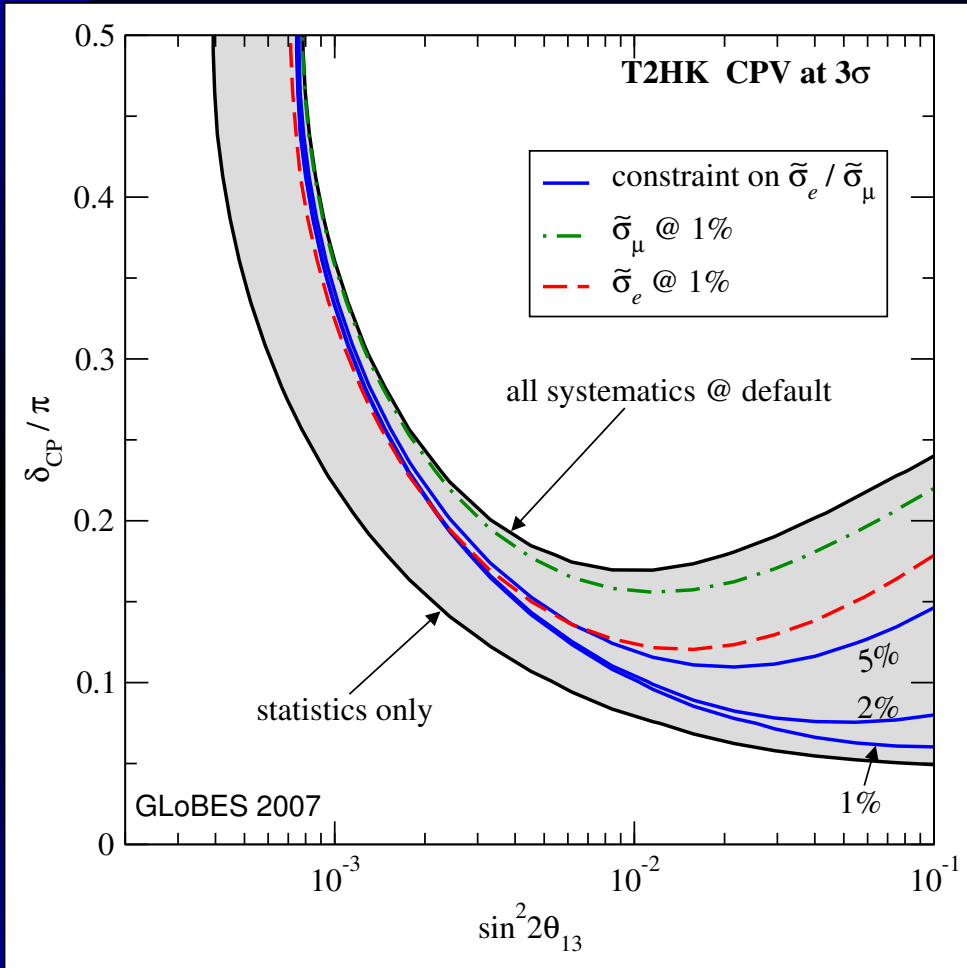
This all works only for disappearance measurements!

$$\frac{R_{\beta}^{\alpha}(\text{far}) L^2}{R_{\beta}^{\alpha}(\text{near})} = \frac{N_{\text{far}} \Phi_{\alpha} \tilde{\sigma}_{\beta} P(\nu_{\alpha} \rightarrow \nu_{\beta})}{N_{\text{near}} \Phi_{\alpha} \tilde{\sigma}_{\alpha} 1}$$

$$\frac{R_{\beta}^{\alpha}(\text{far}) L^2}{R_{\beta}^{\alpha}(\text{near})} = \frac{N_{\text{far}} \tilde{\sigma}_{\beta} P(\nu_{\alpha} \rightarrow \nu_{\beta})}{N_{\text{near}} \tilde{\sigma}_{\alpha} 1}$$

Since  $\tilde{\sigma}$  will be different for  $\nu$  and  $\bar{\nu}$ , this is a serious problem. And we can not measure  $\tilde{\sigma}_{\beta}$  in a beam of  $\nu_{\alpha}$ .

# $\nu_e/\nu_\mu$ total x-sections



Appearance experiments using a (nearly) flavor pure beam can **not** rely on a near detector to predict the signal at the far site!

Large  $\theta_{13}$  most difficult region.

PH, Mezzetto, Schwetz, 2007

Differences between  $\nu_e$  and  $\nu_\mu$  are significant below 1 GeV, see K. McFarland's talk



# Neutrino cross sections

Our detectors are made of nuclei and compared to a free nucleon, the following differences arise

- Initial state momentum distribution
- Nuclear excitations
- Reaction products have to leave the nucleus
- Higher order interactions appear

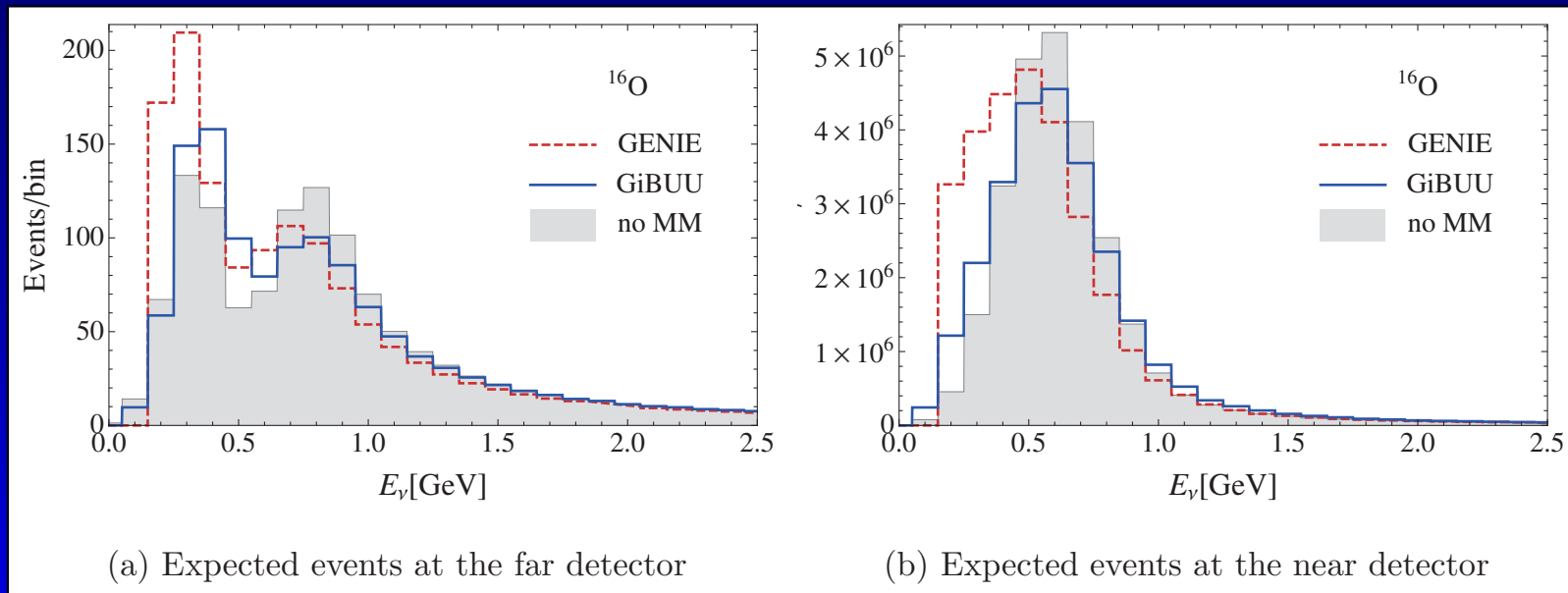
As a function of  $Q^2$  these effects are flavor blind, but we do NOT measure  $Q^2$ .

These effects are NOT the same for neutrinos and antineutrinos.

# Quasi-elastic scattering

QE events allow for a simple neutrino energy reconstruction based on the lepton momentum.

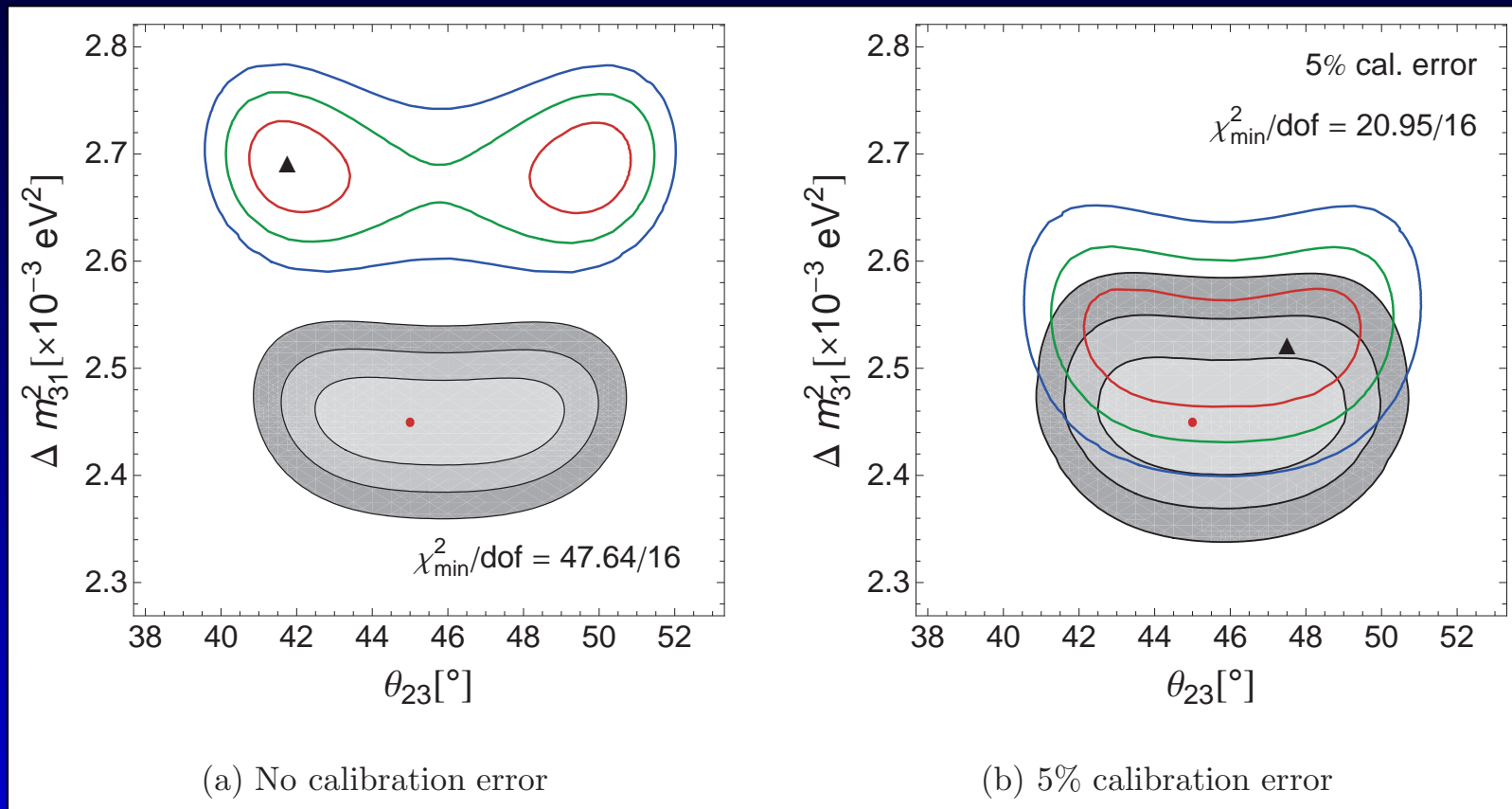
Nuclear effects will make some non-QE events appear to be like QE events  $\Rightarrow$  the neutrino energy will not be correctly reconstructed.



Coloma *et al.* 2013

# Impact on oscillation

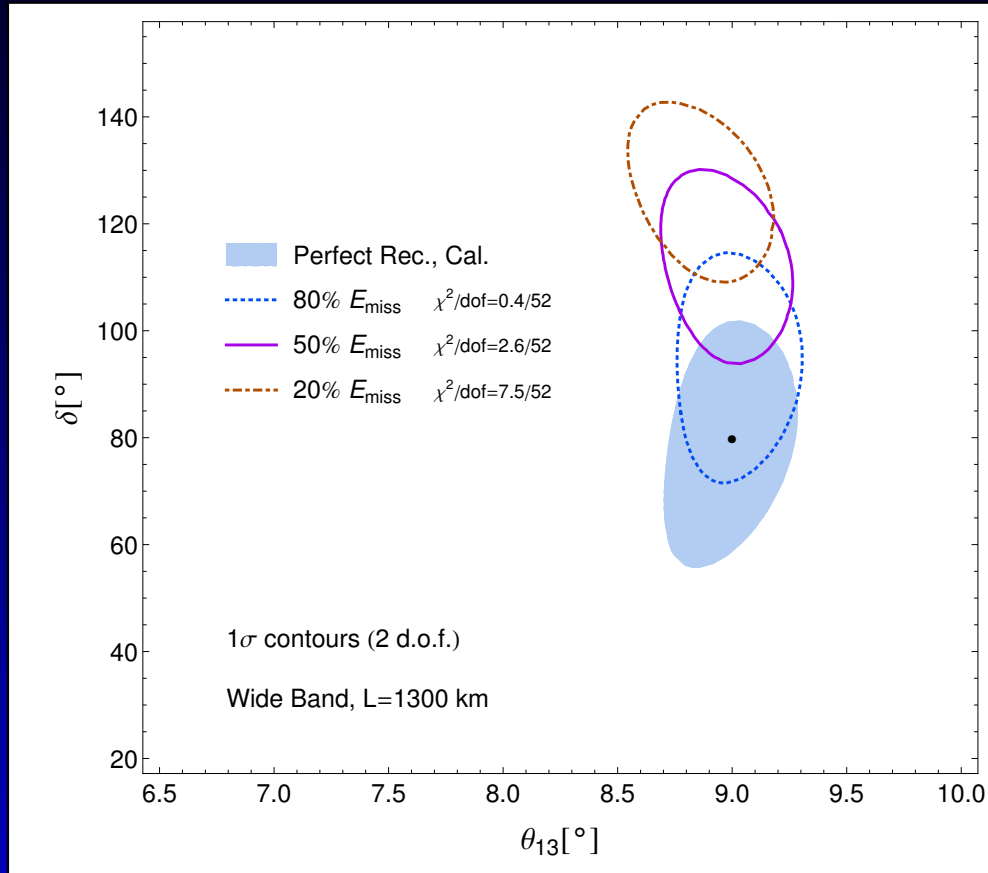
$\nu_\mu \rightarrow \nu_\mu$  in a T2K-like setup with near detector.



Coloma *et al.* 2013

If the energy scale is permitted to shift, tension and bias are reduced, but effects very hard to spot from  $\chi^2$

# Missing energy



In elastic scattering a certain number of neutrons is made

Neutrons will be largely invisible even in a liquid argon TPC

$\Rightarrow$  missing energy

Ankowski *et al.*, 2015

We can correct for the missing energy **IF** we know the mean neutron number and energy made in the event...

# Towards precise cross sections

This will require better neutrino sources, since a cross section measurement is about as precise as the accuracy at which the beam flux is known.

- Percent beam flux normalization
- Very high statistics needed to map phase space
- Neutrinos and antineutrinos
- $\nu_\mu$  and  $\nu_e$

A (the only?) source which can deliver all that is a muon storage ring, aka nuSTORM.