

New physics in IceCube

Iván Martínez Soler

ivanj.m@csic.es

Based on works done with Pilar Coloma, Pedro A.N. Machado and Ian M. Shoemaker,
arXiv:1707.08573

and

M. C. Gonzalez-Garcia, Michele Maltoni, Ivan Martinez-Soler and Ningqiang Song,
Astropart.Phys. 84 (2016) 15-22 arXiv:1605.08055

Summer Institute for Neutrino Theory

July 23rd, 2017



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 - Transition magnetic moment
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Motivation: ν mass

- In SM neutrinos are massless.
 - ▶ No dirac mass term for neutrinos. No right-handed neutrino.
 - ▶ From oscillation experiment ($m_\nu \neq 0$)
- SM can be considered as low energy effective model.

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{\mathcal{L}_{d=5}}{\Lambda} + \frac{\mathcal{L}_{d=6}}{\Lambda^2} + \dots$$

- ▶ For $d = 5$. Weinberg operator.
 - ★ Type-I seesaw
- ▶ For $d = 6$. NSI

Type-I seesaw

- Introduce right-handed neutrinos
- Allow L number violation

$$\mathcal{L}_{mass}^\nu \supset Y_\nu \bar{L}_L \tilde{\phi} N_R + \frac{1}{2} M_R \bar{N}_R^c N_R + h.c.$$

$$m_\nu \sim \frac{Y_\nu^\dagger Y_\nu v^2}{M_R} \quad m_N \approx M_R + \mathcal{O}(m_\nu)$$

- For $M_R \gg v$
 - ▶ Neutrino masses can be smaller than fermion masses
 - ▶ Heavy neutrinos can hardly be tested
 - ▶ Worsen hierarchy problem

Low-scale Type-I seesaw

- Y_ν small
- They may be tested in experiments with meson decays and muon decays
- Right-handed neutrinos with very high masses can be a partial solution for other problems
 - ▶ keV neutrino can be a candidate for dark matter
[A. Kusenko, Phys. Rept. 481(2009) 128]
 - ▶ $m_N \sim \mathcal{O}(1 - 100)$ GeV, majorana neutrinos can generate enough matter-antimatter asymmetry of the Universe
[T. Asaka and M. Shaposhnikov, Phys. Lett. B 620(2005) 1726]

NSI-NC

- Described by effective four-fermion operators

$$\mathcal{L} = -2\sqrt{2}G_F\epsilon_{\alpha\beta}^{fP}(\bar{\nu}_\alpha\gamma^\mu\nu_\beta)(\bar{f}\gamma_\mu Pf)$$

- ▶ Not gauge invariant
 - ▶ Charged lepton flavour violation processes impose tight constraints
[M. B. Gavela, D. Hernandez, T. Ota and W. Winter, Phys. Rev. D79 (2009) 013007]
 - ▶ NSI are generated well below de EW scale
[K. S. Babu, A. Friedland, P. A. N. Machado and I. Mocioiu, arXiv:1705.01822]
- Modify the forward -coherent scattering in regions with matter
 - Can be constrained by measuring neutrino cross section with other fermions
[C. Biggio, M. Blennow and E. Fernandez-Martinez, JHEP 08 (2009) 090,
S. Davidson, C. Pena-Garay, N. Rius and A. Santamaria, JHEP 03 (2003) 011]

Constraints to NSI-NC parameter from oscillation experiments

	90%CL	3 σ CL
$\epsilon_{ee}^q - \epsilon_{\mu\mu}^q$	[+0.02, +0.51]	[-0.09, +0.71]
$\epsilon_{\tau\tau}^q - \epsilon_{\mu\mu}^q$	[-0.01, +0.03]	[-0.03, +0.19]
$\epsilon_{e\mu}^q$	[-0.09, +0.04]	[-0.16, +0.11]
$\epsilon_{e\tau}^q$	[-0.13, +0.14]	[-0.38, +0.29]
$\epsilon_{\mu\tau}^q$	[-0.01, +0.01]	[-0.03, +0.03]

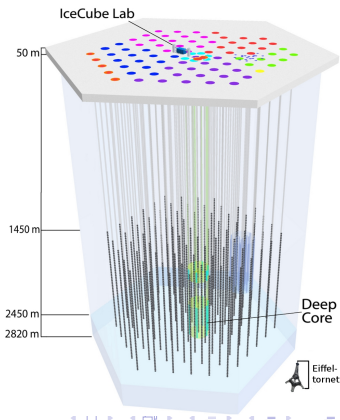
[M. C. Gonzalez-Garcia and M. Maltoni, JHEP09(2013)152]

- IceCube

- ▶ Triangular grid of strings / horizontal distance of 125m
- ▶ 78 vertical strings
- ▶ 60 DOMs per string / vertical distance of 17m

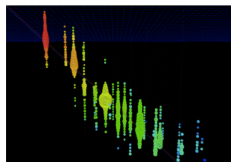
- DeepCore

- ▶ 8 closely-spaced strings + 7 central IceCube strings
- ▶ Horizontal distance of 72m
- ▶ 50 DOMs (vertical distance 7m) + 10 DOMs (10m)



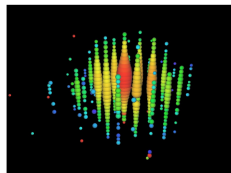
Tracks

- ▶ Through-going muons
- ▶ Direction resolution $\leq 1^\circ$



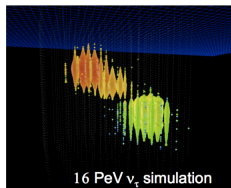
Showers

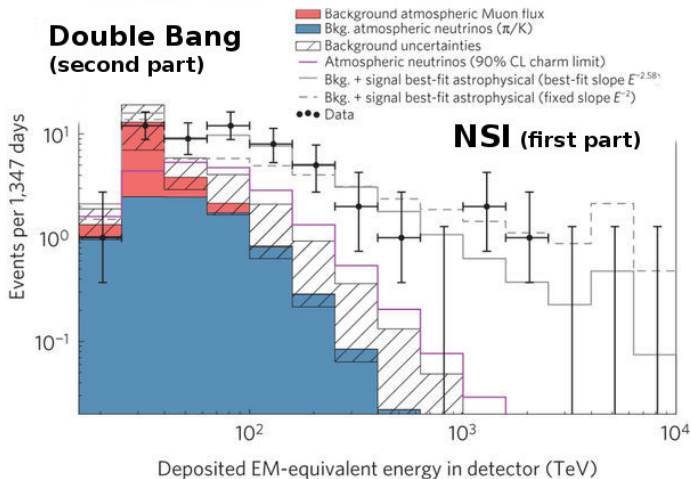
- ▶ ν_α NC + ν_e and ν_τ CC
- ▶ Energy resolution $\sim 10\%$



Composites

- ▶ Starting ν_μ CC and high energy ν_τ CC
- ▶ Good energy and direction resolutions





[Figure taken from: Francis Halzen, Nature Physics 13, 232238 (2017)]

New Physics at high energies in IceCube/DeepCore

New physics at high energies

Considering $d = 6$ effective field operator (NSI)

- Different mechanism can contribute flux of high energy neutrino.
 - ▶ Pion-muon decay chain, single μ -flux, neutron decay, decay charm mesons, ...
 - ▶ Flavor composition at the detector can allow to differentiate.
 - ▶ Standard oscillation modify the flavor content.
 - ★ Uncertainties on oscillation parameters ($\theta_{12}, \theta_{23}, \theta_{31}, \delta_{CP}$) determines the possible flavor content.
- New Physics: non-standard interactions of the neutrinos in the Earth matter.
 - ▶ Introduce new flavor oscillation when neutrino travel through the earth.
 - ▶ We quantify the modification of the neutrino flavor composition at the detector in terms NSI parameters.

Neutrino fluxes at the detector (at the earth surface $d \rightarrow \oplus$)

$$\phi_{\beta}^d(E) = \sum_{\alpha} \int dE' \mathcal{P}_{\alpha\beta}^{s \rightarrow d}(E, E') \phi_{\alpha}^s(E')$$

- Coherent evolution:

$$\mathcal{P}_{\alpha\beta}^{s \rightarrow d}(E, E') = \left| \sum_{\gamma} A_{\alpha\gamma}^{s \rightarrow \oplus} A_{\gamma\beta}^{\oplus \rightarrow d}(E) \right|^2 \delta(E - E')$$

- Incoherent evolution:

- ▶ Considering only the dominant attenuation factors

$$\mathcal{P}_{\alpha\beta}^{s \rightarrow d}(E, E') \simeq \left| \sum_{\gamma} A_{\alpha\gamma}^{s \rightarrow \oplus} A_{\gamma\beta}^{\oplus \rightarrow d}(E) \right|^2 F_{att}^{\oplus \rightarrow d} \delta(E - E')$$

- ▶ No flavor distortion

Non-standard interaction formalism

The matter part of the hamiltonian

$$H_{mat} = \sqrt{2}G_F N_e(r) \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{e\mu}^* & \epsilon_{e\tau} \end{pmatrix} = W D_{mat} W^\dagger$$

$$\epsilon_{\alpha\beta} = \epsilon_{\alpha\beta}^e + Y_u \epsilon_{\alpha\beta}^u + Y_d \epsilon_{\alpha\beta}^d$$

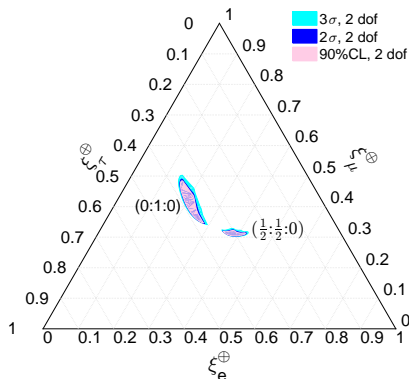
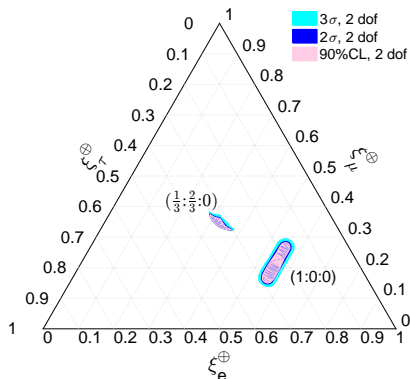
$$H_{vac} = U D_{vac} U^\dagger$$

The oscillation probabilities are

$$\left| \sum_{\gamma} A_{\alpha\gamma}^{s \rightarrow \oplus} A_{\gamma\beta}^{\oplus \rightarrow d}(E) \right|^2 = \sum_i |U_{\alpha i}|^2 |U_{\beta i}|^2 - 2 \sum_{\gamma\eta k l i} \mathcal{R}(W_{\beta k} W_{\beta l}^* W_{\gamma l} W_{\eta k}^* U_{\eta i} U_{\gamma i}^* |U_{\alpha i}|) \sin^2 \left(d_e \frac{\Delta\epsilon_{kl}}{2} \right) \\ + \sum_{\gamma\eta k l i} \mathcal{I}(W_{\beta k} W_{\beta l}^* W_{\gamma l} W_{\eta k}^* U_{\eta i} U_{\gamma i}^* |U_{\alpha i}|) \sin(d_e \Delta\epsilon_{kl})$$

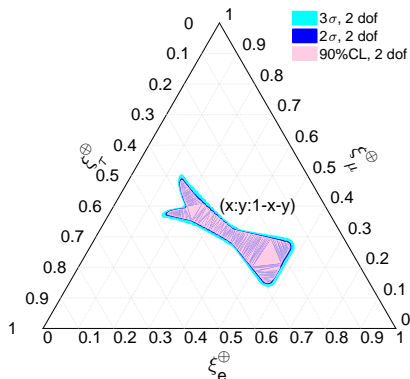
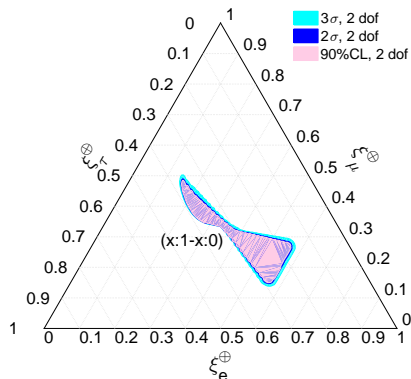
Results: Standard oscillation

- Flavor ratio $\xi_\beta^\oplus = \phi_\beta^\oplus / \sum_\gamma \phi_\gamma^\oplus = \sum_\alpha \sum_i |U_{\alpha i}|^2 |U_{\beta i}|^2 \xi_\alpha^s$
- Projection of six oscillation parameters χ^2 of the global NuFIT analysis.
[M. C. Gonzalez-Garcia, M. Maltoni and T. Schwetz, JHEP11(2014)052]



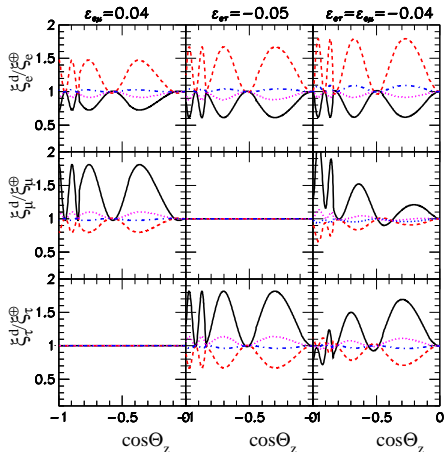
Results: Standard oscillation

Several production processes in the source



NSI in the earth

New oscillations induced by $\epsilon_{\alpha\neq\beta} \neq 0$.



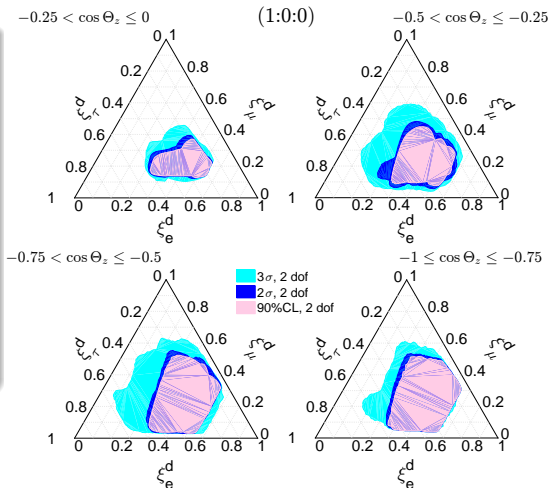
$$\begin{aligned} \xi_{\beta}^d &= \phi_{\beta}^d / \sum_{\gamma} \phi_{\gamma}^d \\ &= \sum_{\alpha} \left| \sum_{\gamma} A_{\alpha\gamma}^{s \rightarrow \oplus} A_{\gamma\beta}^{\oplus \rightarrow d}(E) \right|^2 \end{aligned}$$

- (1:0:0)
- (0:1:0)
- ⋯ (1/2, 1/2, 0)
- ⋯ (1/3, 2/3, 0)

Flavor ratios in presence of NSI

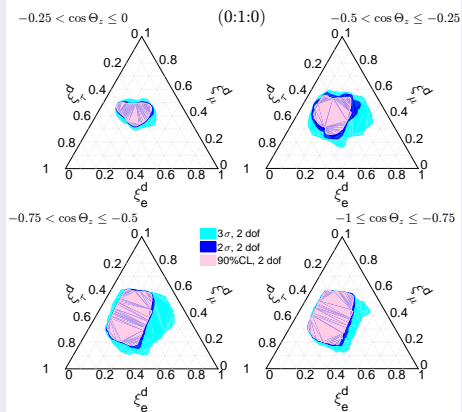
Neutron decay

- The range of flavor ratios increase.
- Included $(1/3, 1/3, 1/3)$.
- The largest values of $\epsilon_{\alpha\beta}$ averaged out.
- The largest effect for "less equal" ratios at the earth surface.

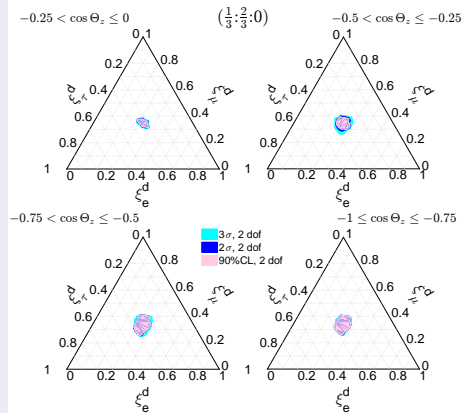


Flavor ratios in presence of NSI

Single μ -flavor flux.



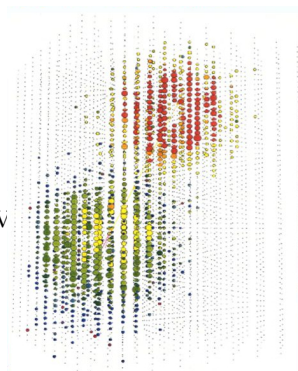
No large deviation from $(1/3, 1/3, 1/3)$ for fluxes produced by $\pi - \mu$ decay chain.



New Physics at low energies in IceCube/DeepCore

Double bang events

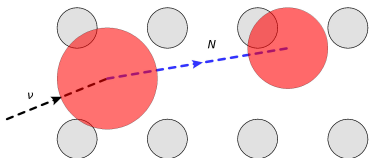
- Standard signature of ν_τ at very high energy
- ν_τ CC interaction produce a τ and a shower (1 shower)
- τ decay (2 shower)
 - ▶ τ emits cherenkov radiation
- For very well-separated showers ($\sim 100\text{m}$) $E_{\nu_\tau} \geq 2\text{PeV}$
- Background negligible
- Not detected yet



New physics at low energies

Double bang signals to look for new physics at low energy

- Two bangs inside the detector
 - ▶ 1st shower ν interaction
 - ▶ 2nd shower N decay
 - ▶ No cherenkov radiation in between



What kind of new physics?

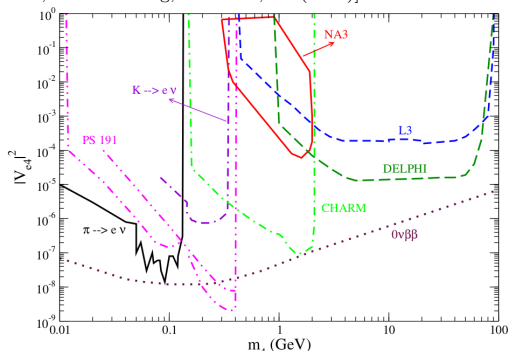
1. BSM: Heavy sterile neutrino

Sterile mass mixing with active neutrinos

$$\nu_{\alpha L} = \sum U_{\alpha m} \nu_{mL} + U_{\alpha 4} N_{4L}$$

In the presence of $\nu - N - Z$ interaction: strong bounds on the mixing between N and ν_e, ν_μ

[A. Atre, T. Han, S. Pascoli, and B. Zhang, JHEP 05,030 (2009)]



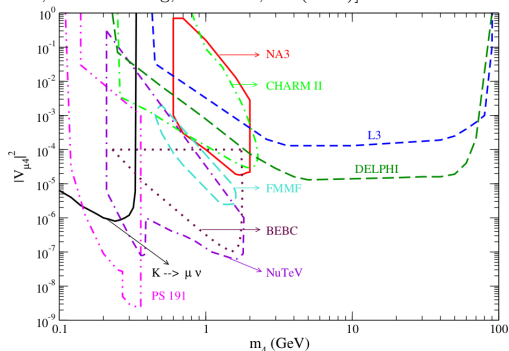
1. BSM: Heavy sterile neutrino

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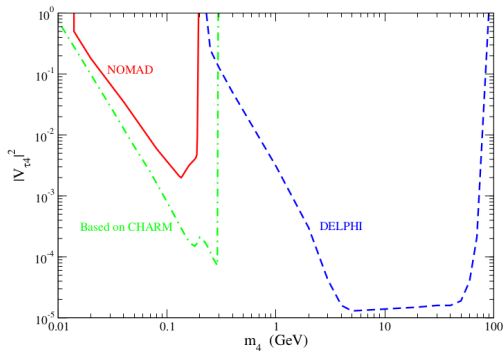
1. Standard scenario: Heavy sterile neutrino

Sterile mass mixing with active neutrinos

$$\nu_{\alpha L} = \sum U_{\alpha m} \nu_{mL} + U_{\alpha 4} N_{4L}$$

In the presence of $\nu - N - Z$ interaction: we are going to constrain $U_{\tau 4}$

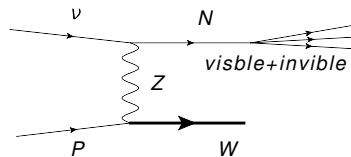
[A. Atre, T. Han, S. Pascoli, and B. Zhang, JHEP 05,030 (2009)]



Sterile neutrino via the Neutral Current

The double bang signal comes from

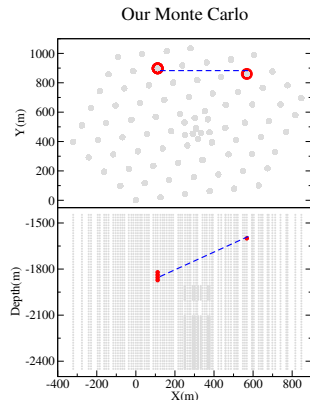
$$\begin{aligned}\nu_\tau + N &\rightarrow N_4 + W \\ N_4 &\rightarrow \text{visible} + \text{invisible}\end{aligned}$$



- The decay length depends on M_4 and on $|U_{\tau 4}|^2$
- Cross section calculated with GENIE (Coherence + Resonance + DIS)
 - ▶ Proportional to mixing parameter $|U_{\tau 4}|^2$

Effective Volume

- Double Pulse (2 separate showers in the full detector)
 - ▶ Minimum distance between showers defined by DOMs resolution wave form
 - ▶ $\geq 20\text{m}$ between showers
- Energy threshold of 5GeV per shower
 - ▶ Minimum energy detected by DeepCore
- Maximum distance covered by light of 36m
- Simulation include DOMs position and triggers
 - ▶ SMT3 for DeepCore
 - ▶ SMT4 for IceCube
- Background
 - ▶ Coincident atmospheric cascades
 - ▶ 0.05/year

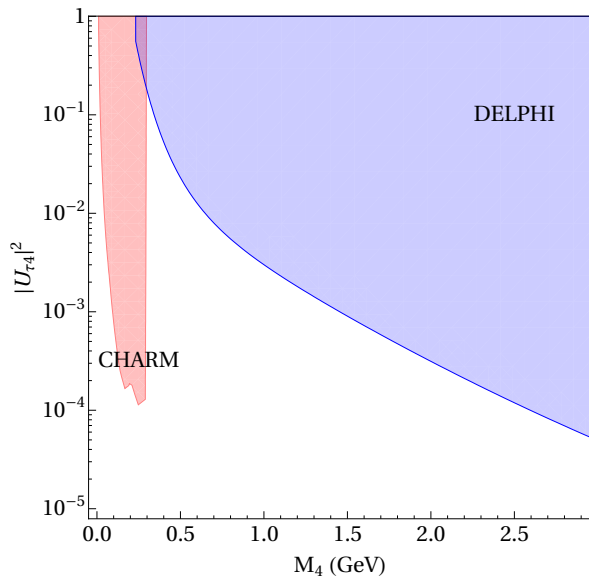


The number of events in the detector is given by

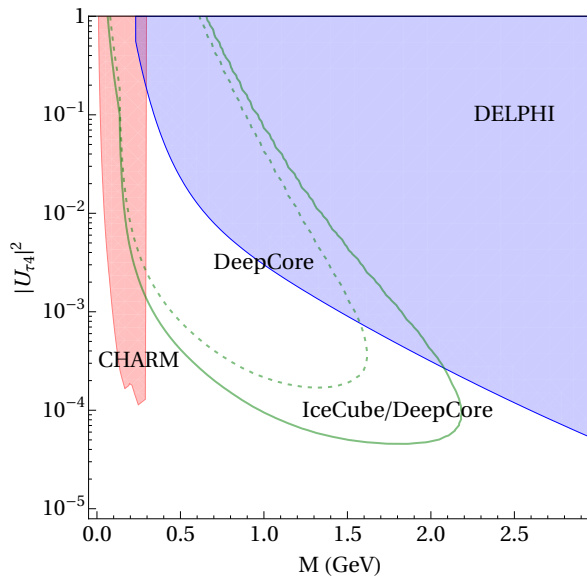
$$N(L) \propto T \int dE d \cos \theta dE' \frac{d\phi_{\nu\mu}}{dE d \cos \theta} P_{\mu \rightarrow \tau}(E, \cos \theta) \frac{d\sigma_{\nu\tau\nu A}}{dE dE'} P_d(L) V_{eff}(L, \cos \theta)$$

- We consider $E \in [10, 100]$ GeV
 - ▶ The energy of the heavy neutrino $5\text{GeV} \leq E' \leq E - 5\text{GeV}$
 - ▶ The showers $\geq 5\text{GeV}$
- $\phi_{\nu\mu}$ atmospheric flux
 - ▶ $\phi \sim E^{-2.7}$ The largest contribution comes from low energy neutrinos
- $P_{\mu \rightarrow \tau}$ 3 neutrino oscillation
- Decay probability $P_d(L) = e^{-L/\Gamma}/\Gamma$
- The results correspond with 6 years.

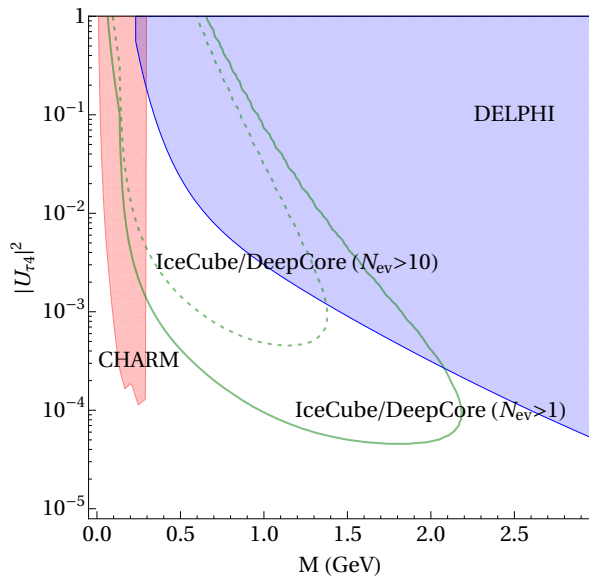
Results: Neutral Currents



Results: Neutral Currents



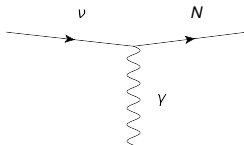
Results: Neutral Currents



2. Neutrino magnetic moment

- We are interested in a transition magnetic moment
- Weak constraints

$$\mathcal{L} \supset -\mu_\nu \bar{N}_4 \sigma_{\mu\nu} P_L \nu_\alpha F^{\mu\nu}$$



The main contribution to our signal events comes from DIS on nucleons

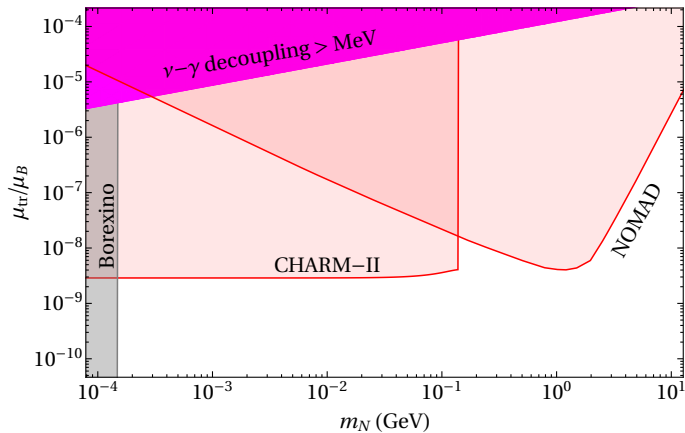
$$\frac{d^2\sigma_N}{dx dy} = g_e^2 \mu_\nu^2 \left(\sum_q e_q^2 f_q(x) \right) \left(\frac{(2-y)^2}{y} - y \right)$$

The decay length $N \rightarrow \nu_i \gamma$

$$\Gamma = \frac{\mu_\nu^2 M_4^3}{16}$$

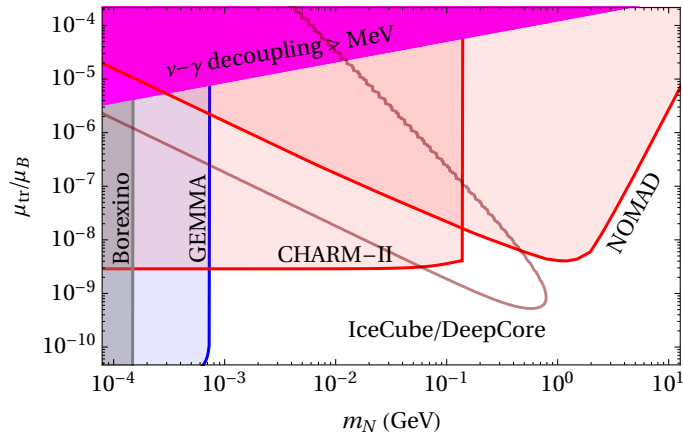
Results: Magnetic moment

$\nu_\mu - N$ transition



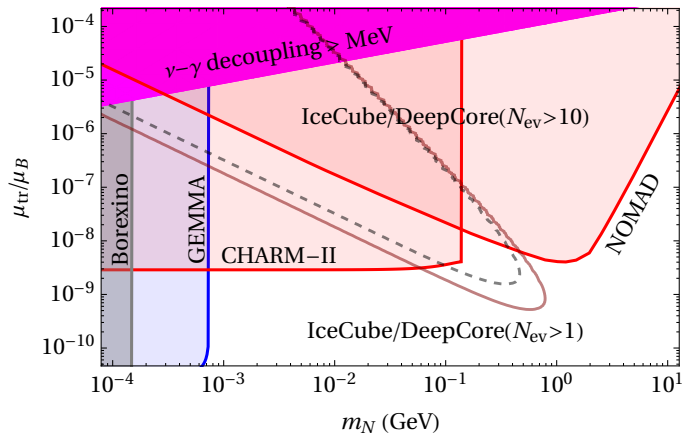
Results: Magnetic moment

$\nu_\mu - N$ transition



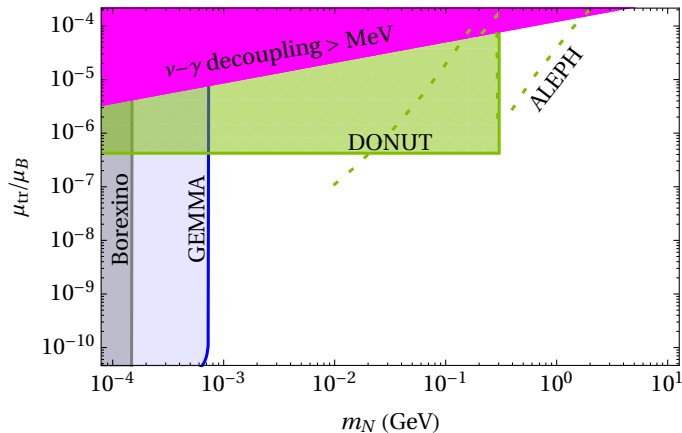
Results: Magnetic moment

$\nu_\mu - N$ transition



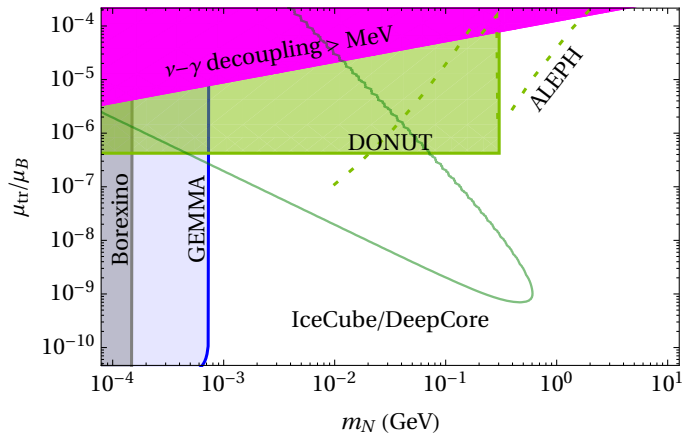
Results: Magnetic moment

$\nu_\tau - N$ transition



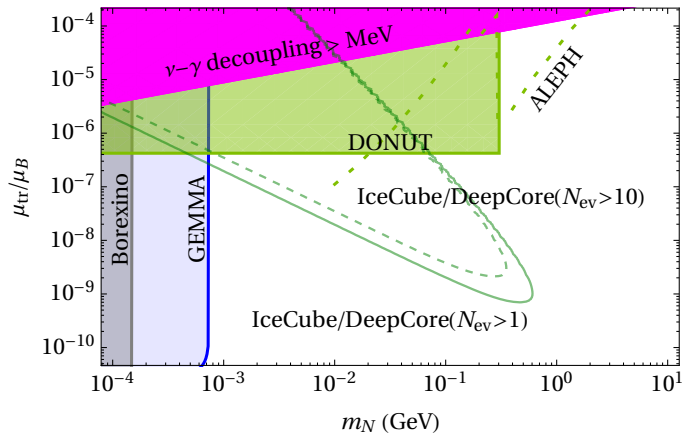
Results: Magnetic moment

$\nu_\tau - N$ transition



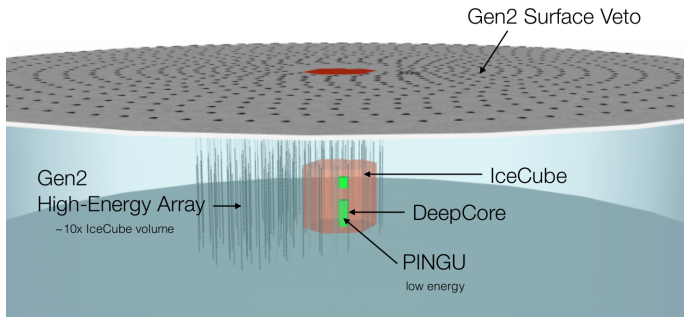
Results: Magnetic moment

$\nu_\tau - N$ transition



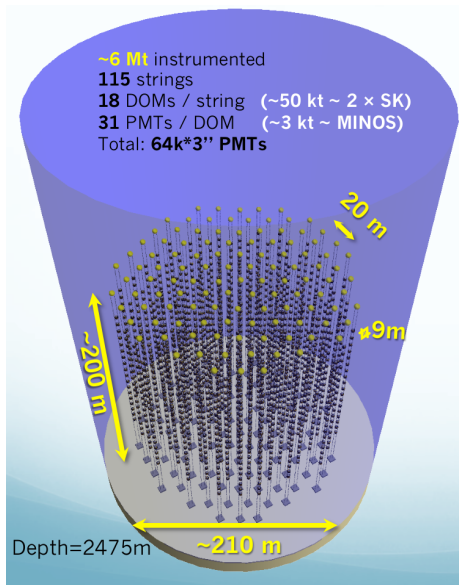
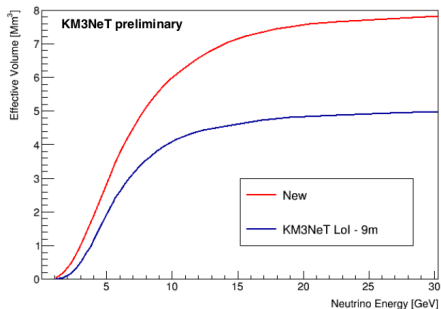
Future Projects: IceCube-Gen2

- PINGU
 - ▶ Effective mass of 6 Mton
 - ▶ Energy threshold 1 GeV
 - ▶ Increase the number of events by a factor 1.6
- High energy array
 - ▶ Fiducial volume 10 km^3 (120 additional strings)
 - ★ Reduce the lower limit of m_N by a factor 1.4
 - ▶ Increase in the effective volume ~ 2.7
- Total increase of the number of events by a factor 3.3



Future Projects: KM3NeT

- Proposed experiment in Mediterranean sea.
- Energy threshold of 3 GeV.
- Effective mass of 3.5 Mton for neutrinos with 10 GeV
- The number of events expected is a factor 0.3 smaller.



Conclusion

- IceCube can probe BSM at different scales of energy

High energy (PeV)

- NSI in the earth can modify the neutrino flavor at the detector.
- The ranges for all neutrino flavor fractions at the detector increase in presence of NSI.
 - ▶ The effect dominates for fluxes with one flavor at source.

Low Energy (GeV)

- Double Bang signals can probe new physics
- Sterile neutrino via neutral current
- Neutrino transition magnetic moment
 - ▶ IceCube can put a competitive bound on μ_ν for ν_τ and ν_μ

Thank you very much!

Constraints to NSI parameter from oscillation experiments + scattering

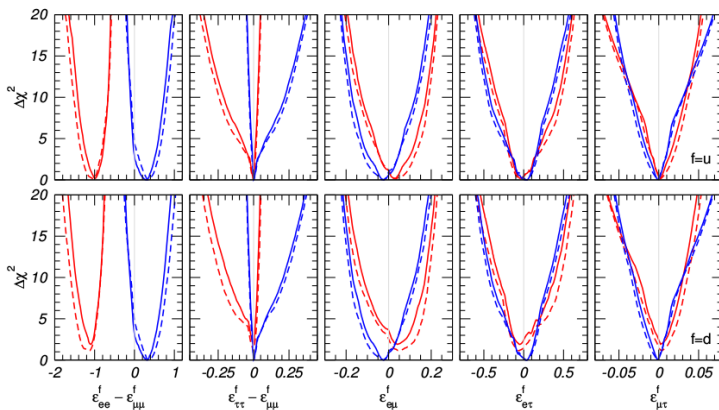
[P. Coloma, P. B. Denton, M. C. Gonzalez-Garcia, M. Maltoni, and T. Schwetz, JHEP 04 (2017) 116]

Light			
PRESENT (OSC)		+COHERENT(SM)	
$\epsilon_{ee}^{u,V} - \epsilon_{\mu\mu}^{u,V}$	$[-1.19, -0.81] \oplus [0.00, 0.51]$	$\epsilon_{ee}^{u,V}$	$[0.002, 0.049] \oplus [0.28, 0.42]$
$\epsilon_{\tau\tau}^{u,V} - \epsilon_{\mu\mu}^{u,V}$	$[-0.03, 0.03]$	$\epsilon_{\mu\mu}^{u,V}$	$[-0.026, 0.033] \oplus [0.36, 0.38]$
		$\epsilon_{\tau\tau}^{u,V}$	$[-0.025, 0.047] \oplus [0.36, 0.39]$
$\epsilon_{\mu\mu}^{u,V}$	$[-0.09, 0.10]$		$[-0.08, 0.04]$
$\epsilon_{e\tau}^{u,V}$	$[-0.15, 0.14]$		$[-0.17, 0.14]$
$\epsilon_{\mu\tau}^{u,V}$	$[-0.01, 0.01]$		$[-0.01, 0.01]$
$\epsilon_{ee}^{d,V} - \epsilon_{\mu\mu}^{d,V}$	$[-1.17, -1.03] \oplus [0.02, 0.51]$	$\epsilon_{ee}^{d,V}$	$[0.022, 0.023] \oplus [0.25, 0.38]$
$\epsilon_{\tau\tau}^{d,V} - \epsilon_{\mu\mu}^{d,V}$	$[-0.01, 0.03]$	$\epsilon_{\mu\mu}^{d,V}$	$[-0.024, 0.029]$
		$\epsilon_{\tau\tau}^{d,V}$	$[-0.023, 0.039]$
$\epsilon_{\mu\mu}^{d,V}$	$[-0.09, 0.08]$		$[-0.07, 0.04]$
$\epsilon_{e\tau}^{d,V}$	$[-0.13, 0.14]$		$[-0.14, 0.12]$
$\epsilon_{\mu\tau}^{d,V}$	$[-0.01, 0.01]$		$[-0.009, 0.007]$
Heavy			
PRESENT (OSC+CHARM+NuTeV)		+COHERENT(SM)	
$\epsilon_{ee}^{u,V}$	$[-0.97, -0.83] \oplus [0.033, 0.450]$		$[0.014, 0.032] \oplus [0.24, 0.41]$
$\epsilon_{\mu\mu}^{u,V}$	$[-0.008, 0.005]$		$[-0.007, 0.005]$
$\epsilon_{\tau\tau}^{u,V}$	$[-0.015, 0.04]$		$[-0.006, 0.04]$
$\epsilon_{\mu\mu}^{u,V}$	$[-0.05, 0.03]$		$[-0.05, 0.03]$
$\epsilon_{e\tau}^{u,V}$	$[-0.15, 0.13]$		$[-0.15, 0.13]$
$\epsilon_{\mu\tau}^{u,V}$	$[-0.006, 0.005]$		$[-0.006, 0.004]$
$\epsilon_{ee}^{d,V}$	$[0.02, 0.51]$		$[0.26, 0.38]$
$\epsilon_{\mu\mu}^{d,V}$	$[-0.003, 0.009]$		$[-0.003, 0.009]$
$\epsilon_{\tau\tau}^{d,V}$	$[-0.001, 0.05]$		$[-0.001, 0.05]$
$\epsilon_{\mu\mu}^{d,V}$	$[-0.05, 0.03]$		$[-0.05, 0.03]$
$\epsilon_{e\tau}^{d,V}$	$[-0.15, 0.14]$		$[-0.15, 0.14]$
$\epsilon_{\mu\tau}^{d,V}$	$[-0.007, 0.007]$		$[-0.007, 0.007]$

Transition magnetic moment. Electron cross section.

$$\frac{d\sigma_e}{d\nu} = \mu_\nu^2 \alpha_{em} \left(\frac{(\nu - M_e) M_4^4}{8\nu^2 E^2 M_e^2} + \frac{(\nu - 2E - M_e) M_4^2}{4\nu E^2 M_e} + \frac{1}{\nu} - \frac{1}{E} \right)$$

Constraints to NSI parameter from oscillation experiments



[M. C. Gonzalez-Garcia and M. Maltoni, JHEP09(2013)152]