

[GLIMPSES of]



# CP violation theory: SM & more

Amarjit Soni  
HET-BNL

Heavy Quarks and Leptons

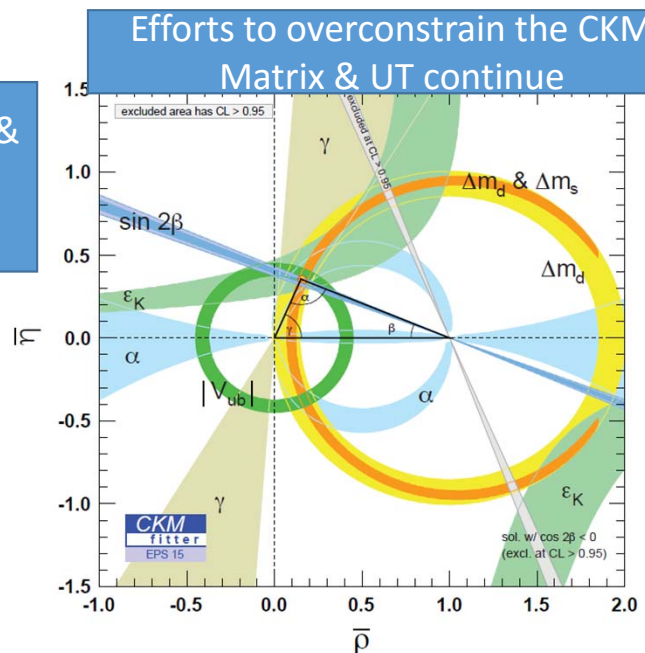
**Virginia Tech, 05/26/16**

[APOLOGIES: HIGHLY SUBJECTIVE]

# Outline

- **SM-CKM**
- **Recent update & near future expectations**
- **BSM**
- **Motivation .....theory**
- **Possible experimental hints.....[some critique]**
- **Strategies for the near future for BSM-CP searches**
- **Summary & Outlook**

Andreas Hoecker &  
Malcolm John  
EW Moriond '16



Compatible with SM-CKM to  
~15% accuracy

O(5-10%) new physics is possible  
and is HUGE

Key new results from LHCb

DATA DRIVEN methods

Precision on  $\sin(2\beta)$  approaches that of  
B-factories:  $0.73 \pm 0.04 \pm 0.02$

ITE ~ 1%  
Mannel et al

- A world-leading measurement of  $\gamma$  is made from a combination of LHCb analysis, concluding with

$$\gamma = 70.9^{+7.1}_{-8.5}$$

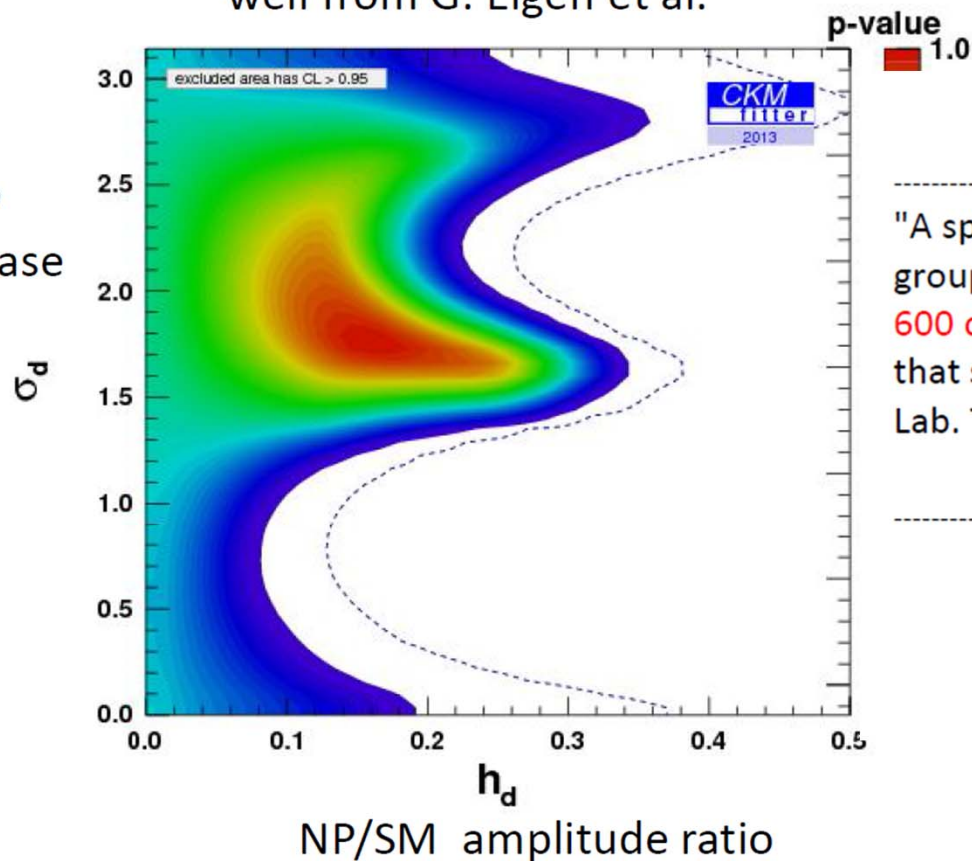
ITE ~  $10^{-7}$  !!  
Brod Zupan'14  
STD. CANDLE

which improved the previous LHCb-only conclusion by  $2^\circ$

- Inline with B-factory conclusions from  $B \rightarrow DK$ ,
  - BaBar:  $\gamma = (70 \pm 18)^\circ$
  - Belle:  $\gamma = (73^{+13}_{-15})^\circ$

BELLE-II & LHCb(upgrade) ~  $\delta\gamma \sim 1^\circ$ , still long  
way to go before ultimate precision

ICHEP2014: Similar results  
from UTFIT (D. Derkach) as  
well from G. Eigen et al.



### A lesson from history (I)

"A special search at Dubna was carried out by E. Okonov and his group. They did not find a single  $K_L \rightarrow \pi^+ \pi^-$  event among 600 decays into charged particles [12] (Anikira et al., JETP 1962). At that stage the search was terminated by the administration of the Lab. The group was unlucky."

-Lev Okun, "The Vacuum as Seen from Moscow"

1964:  $BF = 2 \times 10^{-3}$

A failure of imagination ? Lack of patience ?

# A great personal treat; thanks to LHCb

ADS:  $B^\pm \rightarrow Dh^\pm, D \rightarrow \pi^+ K^-$

$$A_{\text{ADS}(K)}^{\pi K} = -0.403 \pm 0.056 \pm 0.011$$



Malcolm John@EW MORIOND

Huge *direct CP* [tailor made] ~20 years ago!  
ADS PRL'97



*[Recall  $\epsilon' \sim 10^{-6}$ !]* **DESIGNED for MAXIMAL INTERFERENCE**

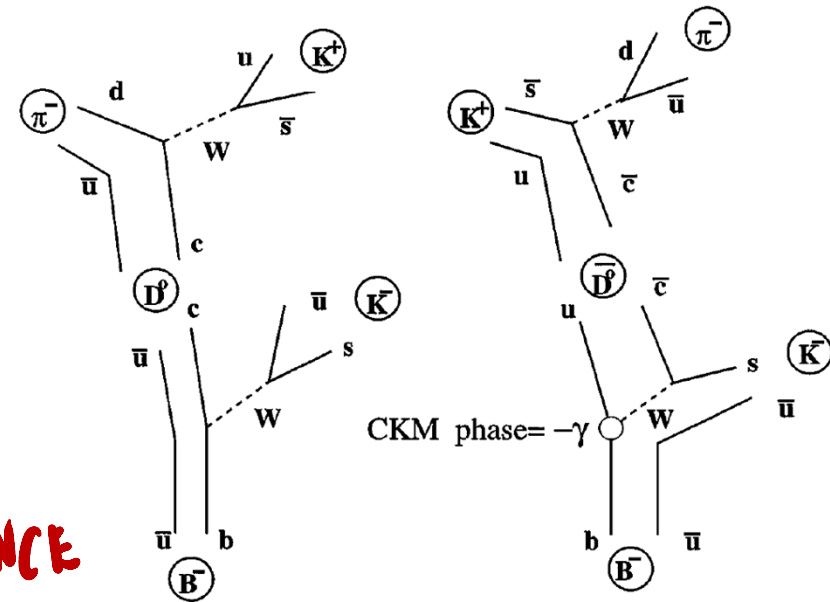
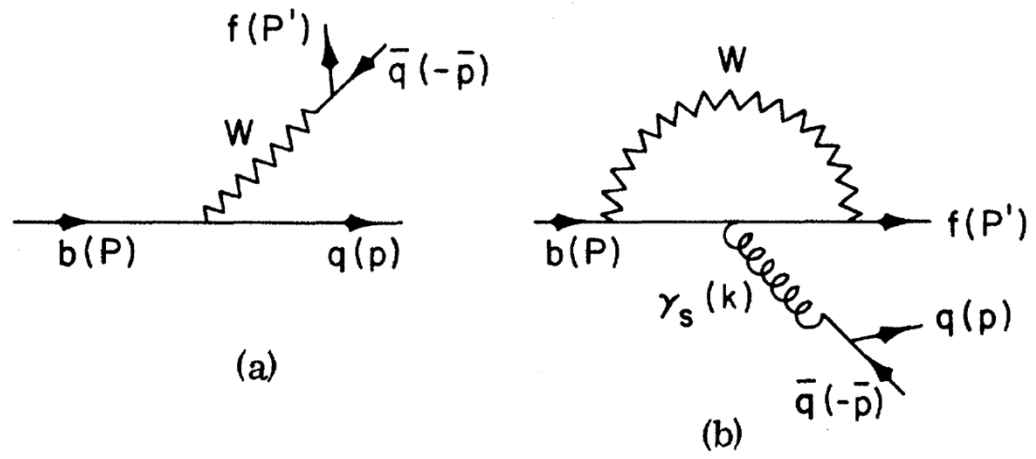


FIG. 1. Diagrams for the two interfering processes:  $B^- \rightarrow K^- D^0$  (color-allowed) followed by  $D^0 \rightarrow K^+ \pi^-$  (double Cabibbo suppressed) and  $B^- \rightarrow K^- \bar{D}^0$  (color-suppressed) followed by  $\bar{D}^0 \rightarrow K^+ \pi^-$  (Cabibbo allowed).

# DIRECT CP: Long-standing challenge for theorists

# Simple example in B-Physics: Tree-Penguin Interference



Bander, Silverman and A. S. PRL '79

**measurable asymmetries may arise.**  
 This would present the first evidence for *CP* noninvariance in charged systems.

$$A = |A_1| \exp[i(\delta_1 + \phi_1)] + |A_2| \exp[i(\delta_2 + \phi_2)]$$

$$\bar{A} = |A_1| \exp[i(\delta_1 - \phi_1)] + |A_2| \exp[i(\delta_2 - \phi_2)] .$$

$$\alpha_{PRA} = \frac{\mathcal{B}(B \rightarrow f) - \mathcal{B}(\bar{B} \rightarrow \bar{f})}{\mathcal{B}(B \rightarrow f) + \mathcal{B}(\bar{B} \rightarrow \bar{f})}$$

$$= \frac{2|A_1| |A_2| \sin \delta \sin \phi}{|A_1|^2 + |A_2|^2 + 2|A_1| |A_2| \cos \delta \cos \phi} ,$$

$$A_{CP}(B^+ \rightarrow \rho^0 K^+) = 0.37 \pm 0.10$$

$$A_{CP}(B^0 \rightarrow K^+ \pi^-) = -0.082 \pm 0.006$$

$$A_{CP}(B_s \rightarrow \pi^+ K^-) = 0.28 \pm 0.04$$

**REGRETTABLY CANNOT BE USED TO  
 RELIABLY TEST THE SM-CKM**

# Update Direct CP in $D^0$ decays

- First measurement using  $0.6 \text{ fb}^{-1}$  of  $D^{*+}$ -tagged  $D^0$  decays [1]

→ LHCb PRL  
2012

$$\Delta A_{CP} = (-0.82 \pm 0.21 (\text{stat}) \pm 0.11 (\text{syst})) \%$$

- New measurement using  $3 \text{ fb}^{-1}$  of  $D^{*+}$ -tagged  $D^0$  decays

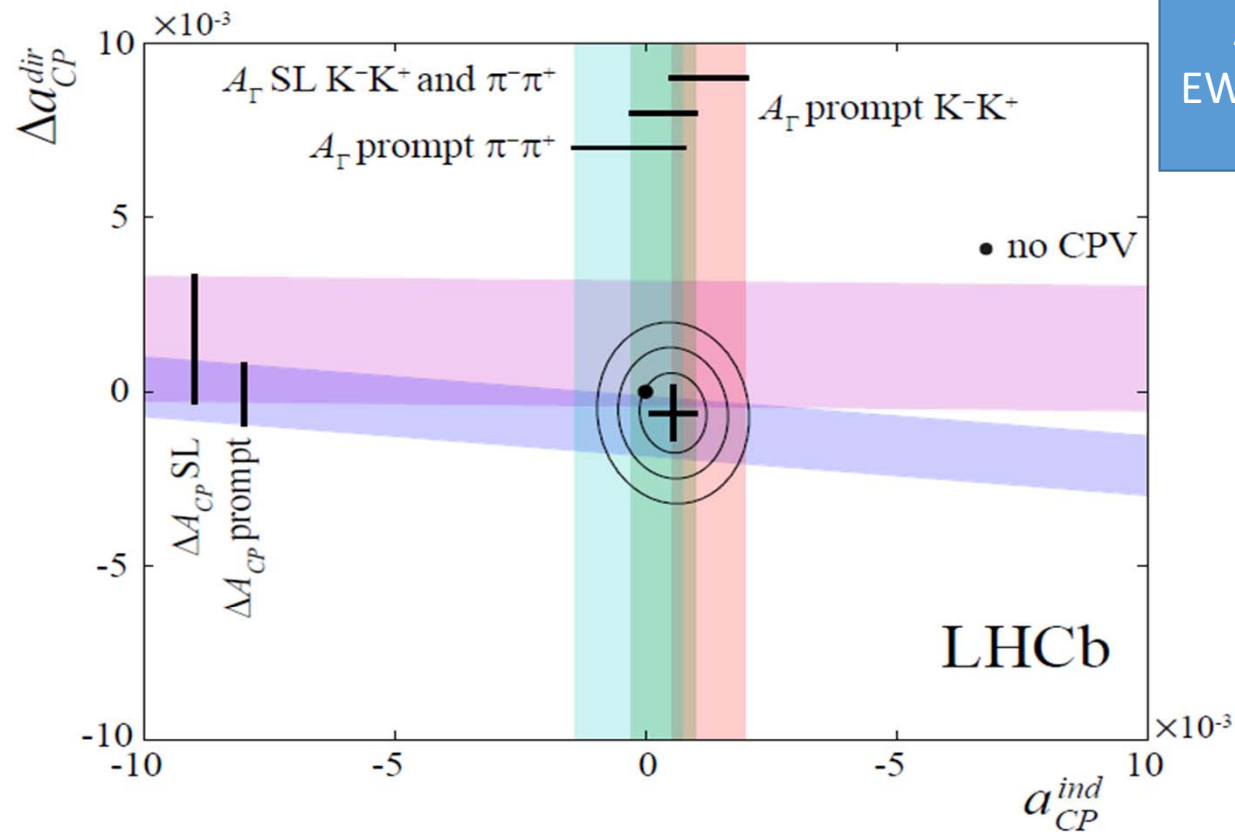
$$\Delta A_{CP} = (-0.10 \pm 0.08 (\text{stat}) \pm 0.03 (\text{syst})) \%$$

arXiv:1602.03160

M. Golden & B. Grinstein, PLB 222,501(1989)  
Feldmann, Nandi and A.S., arXiv:1202.3795  
Brod, Grossman, Kagan & Zupan, arXiv:1203.6659  
Atwood & A. S. arXiv:1211.1026

Alex Pearce, LHCb  
EW Moriond Mar. 2016

THEORY V MURKY  
BSM only if  $\Delta A_{CP} \gtrsim 1\%$



Consistent with  $CP$  conservation hypothesis with  $p = 0.32$

# BSM-CP: Theoretical motivation

- To the extent that SM is not a complete theory, BSM-CP phase(s) are highly likely to exist
- Adding fermions, scalars or gauge bosons entails new phase(s)
- Examples: 4G SM: + 2; LRS : at least + 1; 2HDM : neutral scalar sector as well as charged sector can have new phases; SUSY or WEXD : tens of new CP-odd phases in general may be there
- SM cannot account of baryogenesis.....CKM CP not enough
- Due to all of the above (and some more), searching for BSM phases is just about the best way to look for NP
- **Additional (new) motivation: hints of pretty serious anomalies**

Season of anomalies....countless  
B-physics, K-physics, g-2, ATLAS-  
CMS diboson,  $\sigma(t\bar{t}h)$ ;  $H \Rightarrow \mu\tau$   
tau .....

2 -3.5 sigma effects: while most [and may be  
all] will go away, not following can be a very  
serious mistake

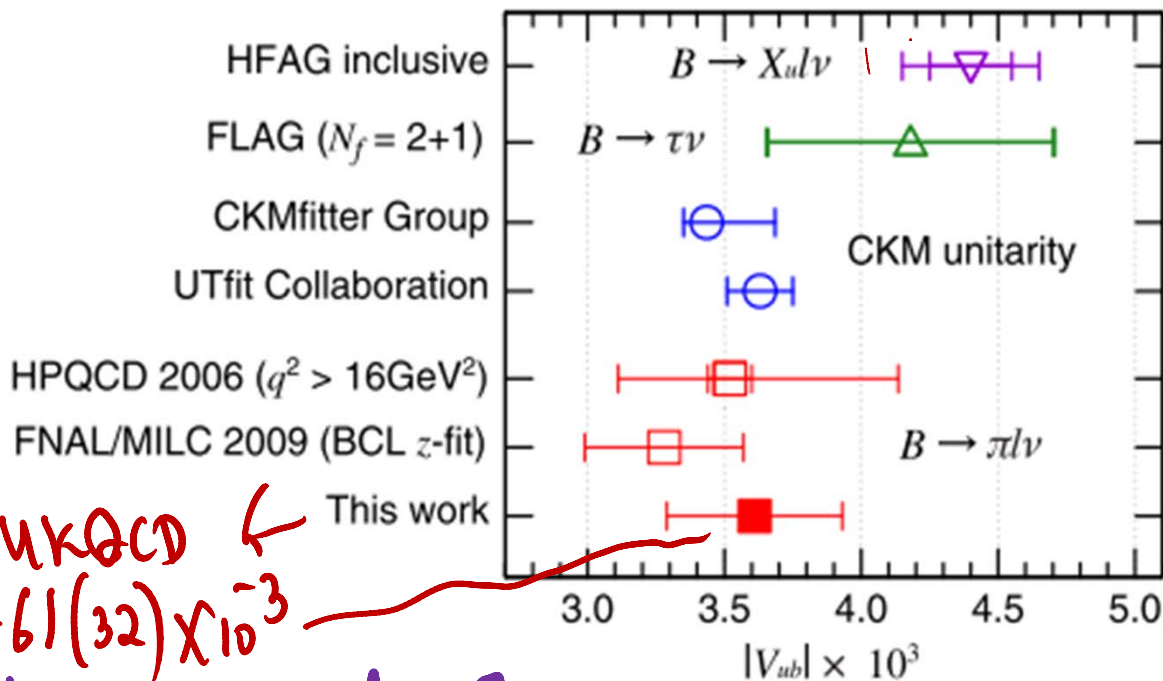
# The long-standing saga of Vub

# Vub “anomaly” i.e inclusive $\Leftrightarrow$ exclusive tension

C also Grinstein

- Exclusive....Experimentally as well as for theory is clean; esp since modes are accessible to the lattice  $\Rightarrow B \Rightarrow \pi \ell \nu$ .....
- Inclusive.....experimental and theoretical difficulties
- Large background; need to discriminate from cascade charm decays.....[Recall  $V_{ub}/V_{cb} \sim 0.08$ ]
- Absence of symmetry [unlike  $b \Rightarrow c \ell \nu$ ]

$B \rightarrow X_u \ell \nu$   
 $\downarrow$  No Charm



$$(4.40 \pm 0.22) \times 10^{-3}$$

)  $\Rightarrow$

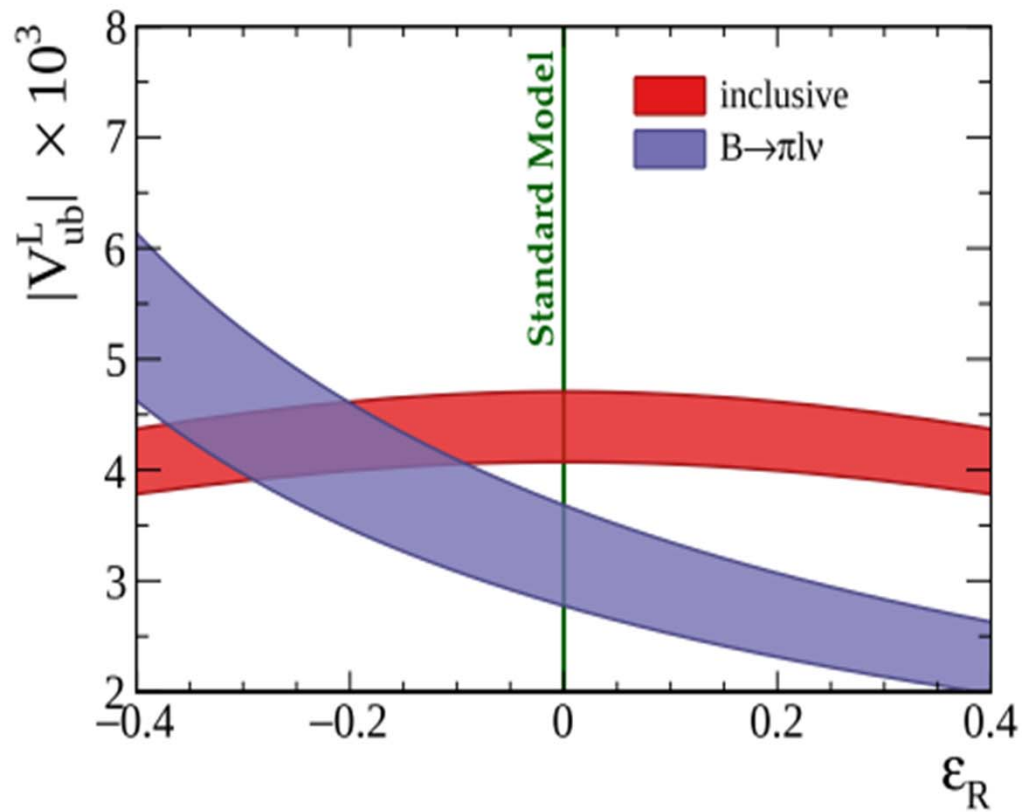
$$\text{MILC 2015}$$

$$3.72(16) \times 10^{-3}$$

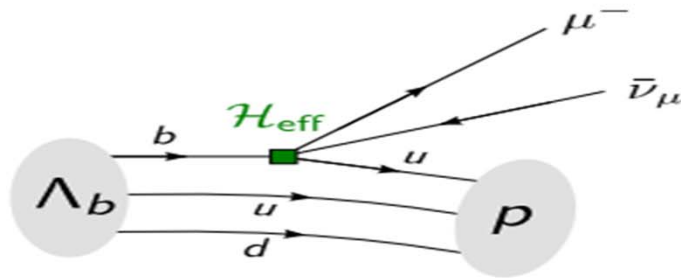
RBC-UKQCD  $\leftarrow$   
 $V_{ub} = 3.61(32) \times 10^{-3}$   
 [USED in pheno analysis]

FIG. 22 (color online). Determinations of  $|V_{ub}|$  from Table XIV. For points with double error bars, the inner error bars are experimental while the outer error bars show the total experimental plus theoretical uncertainty added in quadrature.

Fig from RBC-UKQCD PRD 2015



SOME introduce  
NP [RHC] to  
explain the  
diff



LHCb

+

Lattice

- At LHCb,  $p\mu\bar{\nu}$  final state easier to identify than  $\pi\mu\bar{\nu}$
- Complementary constraints on right-handed coupling

" $\Lambda_b \rightarrow p \ell^- \bar{\nu}_\ell$  and  $\Lambda_b \rightarrow \Lambda_c \ell^- \bar{\nu}_\ell$  form factors from lattice QCD with relativistic heavy quarks"

[W. Detmold, C. Lehner, S. Meinel, arXiv:1503.01421 (~~to appear in PRD~~)]

MIT

HET-BNL

→ RBRC-BNL

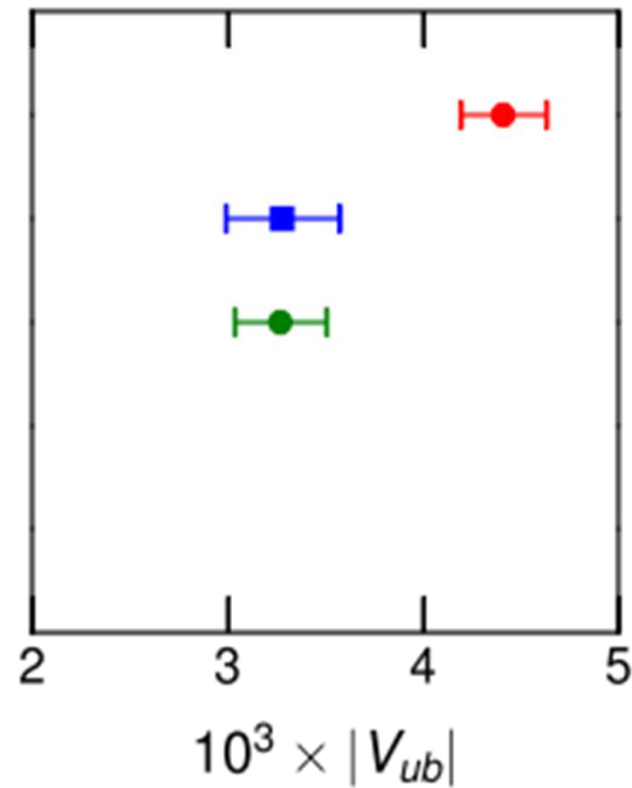
HQL May 2016; V-I ECH; soni[BNL]

$B \rightarrow X_u \ell \bar{\nu}_\ell$  (PDG 2014)

$B \rightarrow \pi \ell \bar{\nu}_\ell$  (PDG 2014)

$\Lambda_b \rightarrow p \ell \bar{\nu}_\ell$  (this work)

Striking consistency between  
completely independent  
exclusive determinations



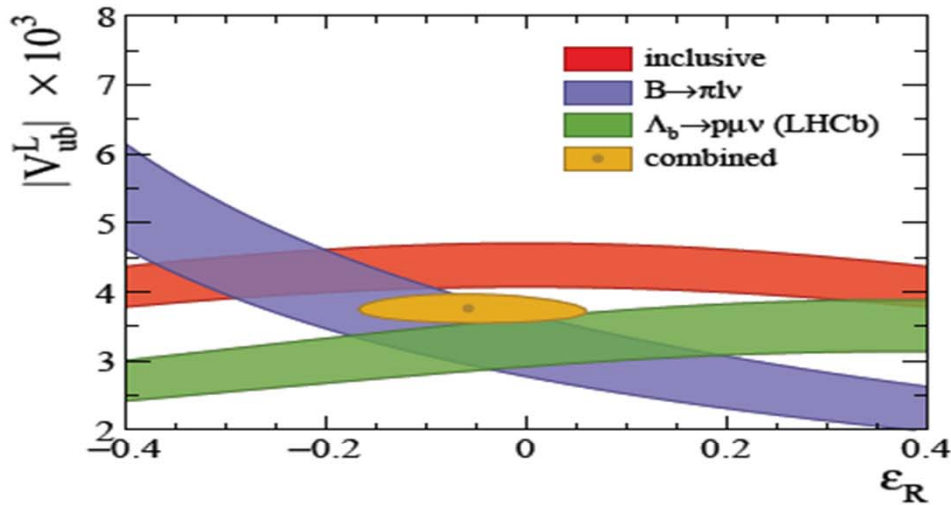


Figure 4: Experimental constraints on the left-handed coupling,  $|V_{ub}^L|$  and the fractional right-handed coupling,  $\epsilon_R$ . While the overlap of the 68% confidence level bands for the inclusive [14] and exclusive [7] world averages of past measurements suggested a right handed coupling of significant magnitude, the inclusion of the LHCb  $|V_{ub}|$  measurement does not support this.

LHCb 1504.01586  
Nature

RULES OUT RHC

# My (subjective) summary on $V_{ub}$

Bearing in mind the clean signal for experiment and also confirmation of exclusive meson versus baryonic extractions of  $V_{ub}$  and the simplicity of exclusive lattice studies

- **Most likely the resolution lies in some small underestimate of errors in both experiment and in theory in inclusive extraction rather than new physics**

## Outlook for next 3-5 years

- **~40 times more data from BELLE II.....**
- **Similarly...lot more data from LHCb anticipated**
- **Factor ~20 or even more computing power now for lattice**
- **Significant progress on this issue expected in ~3 years i.e. expect  $V_{ub}$  errors from lattice studies to go down from ~7% now to around 3-4%**

$R(D(*))$ : concerns

→ BABAR '12

Independent of  $V_{cb}$ !

- To test the SM Prediction, we measure

$$R(D) = \frac{\Gamma(\bar{B} \rightarrow D \tau \nu)}{\Gamma(\bar{B} \rightarrow D \ell \nu)} \quad R(D^*) = \frac{\Gamma(\bar{B} \rightarrow D^* \tau \nu)}{\Gamma(\bar{B} \rightarrow D^* \ell \nu)}$$

Leptonic  $\tau$   
decays only

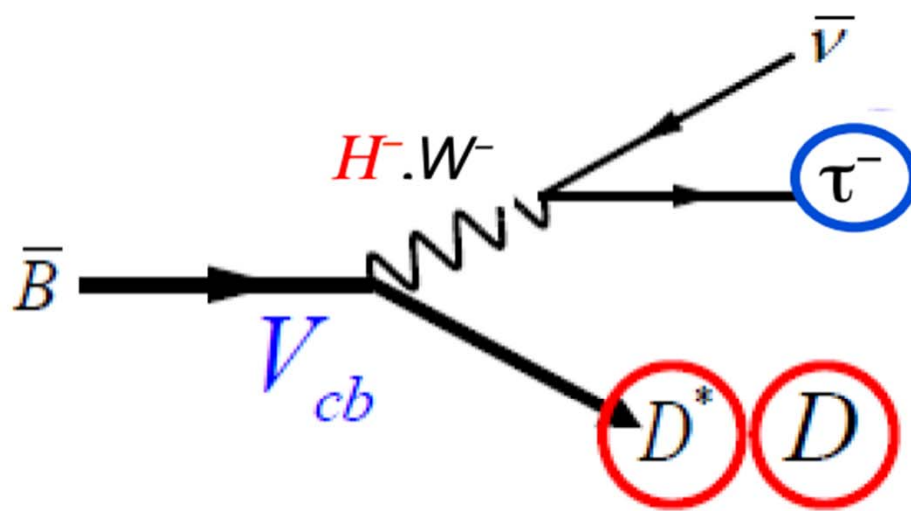
Several experimental and theoretical uncertainties cancel in the ratio!

- $DD^*$  events are fully reconstructed.

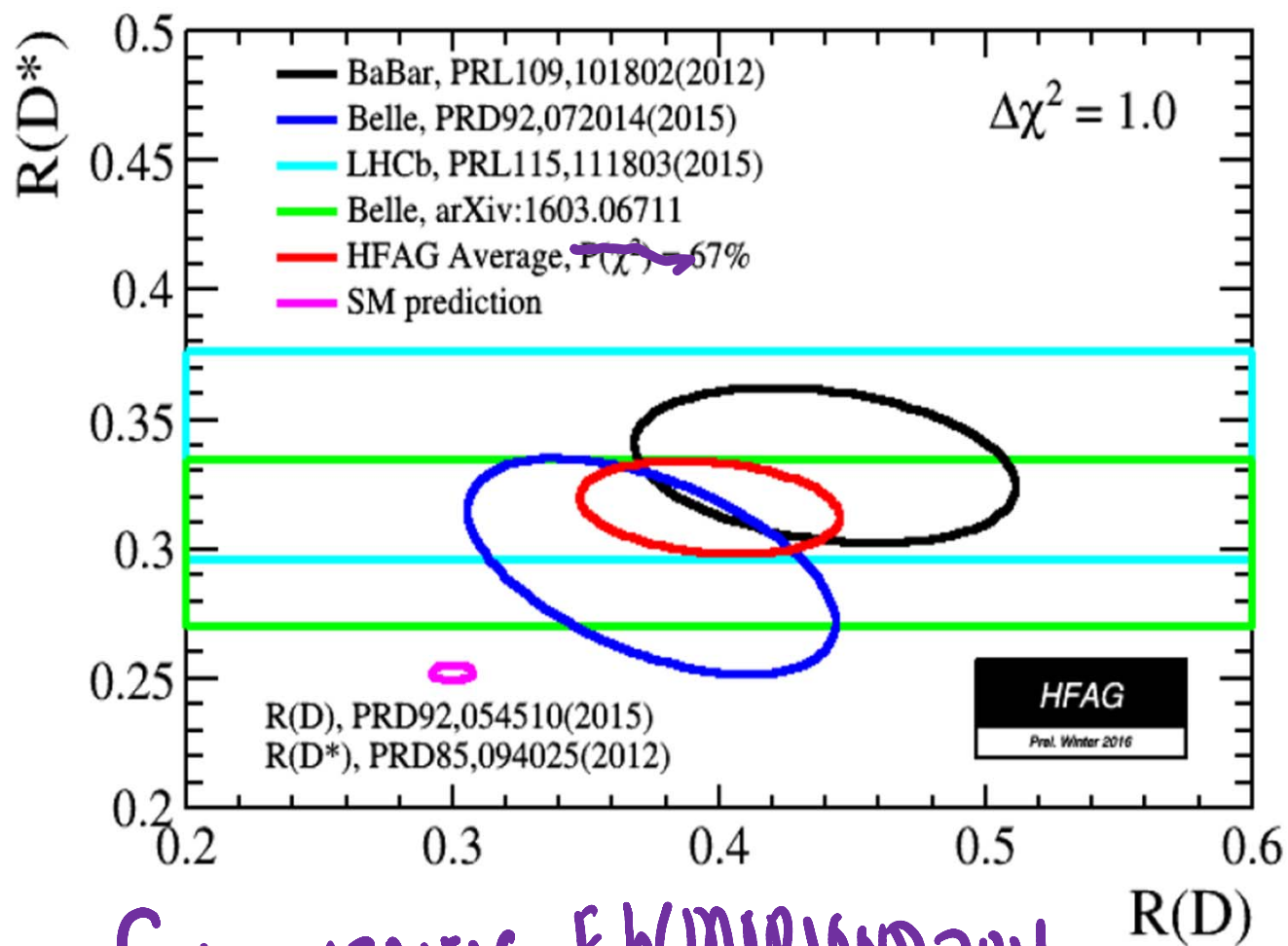
Suggested in theory papers for a long time  
See e.g. Kiers & AS PRD'97; Kamenik et al  
PRD'08; Nierste et al PRD'08

RA LUTH (BABAR)  
CP May 2012  
(HEFEI, China)

Exclusive  $B \rightarrow D^{(*)}\tau\nu$



MANUEL FRANCO  
SEVILLA  
PHD Thesis  
SLAC



GOLDENZWEIG, EW MIRIAD 2016

Lots of excitement  
 for  
 past  $\sim 3$  years!  
 $\sim 3.9\sigma$ ! BUT.....

Decay	$N_{\text{sig}}$	$N_{\text{norm}}$	$R(D^{(*)})$	$\mathcal{B}(B \rightarrow D^{(*)}\tau\nu)$ (%)	$\Sigma_{\text{tot}}(\sigma)$
$D^0\tau^-\bar{\nu}_\tau$	$314 \pm 60$	$1995 \pm 55$	$0.429 \pm 0.082 \pm 0.052$	$0.99 \pm 0.19 \pm 0.13$	4.7
$D^{*0}\tau^-\bar{\nu}_\tau$	$639 \pm 62$	$8766 \pm 104$	$0.322 \pm 0.032 \pm 0.022$	$1.71 \pm 0.17 \pm 0.13$	9.4
$D^+\tau^-\bar{\nu}_\tau$	$177 \pm 31$	$986 \pm 35$	$0.469 \pm 0.084 \pm 0.053$	$1.01 \pm 0.18 \pm 0.12$	5.2
$D^{*+}\tau^-\bar{\nu}_\tau$	$245 \pm 27$	$3186 \pm 61$	$0.355 \pm 0.039 \pm 0.021$	$1.74 \pm 0.19 \pm 0.12$	10.4
$D\tau^-\bar{\nu}_\tau$	$489 \pm 63$	$2981 \pm 65$	$0.440 \pm 0.058 \pm 0.042$	$1.02 \pm 0.13 \pm 0.11$	6.8
$D^*\tau^-\bar{\nu}_\tau$	$888 \pm 63$	$11953 \pm 122$	$0.332 \pm 0.024 \pm 0.018$	$1.76 \pm 0.13 \pm 0.12$	13.2

Comparison with SM calculation:

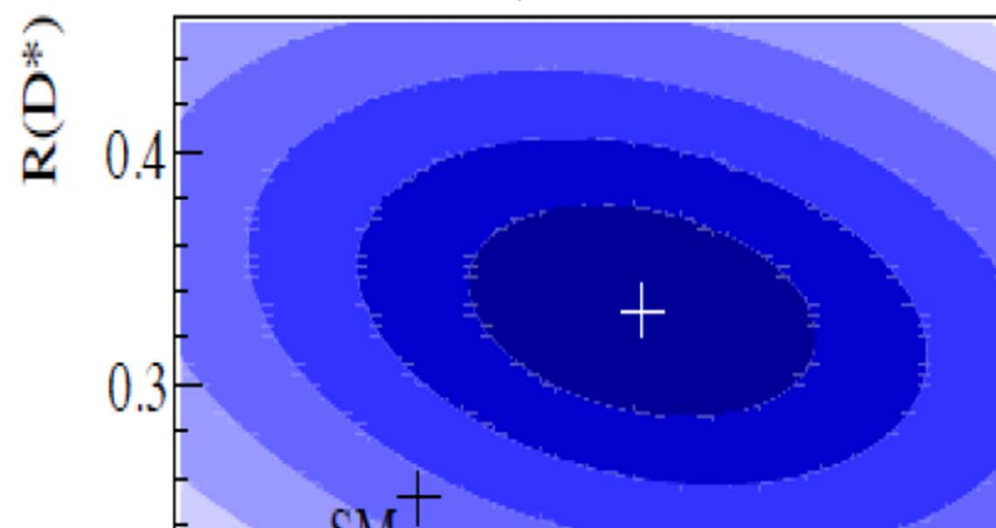
	$R(D)$	$R(D^*)$
BABAR	$0.440 \pm 0.071$	$0.332 \pm 0.029$
SM	$0.297 \pm 0.017$	$0.252 \pm 0.003$
Difference	<b><math>2.0 \sigma</math></b>	<b><math>2.7 \sigma</math></b>

The combination of the two measurements (0.27 correlation) yields  $\chi^2/\text{NDF}=14.6/2$ ,  
 $\text{Prob} = 6.9 \times 10^{-4}$

LATH

Combined 3.46

BABAR



Combination to 3.4 sigma is too aggressive and worrbersome

$$H_t^{2\text{HDM}} = H_t^{\text{SM}} \times \left( 1 - \frac{\tan^2 \beta}{m_{H^\pm}^2} \frac{q^2}{1 \mp m_c/m_b} \right)$$

- for  $D\tau\nu$   
+ for  $D^*\tau\nu$

This could enhance or decrease the ratios  $R(D^*)$  depending on  $\tan\beta/m_H$

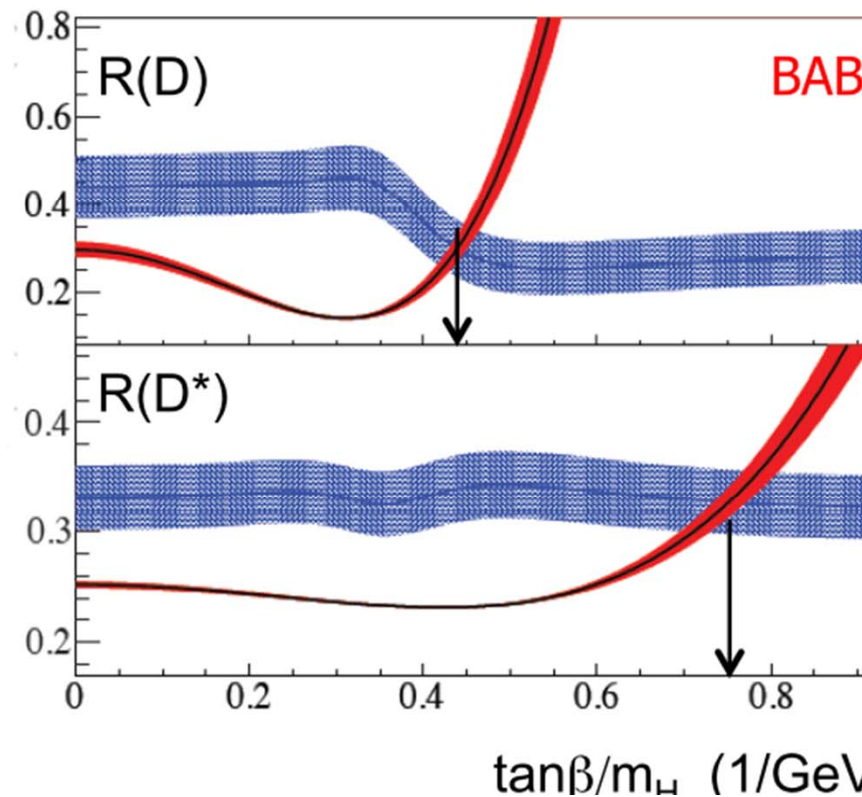
We estimate the effect of 2DHM, accounting for difference in efficiency, and its uncertainty

The data match 2DHM Type II at

$$\tan\beta/m_H = 0.44 \pm 0.02 \quad \text{for } R(D)$$

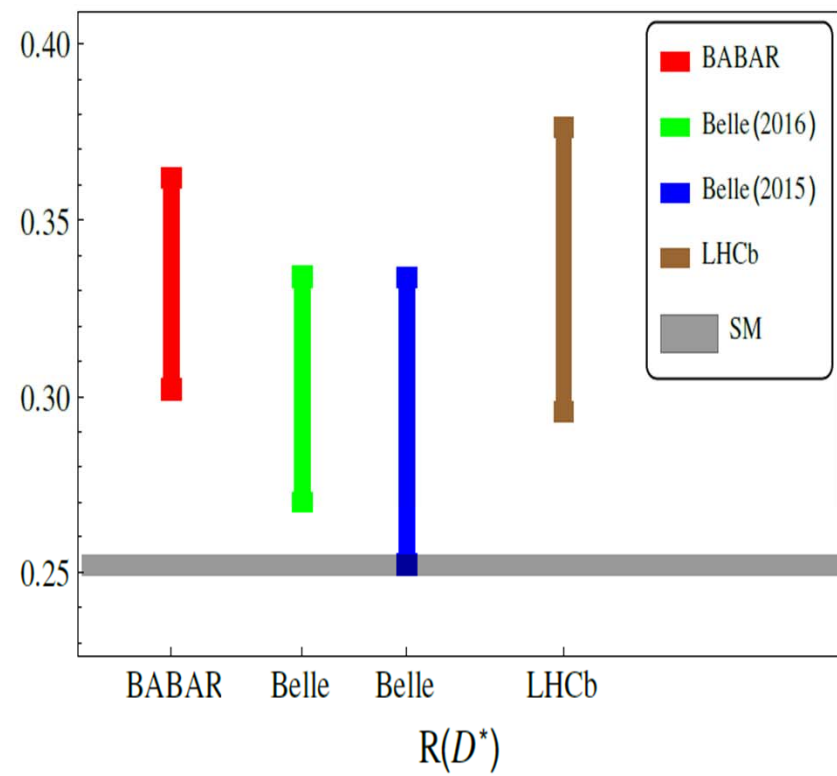
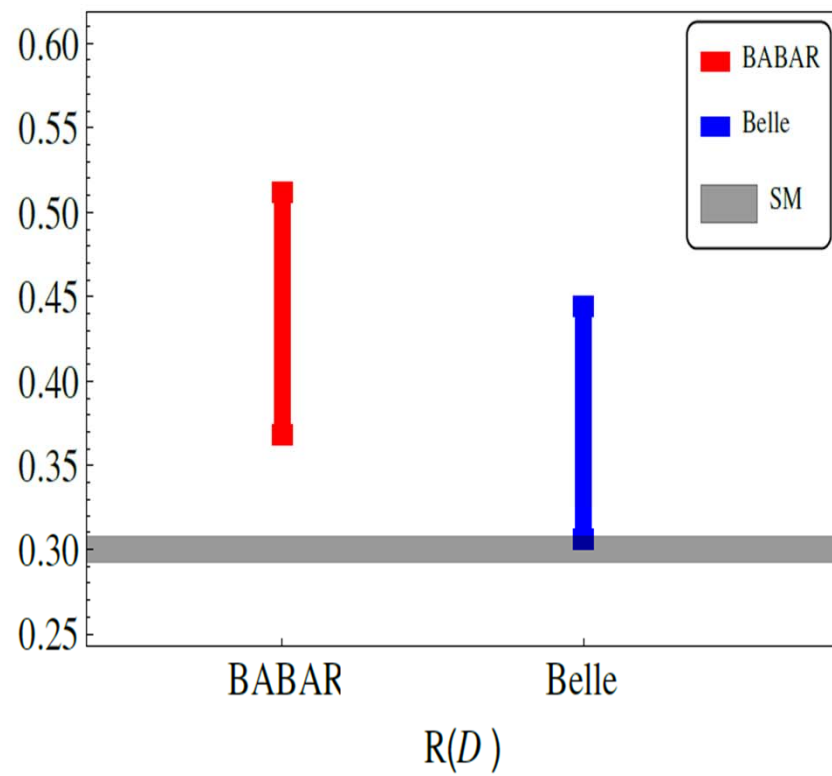
$$\tan\beta/m_H = 0.75 \pm 0.04 \quad \text{for } R(D^*)$$

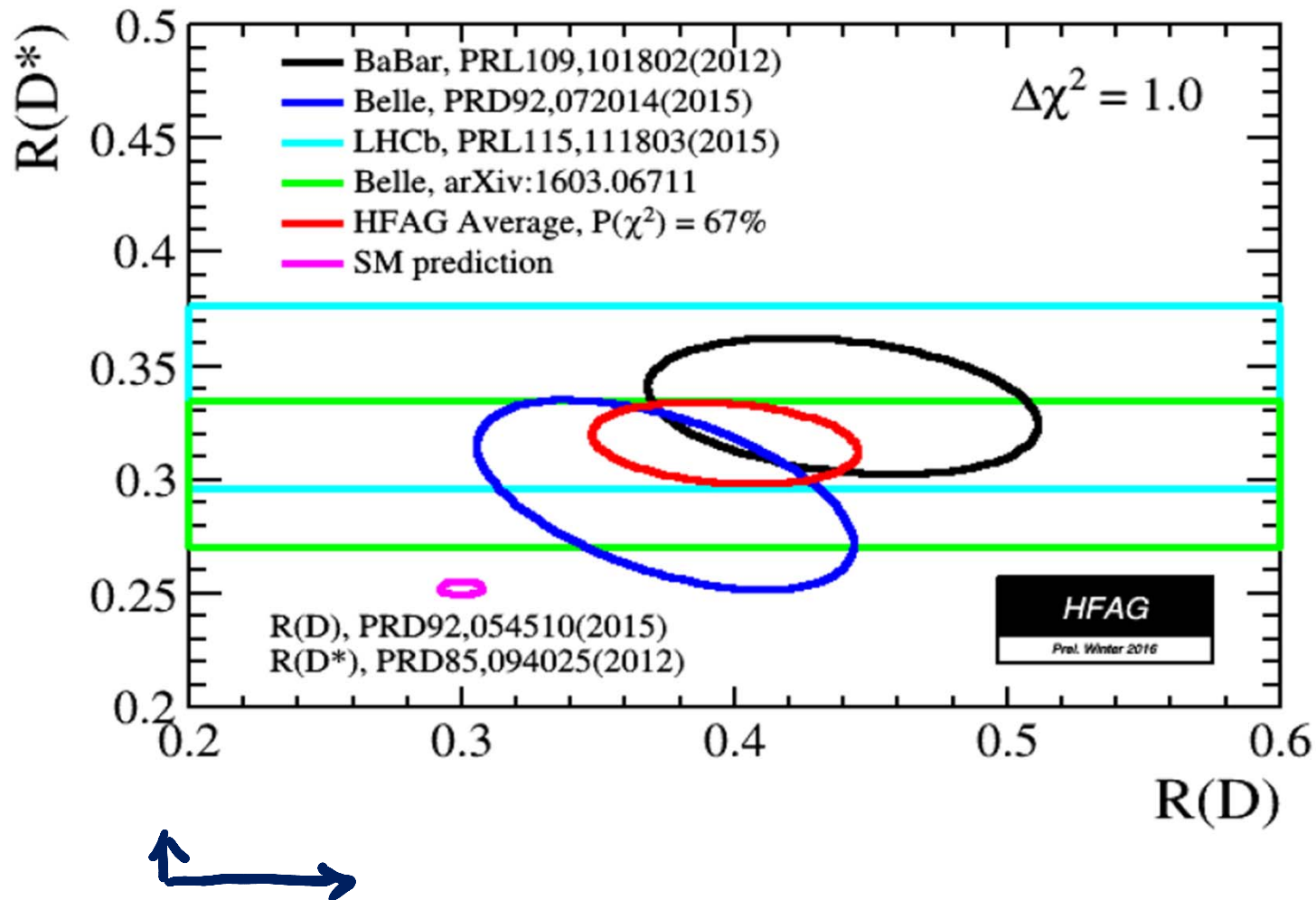
However, the combination of  $R(D)$  and  $R(D^*)$  excludes the Type II 2HDM in the full  $\tan\beta$ - $m_H$  parameter space with a probability of  $>99.8\%$ , provided  $M_H > 10\text{GeV}$  !



The conclusion that type II 2HDM  
is ruled out seems premature


# Enter Belle & LHCb





# My (subjective) summary Part I

- Theory is relatively clean; though 1 form factor does not get checked by light lepton as its contribution is mass suppressed...theory error
- A WEAK LINK: on  $D^*$  theory error look anomalously small....reason NO LATTICE CALCULATION SO FAR
- It seems really impressive that all three experiments are lending support that this may be a problem.....but significance of each is quite low . [Remember DO-dir CP...]
- **Tau's have always been extremely difficult historically.....many 1<sup>st</sup> results were problematic....this may share the same fate.**

A critique [SN+SP+AS  NPS  
arxiv:1605.07191; WHEPP XIV,  
IITK,India, Dec 2015]

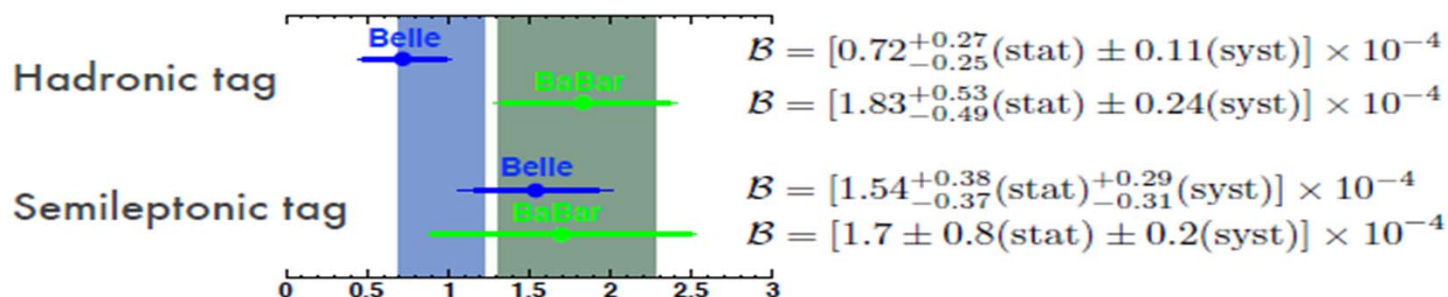
# My concerns

- Esp for  $RD^*$  theory errors look way too small
- Method of detecting tau is common to both  $RD$  AND  $RD^*$
- So  $RD$  and  $RD^*$  are not really independent observable
- Moreover, before  $RD^*$ , there in fact for about several years was a  $BR [B \rightarrow \tau \nu]$  anomaly originally seen in 2 ways in BABAR and also confirmed in both ways by BELLE
- However, around 07/04/12, BELLE made an important announcement

[Y KWON]

2012 Results

# $B^+ \rightarrow \tau^+ \nu_\tau$ Belle/BaBar comparison

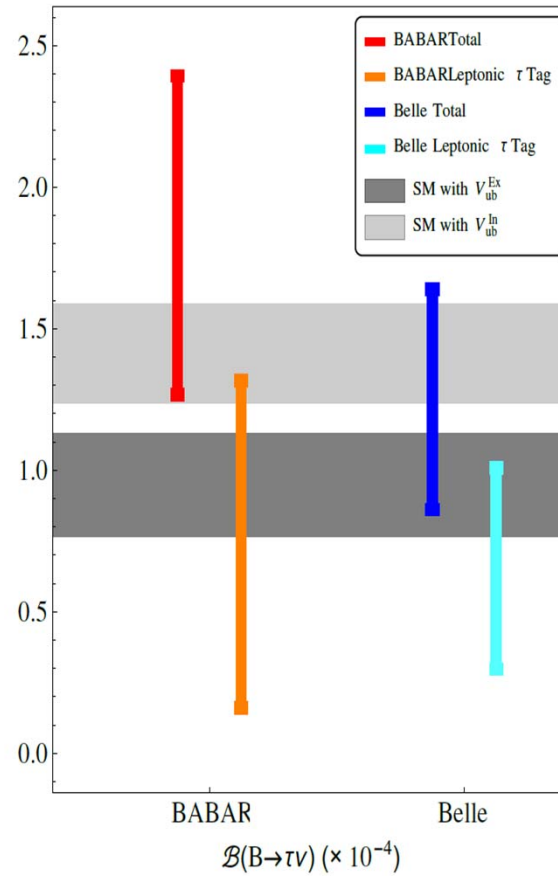


Belle combined:  $\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = (0.96 \pm 0.26) \times 10^{-4}$   
 BaBar combined:  $\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = (1.79 \pm 0.48) \times 10^{-4}$

- Belle vs. BaBar – consistent within  $\sim 1.5\sigma$
- The results are consistent with  $\mathcal{B}_{\text{SM}} = (1.10 \pm 0.30) \times 10^{-4}$ , which is based on

- \*  $f_B = (190 \pm 13) \text{ MeV}$  from HPQCD, PRD 80, 014503 (2009)
- \*  $|V_{ub}| = (4.15 \pm 0.49) \times 10^{-3}$  from PDG 2012 (via  $B \rightarrow X_u \ell \nu$ , incl.+excl.)

SM  $\begin{cases} V_{ub}^{inc} \rightarrow \\ V_{ub}^{exc} \rightarrow \end{cases}$



RECENT Belle UPDATE  
ON  
 $B \Rightarrow T \nu$

Within largish errors at this  
point pretty  
consistent with the SM and  
with each other

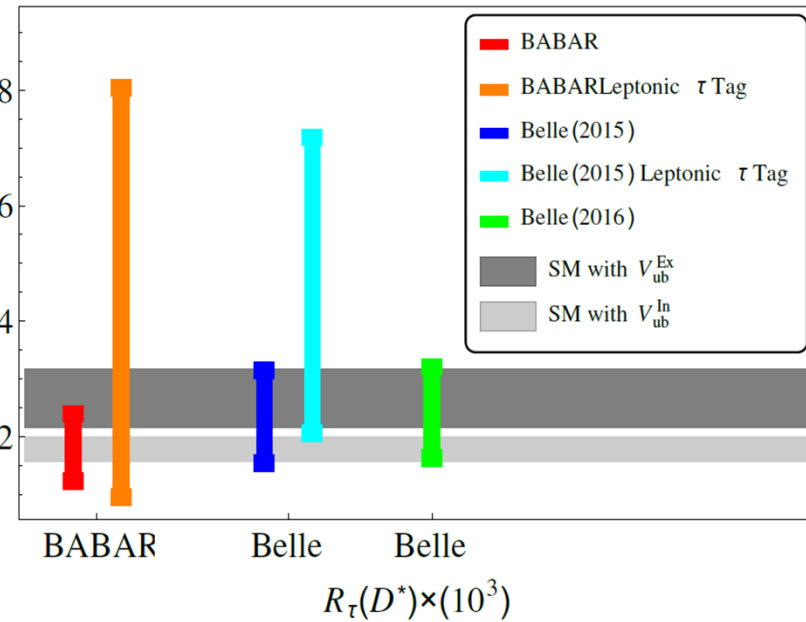
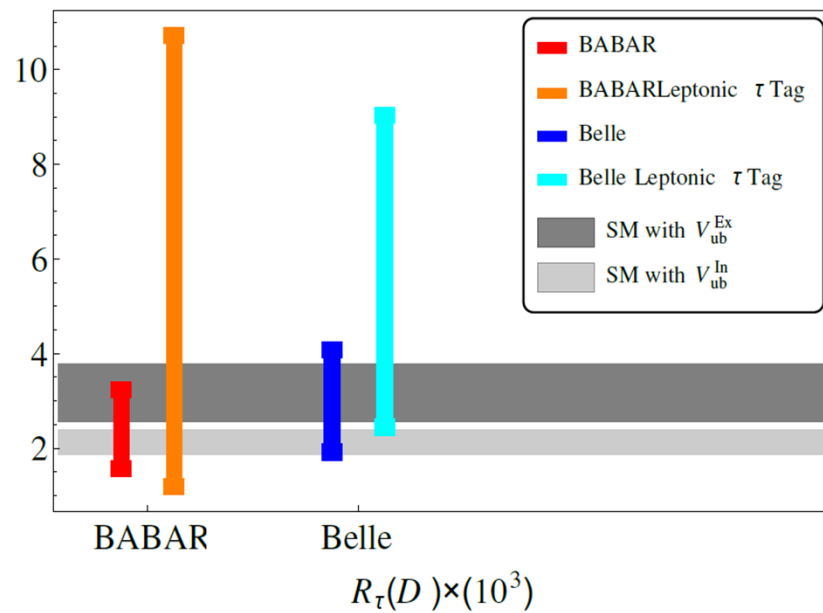
Bearing in mind all of the above, try a different route

- Introduce a new observable:

$$R_\tau(D^{(*)}) = \frac{\mathcal{R}(D^{(*)})}{\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau)}$$

- This is designed to be less sensitive to possible systematics afflicting tau detection though this depends on  $V_{ub}$  but...

SN, SP+AS [MIP]



REMARKABLY: NO deviation from the SM; all the few sigma problems are gone in here!

IN SHARP CONTRAST TO  $R_{D^{(*)}th}$

## Two conclusions

Hardly any indication of NP in here.

If there is any new physics then it seems largely cancelling away and an important consequence then is that Type II 2HDM may well be well and alive

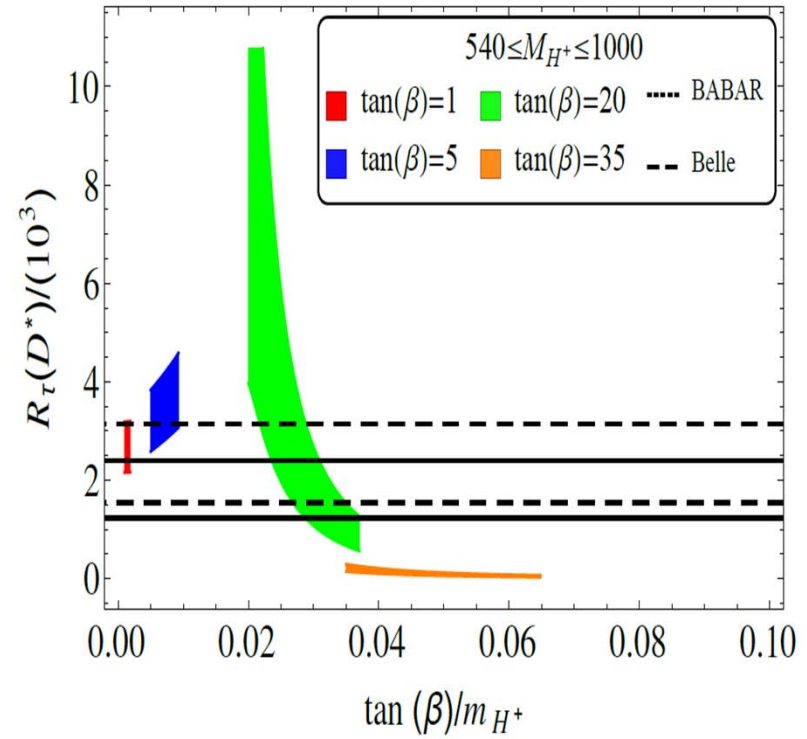
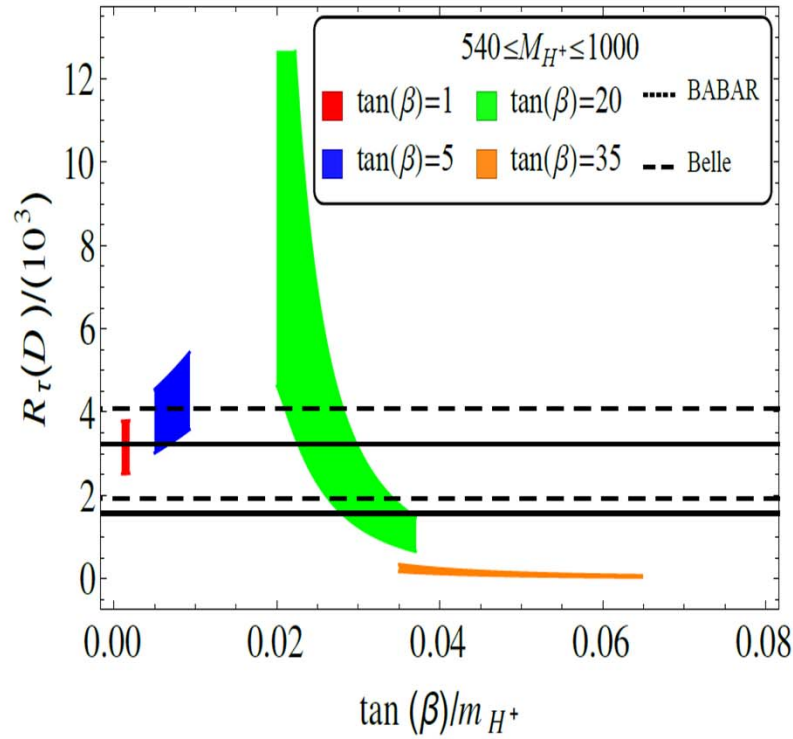


FIG. 7: Variations of  $R_\tau(D)$  (left) and  $R_\tau(D^*)$  (right) with the 2HDM-II parameter  $r = \tan \beta / m_{H^+}$  for different values of  $\tan \beta$ . The  $1\sigma$  experimental ranges are shown by the dotted (BABAR) and dashed (Belle) lines.

NPS

$\tan \beta \gtrsim 30$  Ruled out

# Take home

- **The best test for NP is  $B \Rightarrow \tau \nu$  [even though RD and RD\* do also give impt info]**
- **Interpreted in terms of NP, type II-2HDM with  $\tan \beta$  less than about 30 is a viable candidate [good news for SUSY]**
- **Belle-II with 20-40 X stat should be able to decisively test the SM with  $B \Rightarrow \tau \nu$ ; improved lattice calculations in the next 3 years should also be extremely useful for this purpose.**
- **Improved determination of  $B \Rightarrow \tau \nu$  also has important repercussion for UT ...see Lunghi + AS, PRL' 10**

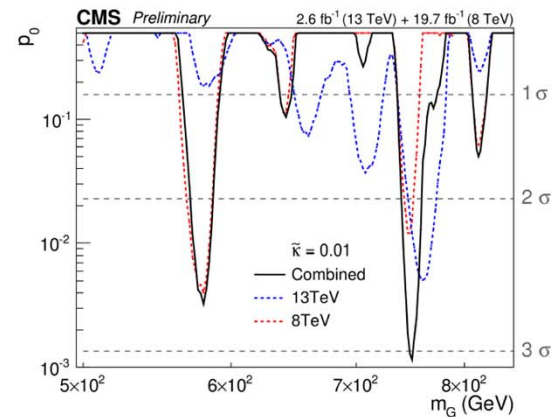
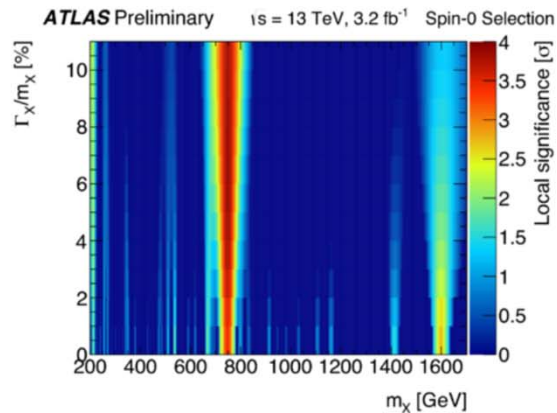
# Possibly exciting news from the collider front

## Diphoton anomaly @ 750 GeV

# local significance @ 750 GeV

ATLAS

CMS



Chances of sighting BSM seem high!

A highly plausible explanation of  
750 GeV and or flavor anomalies is  
XHM---typell-2HDM is a nice  
prototype =>FV & BSM-CP phase(s)

# BSM-CPV: promising menu; few prime examples

- In the SM CPA in  $t$ -decays are zilch due GIM
- BSM-CPV phase(s) in 2HDM....Neutral and or charged H exchanges
- $t \bar{t}$ .....very sensitive to phase(s) from neutral higgs
- $t\bar{t}h$ ..... very sensitive to neutral higgs
- $t \Rightarrow b \tau \nu$  [also  $b \Rightarrow c \tau \nu$ ] .....;  $\tau$  trans. Pol.very sensitive to charged higgs phase
- TDCPV  $B \Rightarrow \gamma K \pi$  ....
- $K : \epsilon_s' ; K_L \Rightarrow \pi^0 \nu \nu ; K_{UT}$

# Possibilities of large tree-level CP violation

- $e^+ e^- \Rightarrow t \bar{t} h (Z)$
- Simple example used 2HDM
- See Atwood et al PRD 96 ; PRD 98
- See also Gunion, Grzadkowski, He PRL96

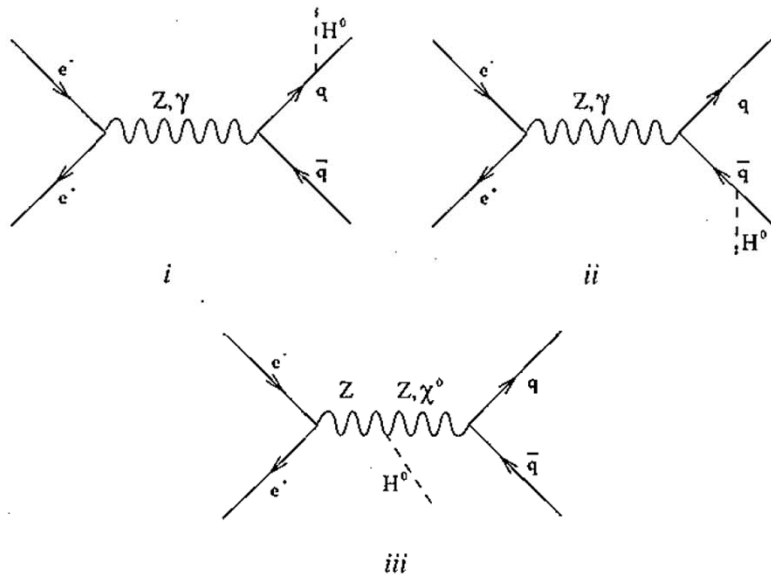


FIG. 1. Tree-level Feynman diagrams contributing to  $e^+e^- \rightarrow t\bar{t}H^0$  within the two Higgs doublet model.

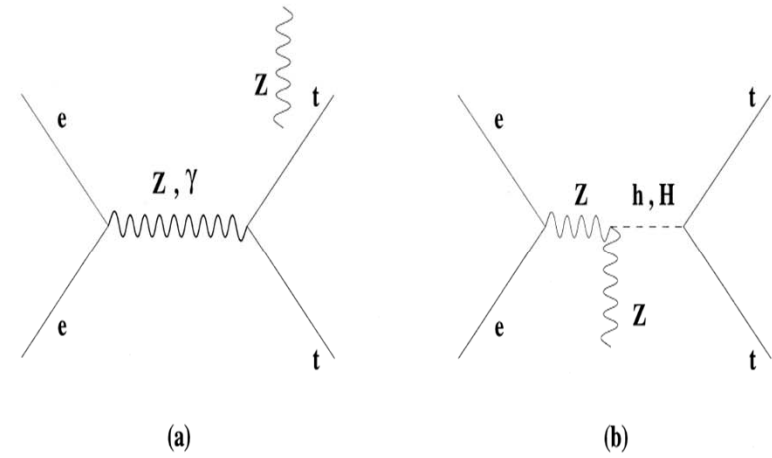


Fig. 1. Tree-level Feynman diagrams contributing to  $e^+e^- \rightarrow t\bar{t}Z$  in a two Higgs doublet model. Diagram *a* represents 8 diagrams in which either  $Z$  or  $\gamma$  are exchanged in the  $s$ -channel and the outgoing  $Z$  is emitted from  $e^+, e^-, t$  or  $\bar{t}$ .

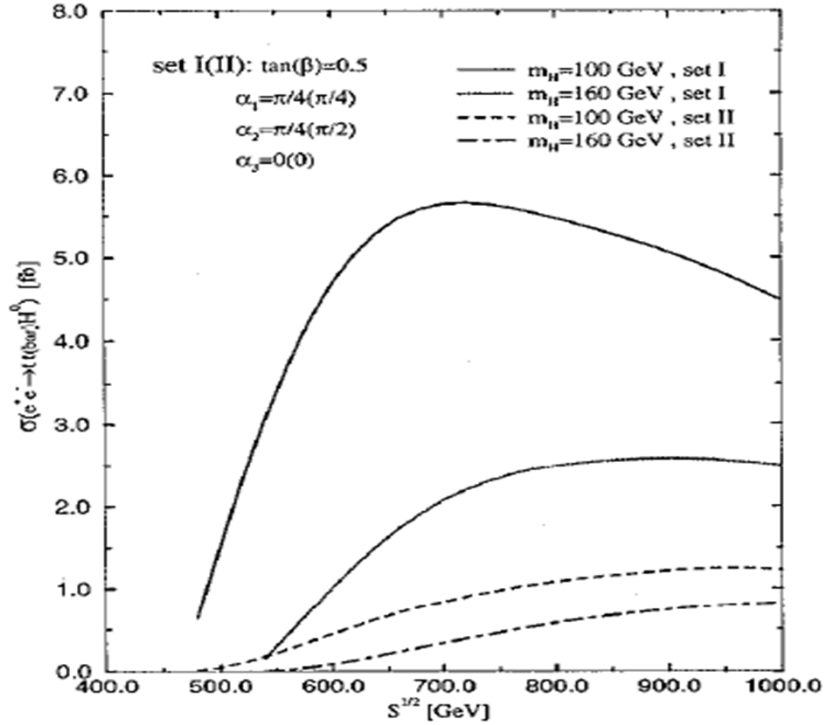


FIG. 2. The cross section (in [fb]) for the reaction  $e^+e^- \rightarrow t\bar{t}H^0$ , for sets I and II of the parameters  $\alpha_t$ ,  $b_t$ , and  $c$  and for  $m_{H^0} = 100$  and 160 GeV assuming unpolarized electron and positron beams.

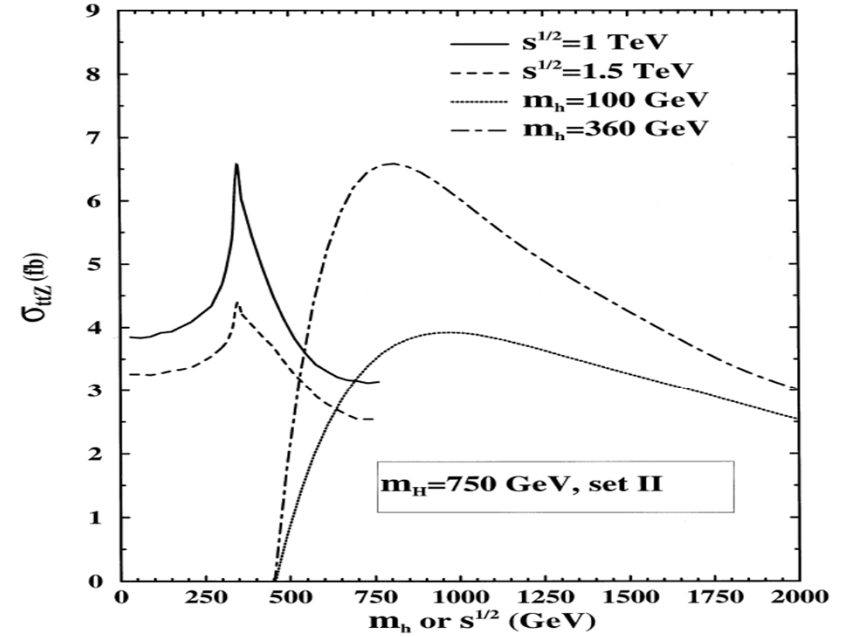


Fig. 2. The cross section (in [fb]) for the reaction  $e^+e^- \rightarrow \tau\bar{\tau}Z$ , assuming unpolarized electron and positron beams, for Model II with set II and as a function of  $m_h$  (solid and dashed lines) and  $\sqrt{s}$  (dotted and dotted-dashed lines). Set II means  $\{\tan\beta, \alpha_1, \alpha_2, \alpha_3\} \equiv \{0.3, \pi/2, \pi/4, 0\}$ .

#### IV. *CP*-VIOLATING OBSERVABLES AND NUMERICAL RESULTS

Of course, at tree level there are no absorptive phases. Thus, the *CP*-violating term  $\Sigma_-^0$  can probe only *CP* asymmetries of the  $T_N$ -odd type. This leads us to consider the following *CP*-odd,  $T_N$ -odd, triple correlation product

$$O = \vec{p}_- \cdot (\vec{p}_t \times \vec{p}_{\bar{t}}) / s^{3/2} . \quad (22)$$

*top decays analyze its spin.*

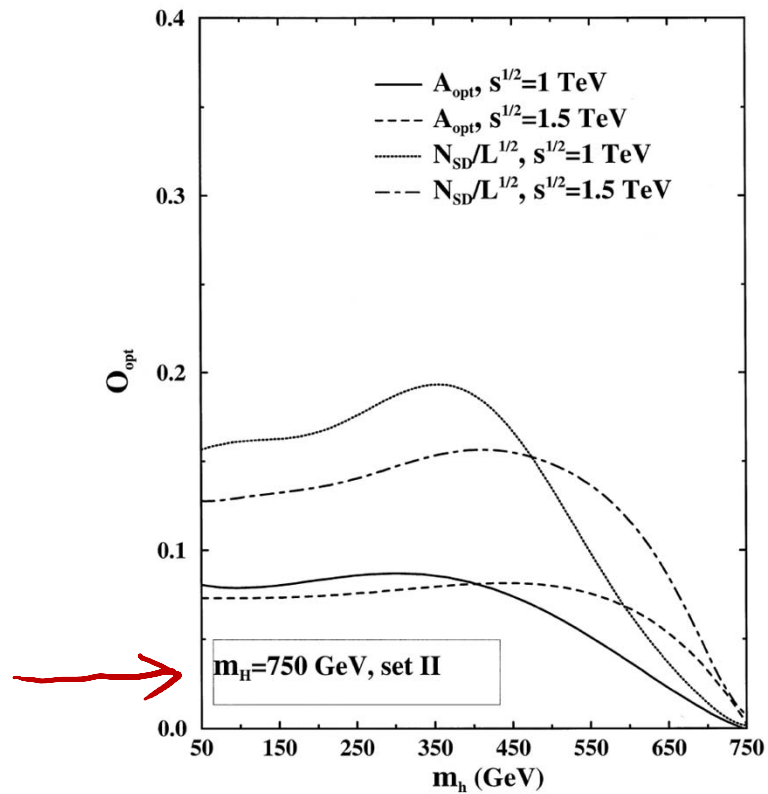


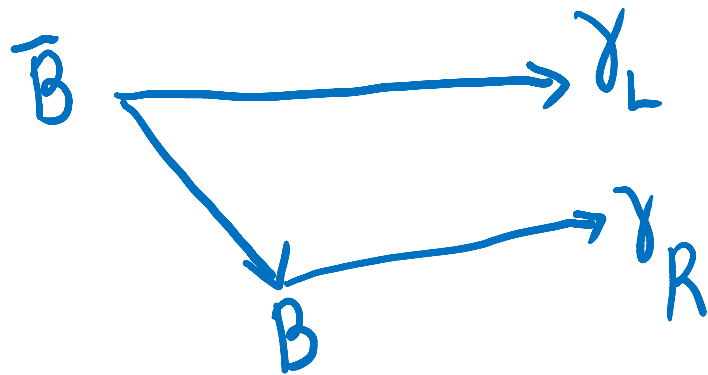
Fig. 3. The asymmetry,  $A_{\text{opt}}$ , and scaled statistical significance,  $N_{\text{SD}}/\sqrt{L}$ , for the optimal observable  $O_{\text{opt}}$  as a function of the light Higgs mass  $m_h$ , for  $\sqrt{s} = 1 \text{ TeV}$  and  $1.5 \text{ TeV}$ . See also caption to Fig. 2.

ASYMMETRY  
 $O(10\%)$

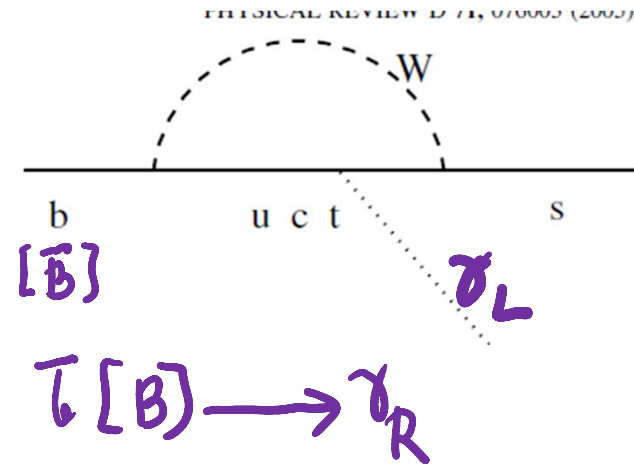
# Mixing induced (i.e. TD) CPV in radiative B-decays [Atwood, Gronau, AS, PRL'97]

$$H_{\text{eff}} = -\sqrt{8} G_F \frac{e m_b}{16\pi^2} F_{\mu\nu} \left[ \frac{1}{2} F_L^q \bar{q} \sigma^{\mu\nu} (1 + \gamma_5) b + \frac{1}{2} F_R^q \bar{q} \sigma^{\mu\nu} (1 - \gamma_5) b \right].$$

*SM*  
 $F_R^q / F_L^q \sim \frac{m_q}{m_b}$



*NO interference  
in SM*



*QCD Complications emphasized in :*

B. Grinstein, Y. Grossman, Z. Ligeti, and D. Pirjol, Phys. Rev. D **71**, 011504 (2005).

Dalitz study of mixing induced CPV in  $B, B_s \Rightarrow P1 P2 \gamma$ ; **a data driven *clean* way to constrain NP** [Atwood, Gershan, Hazumi & AS PRD'05]

- **Note P1, P2 are CP-eigenstates e.g  $K_S^0 \pi^0 \gamma$  ;  $\pi^+ \pi^- \gamma$  .....**

TABLE I. Final states which can be used to probe  $b \rightarrow s \gamma$  and  $b \rightarrow d \gamma$  transitions in  $B_d$  and  $B_s$  decays. This list is not exhaustive; in particular, other neutral (pseudo-)scalar particles ( $\eta$ ,  $\eta'$ ,  $f_0$ ) may be used in the place of  $\pi^0$ .

	$K_S \pi^0 \gamma$	$K_S K_S \gamma$	$\pi^+ \pi^- \gamma$	$K^+ K^- \gamma$	$K_S K_L \gamma$
$B_d / \bar{B}_d$	$b \rightarrow s \gamma$	$b \rightarrow d \gamma$	$b \rightarrow d \gamma$	$b \rightarrow d \gamma$	$b \rightarrow d \gamma$
$B_s / \bar{B}_s$	$b \rightarrow d \gamma$	$b \rightarrow s \gamma$	$b \rightarrow s \gamma$	$b \rightarrow s \gamma$	$b \rightarrow s \gamma$

**LOOK FORWARD to BELLE-2 & possibly LHCb Cashing in on these.**

$\epsilon'$ : a possible gem in  
search of new phenomena

*c also NORMAN CHRIST*

## $\epsilon' / \epsilon$ : Direct CPV

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L \rightarrow \pi^+ \pi^-)}{A(K_S \rightarrow \pi^+ \pi^-)}$$
$$\eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L \rightarrow \pi^0 \pi^0)}{A(K_S \rightarrow \pi^0 \pi^0)}$$

$$\eta_{+-} = \epsilon + \epsilon', \quad \eta_{00} = \epsilon - 2\epsilon'$$

$$\text{Re}\left(\frac{\epsilon'}{\epsilon}\right) = \frac{\omega}{\sqrt{2}|\epsilon|} \left[ \frac{\text{Im}(A_2)}{\text{Re}(A_2)} - \frac{\text{Im}(A_0)}{\text{Re}(A_0)} \right]; \quad \omega \approx \frac{\text{Re}A_2}{\text{Re}A_0}$$

BNL '64  
CRONIN +  
FITCH N.P.

DIRECT ~~CP~~

$$|\epsilon| = 2.228(11) \times 10^{-3},$$

$$\text{Re}(\epsilon'/\epsilon) = 1.65(26) \times 10^{-3}.$$

$$\epsilon' \sim 0(10^{-6})!$$

Indirect CP

## Its presumed importance:

- lies in its very small size => Perhaps new phenomena has a better chance of showing up
- Exceedingly important monitor of flavor –alignment
- Simple naturalness arguments strongly suggest  $\epsilon$ s' very sensitive to BSM – CP odd phases
- In many ways  $\epsilon$ s' is rather analogous to nedm.....both being very sensitive to BSM phases
- Understanding  $\epsilon$ s', nedm is extremely important for learning how naturalness really works in nature

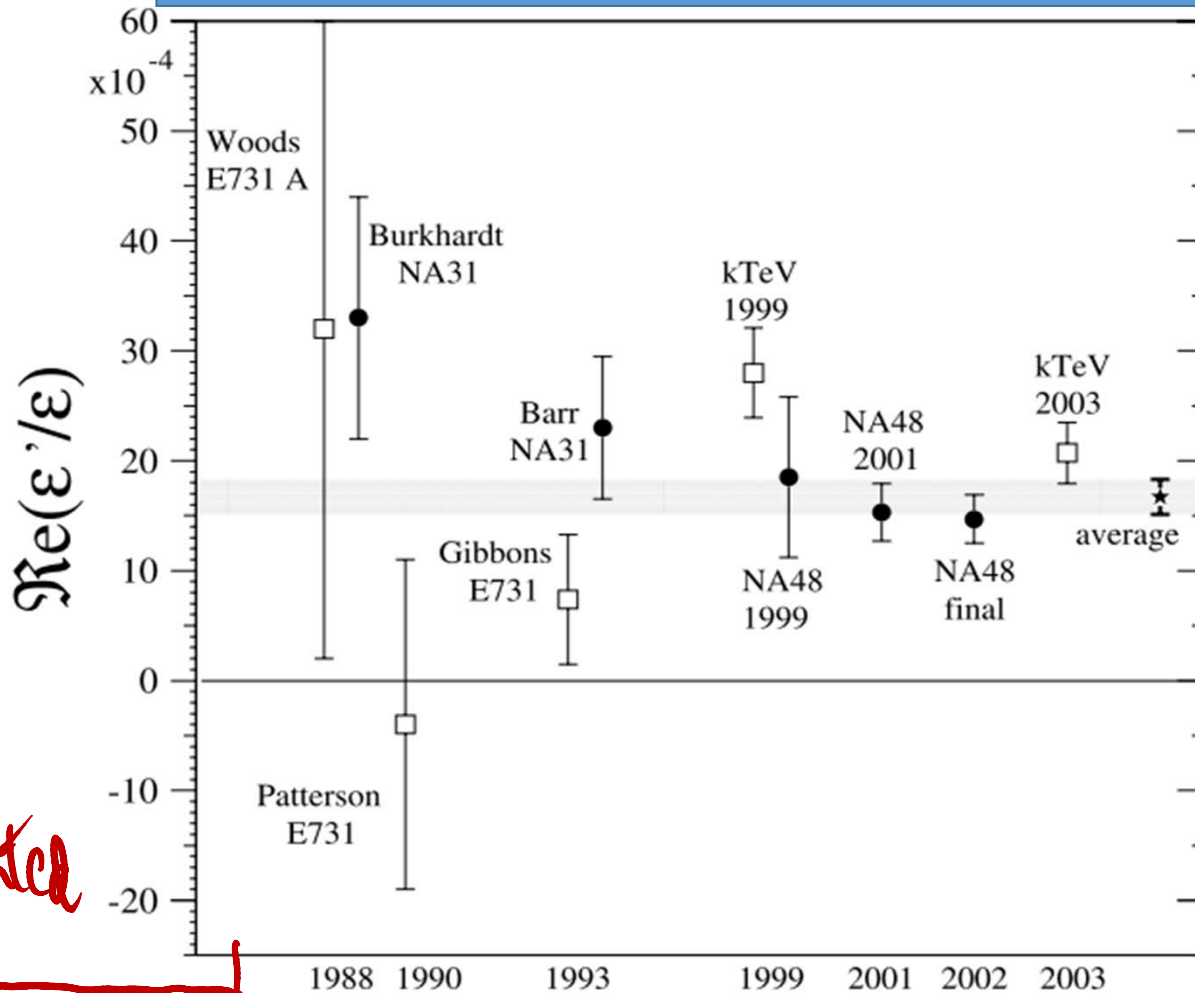
## A.S. in Proceedings of Lattice '85 (FSU)..1<sup>st</sup> Lattice meeting ever attended

The matrix elements of some penguin operators control in the standard model another CP violation parameter, namely  $\epsilon'/\epsilon$ .<sup>6,8)</sup> Indeed efforts are now underway for an improved measurement of this important parameter.<sup>10)</sup> In the absence of a reliable calculation for these parameters, the experimental measurements, often achieved at tremendous effort, cannot be used effectively for constraining the theory. It is therefore clearly important to see how far one can go with MC techniques in alleviating this old but very difficult

With C. Bernard  
[UCLA]

L May 2016; V-TECH; soni[BNL]

A monumental experimental achievement!



Konrad  
kleinknecht  
"Uncovering CPV"

$16.6(2.3) \times 10^{-4}$   
PDG 2014

LATTICE  
WORK started

# MOTHER of all (lattice) calculations to date: A Personal Perspective

- **As mentioned, calculation  $K \Rightarrow \pi \pi$  &  $\epsilon \pi'$  were the reasons I went into lattice  $\sim 1/3$  of a century ago**
- **9 PhD thesis: Terry Draper (UCLA), George Hockney(UCLA), Cristian Calin (Columbia=CU), Jack Laiho(Princeton), Sam Li(CU), Matthew Lightman(CU), Elaine Goode(Southampton), Qi Liu(CU), Daiqian Zhang(CU)**

I. Wilson Fermions with Bernard ~'82 See also Martinelli et al [WF] Giusti et al [WF] Sharpe et al [Stag F]	Lattice $\chi$ S is a pre-requisite for this physics Off-shoot B-physics important observables identified & studied=> evolved into UT		
II (a) DWF with Blum ~ '95 II(b) DWF with RBC[with Blum, Christ and Mawhinney became "flagship" project of RBC] ~'97.	LO $\chi$ PT; Quenched approx.[QA] Same QA is disastrous for this physics [Golterman-Pallante] pathologies; NPR of full $\Delta S=1$ accomplished for the 1 <sup>st</sup> time used since then.	CRAY @ NERSC  QCDSP ~ 1 TF	
III. DWF with full QCD RBC, ~ '02	Used LO $\chi$ PT + full QCD Large chiral corrections	QCDSP ~ 1TF	
IV. DWF with full QCD RBC + UKQCD, ~ '06	<b>Direct <math>K \Rightarrow \pi\pi</math>, [Lellouch-Luscher method] @ threshold</b>	QCDOC ~ 10 TF	
V. DWF with full QCD, RBC + UKQCD ~ '11	Direct $K \Rightarrow \pi\pi$ , [Lellouch-Luscher method] ; physical kinematics	BG/Q ~ 100TF@BNL; RBRC;ANL; Edinburgh	
Vi. Same ~now	Same	new hardware  ~1.5PF;NERSC;ANL;UK	

HUGE # of  
OBSTACLES HAS  
to be overcome

~2006  
↓

## Results for $\varepsilon'$

- Using  $\text{Re}(A_0)$  and  $\text{Re}(A_2)$  from experiment and our lattice values for  $\text{Im}(A_0)$  and  $\text{Im}(A_2)$  and the phase shifts,

and our lattice values for

$$\text{Re} \left( \frac{\varepsilon'}{\varepsilon} \right) = \text{Re} \left\{ \frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}\varepsilon} \left[ \frac{\text{Im}A_2}{\text{Re}A_2} - \frac{\text{Im}A_0}{\text{Re}A_0} \right] \right\}$$

$$= \frac{1.38(5.15)(4.43) \times 10^{-4}}{16.6(2.3) \times 10^{-4}}, \quad \begin{array}{l} \text{(this work)} \\ \text{(experiment)} \end{array}$$

RBC-UKQCD PRL'15  
EDITOR'S CHOICE

EWP

QCDP

PARTIAL  
Cancellation

Bearing in mind the largish errors in this first calculation, we interpret that our result are consistent with experiment at  $\sim 2\sigma$  level

$$\omega = \frac{\text{Re} A_2}{\text{Re} A_0} \sim 0.045$$

$\text{Re} A_0, A_2$  from expt  
band  $\omega$ .

Tree

$$Q_1 = (\bar{s}_\alpha d_\alpha)_L (\bar{u}_\beta u_\beta)_L,$$

$$Q_2 = (\bar{s}_\alpha d_\beta)_L (\bar{u}_\beta u_\alpha)_L,$$

$$Q_3 = (\bar{s}_\alpha d_\alpha)_L \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_L,$$

$$Q_4 = (\bar{s}_\alpha d_\beta)_L \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_L,$$

$$Q_5 = (\bar{s}_\alpha d_\alpha)_L \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_R,$$

$$Q_6 = (\bar{s}_\alpha d_\beta)_L \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_R,$$

$$Q_7 = \frac{3}{2} (\bar{s}_\alpha d_\alpha)_L \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_R,$$

$$Q_8 = \frac{3}{2} (\bar{s}_\alpha d_\beta)_L \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_R,$$

$$Q_9 = \frac{3}{2} (\bar{s}_\alpha d_\alpha)_L \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_L,$$

$$Q_{10} = \frac{3}{2} (\bar{s}_\alpha d_\beta)_L \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_L,$$

EWP

~~I=2~~

QCDP

$I=0 \Rightarrow$

$\rightarrow 0$   
 $m_q \rightarrow 0$

$\rightarrow \text{const}$   
 $m \rightarrow 0$

$\frac{S \bar{M}_d}{E_9}$   
QCDP

$\frac{S \bar{M}_d}{E_9}$   
 $\frac{S \bar{M}_d}{E_9}$

EWP

## More demands on the calculation

- ~ The 1995 discovery of the huge top mass has accentuated the cancellation of  $I=0$  and  $I=2$  contributions to  $\epsilon_s$  significantly, putting additional demands on the calculation but also enhancing the potential for discovery of new physics

$$c_8 \propto m_t^2 / M_W^2$$



**IF YOU BUILD IT THEY WILL COME**

If there is new physics around below  $\sim 5$  TeV, there is an excellent chance that eps' will find it!

[of course requires accurate theory calculation... RBC-UKQCD plans for X4 in stat and appreciable improvements in systematic in  $\sim 2$  years ]

Kaon Unitarity Traingle (KUT)  
Lehner,Lunghi +AS [LLS],  
arxiv:1508.01801

# A dream for some

Blucher, Winstein and  
Yamanaka '09

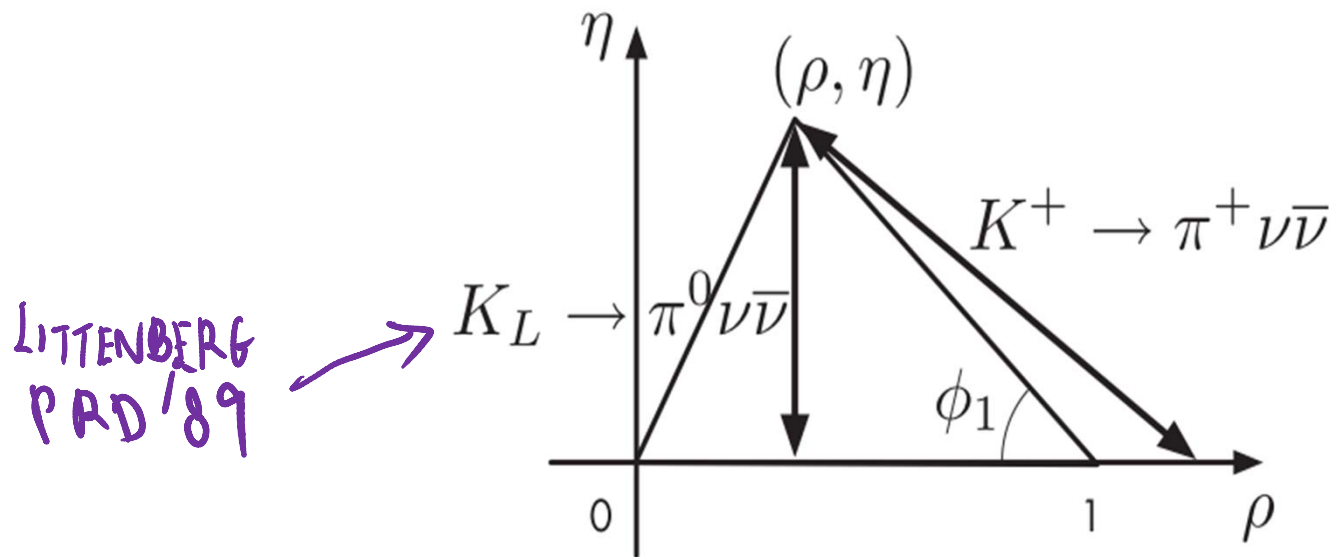
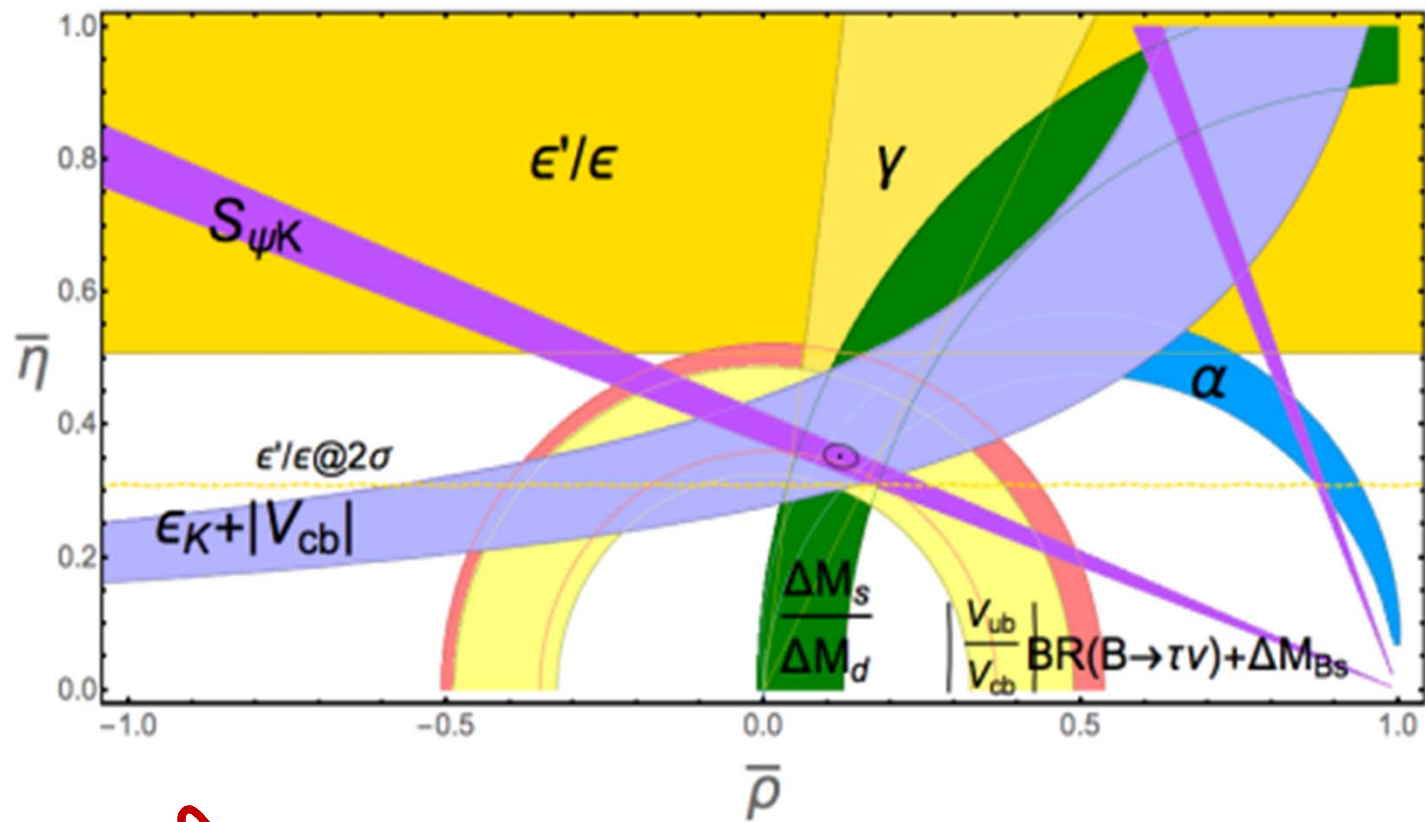


Fig. 14. Unitarity triangle.

A Faster way in the offing?

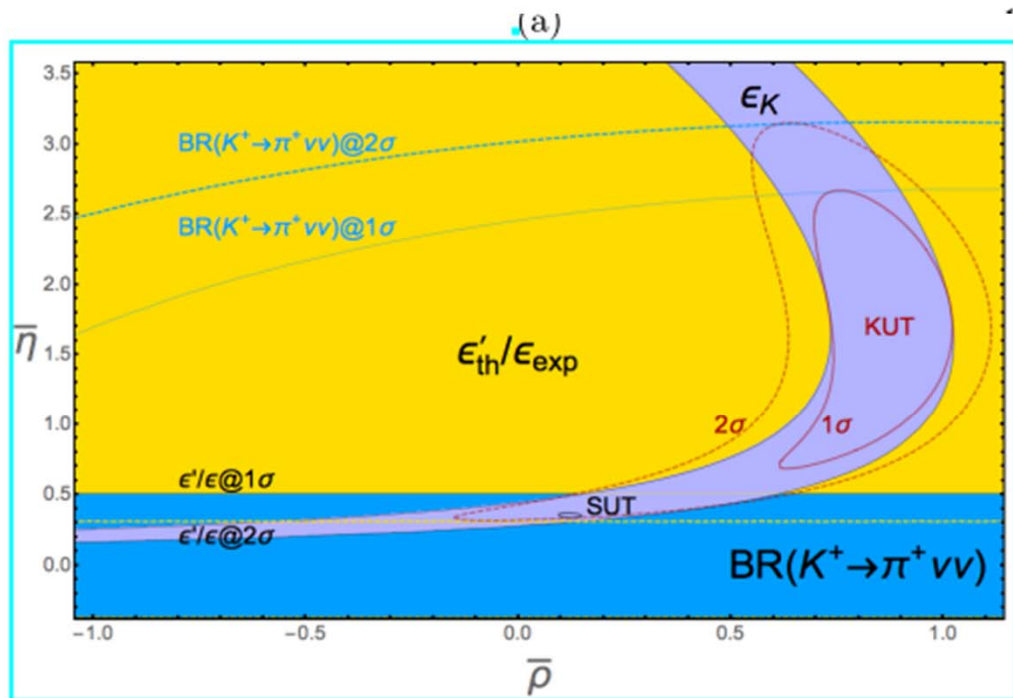
Lattice  $\epsilon'/\epsilon$  & SUT  $\equiv$  B-UT



LLS

Sketch of an emerging K-UT

$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = \begin{cases} (8.64 \pm 0.60) \times 10^{-11} & \text{SM} \\ (17.3^{+11.5}_{-10.5}) \times 10^{-11} & \text{E949} \end{cases} \quad \text{BNL}$$

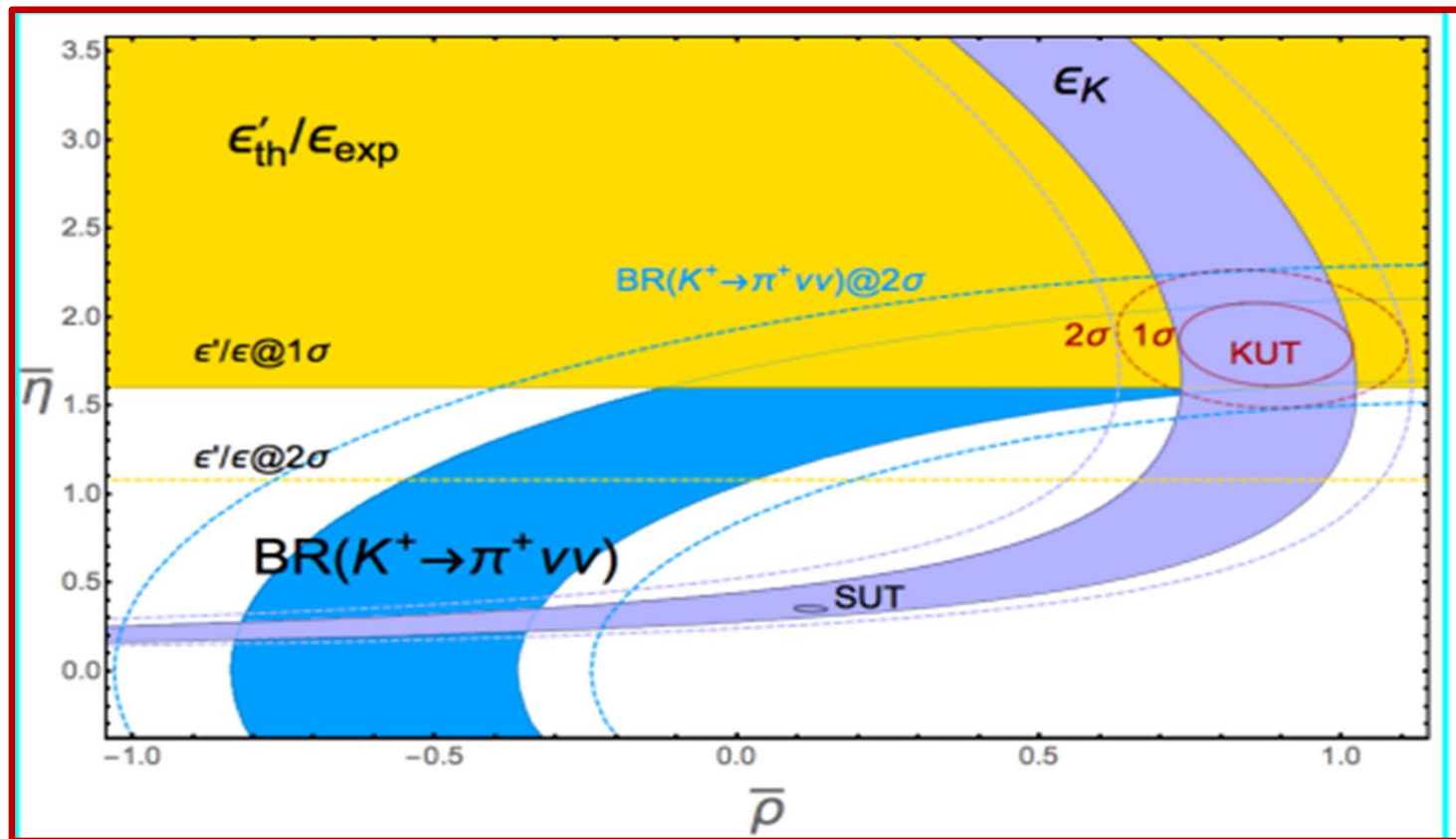


$$\text{Re}\left(\frac{\epsilon'}{\epsilon}\right)_K = \begin{cases} (16.7 \pm 1.6) \times 10^{-4} \\ (1.36 \pm 5.21_{\text{stat}} \pm 4.49_{\text{syst}}) \times 10^{-4} \end{cases} \quad \text{PDG 2015}$$

← RBC+UKQCD '15

LHS '15

POSSIBLE KUT CIRCA 2020: DUE NA62+RBC-UKQCD



NO unique  
 $\rho, \eta$

Assumed: NA62, 100  
events with  $\sim 7\%$  error  
RBC-UKQCD,  
 $\delta(\text{Im}A_0) \sim 18\%$   
[current  $\sim 60\%$ ]

LLS

# Summary & Outlook [p1 of 2]

- SM-CKM works to  $\sim 15\%$   $\Rightarrow$  we need to work harder to unearth new phases as it is extremely difficult to understand why they are not there...
  - Current hints of BSM in experiments clearly need scrutiny & clarification
  - $V_{ub}$  diff bet exc and inc unlikely due NP
  - $R_{D^{(*)}}$  indications of NP are very interesting, careful understanding of experimental systematics of tau detection are essential.
  - However,  $R_{D^{(*)}}$  anomaly is NOT different than NP in the simpler  $B \Rightarrow \tau \nu$ ; better measurement of  $Br$  of  $B \Rightarrow \tau \nu$  is extremely important.
  - Be that as it may it does not appear that type II, 2 HDM is ruled out....good news for SUSY
  - Reg eps', if there is new physics around, then its very likely that eps' will find it.
- We believe we'll be able to reduce theory errors significantly in 3-5 years;
- . ....Improvement in Expt may also be worth thinking.

# Summary & Outlook [p.2]

- LHC  $\sim 750$  GeV diphoton signal is perhaps the strongest indication of new physics in a long long time
- If true very likely extended Higgs sector will be needed [possibly along with other things..]
- Strong renewed motivation for searches for BSM-CP phase(s) in: TDCPV in radiative B,Bs decays,  $B \rightarrow M \tau \nu$ ,  $\epsilon_s'$ ,  $KL = \pi^0 \nu \nu$ ; in production of top pairs,  $t\bar{t}$ ; and in top decays are very worthwhile in near future experiments
- **HIGH HOPES FOR RUN II, [upgraded] LHCb & Belle-II!**

# EXTRAS

## II. TWO-HIGGS-DOUBLET MODEL WITH $CP$ VIOLATION IN THE NEUTRAL HIGGS SECTOR

The model we use as the source for  $CP$  violation is a THDM [5,6] namely, a nonminimal SM with the two complex Higgs fields  $\Phi_1$  and  $\Phi_2$ , in which  $CP$  violation arises from exchanges of neutral Higgs particles. Flavor-changing neutral currents (FCNC's) at the tree level, will appear in the theory if the two vacuum expectation values (VEV's) contribute to the quark mass matrices. To avoid FCNC's one can impose the discrete symmetry  $D$ ,

$$D : \Phi_2, u_{iR} \rightarrow -\Phi_2, -u_{iR} , \quad (1)$$

which gives the coupling scheme [10]

$$u_R \leftrightarrow \Phi_2, d_R, l_R \leftrightarrow \Phi_1 . \quad (2)$$

The invariant quark Yukawa interactions then read

$$\mathcal{L}_Y = -(\bar{u}_i, \bar{d}_i)_L \Gamma_u^{ij} \tilde{\Phi}_2 u_{jR} - (\bar{u}_i, \bar{d}_i)_L \Gamma_d^{ij} \Phi_1 d_{jR} + \text{H.c.} , \quad (3)$$



In general the Higgs potential can be written as:

$$V(\Phi_1, \Phi_2) = V_{\text{symm}}(\Phi_1, \Phi_2) + \delta V_{\text{soft}}(\Phi_1, \Phi_2) , \quad (4)$$

where  $V_{\text{symm}}(\Phi_1, \Phi_2)$  is the part of the potential which is symmetric under  $D$  while  $\delta V_{\text{soft}}$  breaks this discrete symmetry and depends on the phases of  $\Phi_1, \Phi_2$ . In particular,

$$\delta V_{\text{soft}} = \mu \Phi_1 \Phi_2^+ + k(\Phi_1^+ \Phi_2)^2 + \text{H.c.} \quad (5)$$

If  $\text{Im}(k/\mu^2) \neq 0$ , which is being assumed throughout this paper, then  $CP$  is no longer a symmetry of the Higgs potential. This will induce mixing between real and imaginary parts of the Higgs fields in the mass matrix of the Higgs boson, which means that the mass eigenstates do not have a definite  $CP$  property. Therefore, in the THDM  $CP$  violation may emanate from the neutral Higgs sector even when there is none in the charged Higgs sector. In general, the manifestation of such  $CP$  violation is that the neutral Higgs boson mass eigenstates couple to fermions with both scalar and pseudoscalar couplings.

For  $e^+e^- \rightarrow t\bar{t}H^0$  the following interaction terms (which appear in a THDM) in  $\mathcal{L}$  are required [5]:

$$\begin{aligned} \mathcal{L}_{H_j^0} = & H_j^0 \bar{f}(a_{fj} + ib_{fj}\gamma_5)f + H_j^0 c_j g_{\mu\nu} Z^\mu Z^\nu \\ & + \frac{c_j}{2M_Z} [\chi^0(\partial_\mu H_j^0) - (\partial_\mu \chi^0)H_j^0] Z^\mu , \end{aligned} \quad (6)$$



$\Delta S=1$   $H_W$

W L & NLO

Buchalla, Buras, Lautenbacher  
RMP 196; Cinquini et al 95

$$H_W = \frac{G_F}{\sqrt{2}} V_{us}^* V_{ud} \sum_{i=1}^{10} [z_i(\mu) + \tau y_i(\mu)] Q_i(\mu).$$

$$\tau = -V_{ts}^* V_{td} / V_{us}^* V_{ud}.$$