

Kaons on the Lattice

The XIII International Conference on
Heavy Quarks and Leptons

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RBC and UKQCD Collaborations

Kaons on the lattice

- Lattice QCD – 2016
 - Kaons from Euclidean space
 - Multi-particle final states
 - 2nd order long distance effects
- $K \rightarrow \pi \pi$ decay: $\Delta I = 1/2, \epsilon'$
- $K_L - K_S$ mass difference
- Long distance contribution to ϵ_K
- Long distance contribution to rare kaon decay: $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

Lattice QCD – 2016

- Physical quark masses (ChPT not needed)
- Chiral quarks (doubling problem solved)
- Large physical volumes: $(6 \text{ fm})^3$
- Small lattice spacing: $1/a = 2.4 \text{ GeV}$
 - $(\Lambda_{\text{QCD}} a)^2$ effects $< 1\%$ 😊
 - $(m_{\text{charm}} a)^2$ effects $\sim 20\%$ ☹️

RBC Collaboration

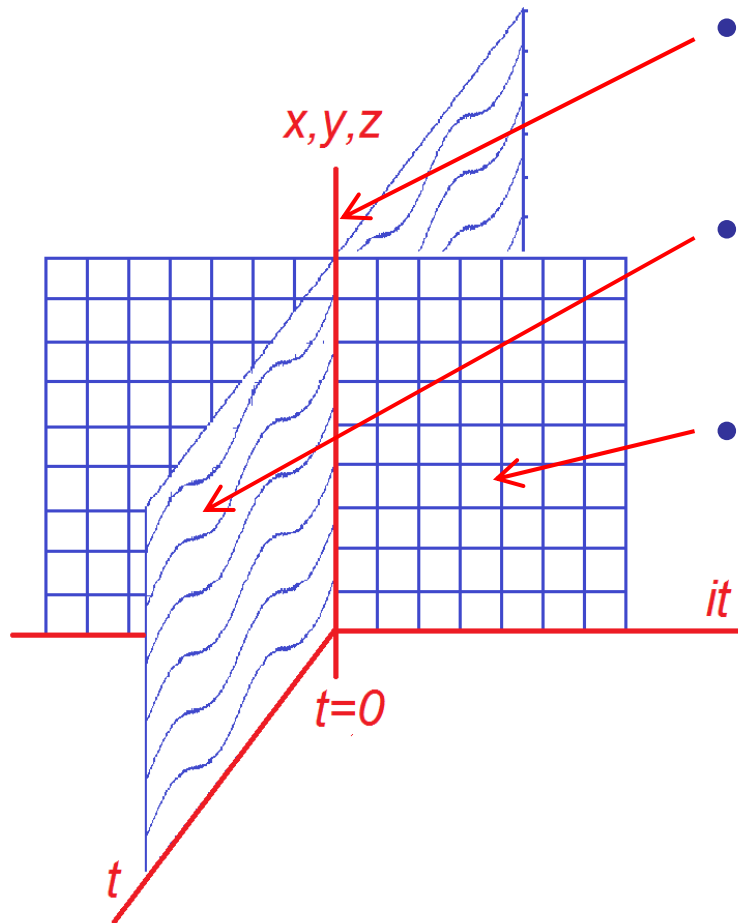
- BNL
 - Chulwoo Jung
 - Taku Izubuchi
 - Christoph Lehner
 - Meifeng Lin
 - Amarjit Soni
- RBRC
 - Mattia Bruno
 - Chris Kelly
 - Tomomi Ishikawa
 - Hiroshi Ohki
 - Shigemi Ohta (KEK)
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 - Ziyuan Bai
 - Xu Feng
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 - Luchang Jin
 - Robert Mawhinney
 - Greg McGlynn
 - David Murphy
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UKQCD Collaboration

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 - Jonathan Flynn
 - Vera Guelpers
 - Tadeusz Janowski
 - Andreas Juttner
 - Andrew Lawson
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 - Chris Sachrajda
 - Francesco Sanfilippo
 - Matthew Spraggs
 - Tobi Tsang
- Edinburgh
 - Peter Boyle
 - Julien Frison (KEK)
 - Nicolas Garron (Plymouth)
 - Ava Khamseh
 - Antonin Portelli
 - Oliver Witzel

Kaons from Euclidean Space

Kaons from Euclidean space



- Begin with standard Hilbert space QM at $t = 0$
- Use e^{-iHt} for physical time development.
- Lattice QCD requires the use of e^{-Ht} without the i
 - Low lying eigenvalues and eigenstates of H are easily accessible
 - More massive unstable states are increasingly difficult

Kaons from Euclidean space

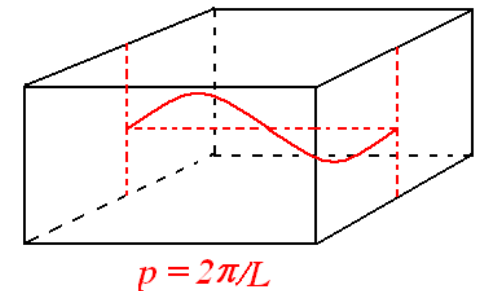
- Matrix elements between stable states are easy: [JHEP 1506 (2015) 164]

$$\langle \pi(p_\pi) | V_\mu(\vec{0}) | K(p_K) \rangle = f_+^{K\pi}(q^2)(p_K - p_\pi)_\mu + f_-^{K\pi}(q^2)(p_K + p_\pi)_\mu$$

$$f_+^{K\pi}(0) = 0.9685(34)(14) \quad |V_{us}| = 0.2233(5)(9)$$

- Final two-pion state more difficult:

- e^{-Ht} projects onto the ground state
- can obtain excited states by using finite volume quantization.
- Lellouch-Lüscher finite volume correction.



- Final three-pion state is hard [Hansen and Sharpe]

Kaons from Euclidean space

- Long-distance parts of second-order processes: ΔM_K , ε_K , and $K^+ \rightarrow \pi^+ \nu \nu$ are possible
- Exploit connection between time-dependent and time-independent perturbation theory:

$$E_n^{(2)} = \sum_{n'} \frac{|V_{n'n}|^2}{E_n - E_{n'}} \quad 1 - iE_n^{(2)}(T_a - T_b) = \langle n | T \left\{ e^{-i \int_{T_b}^{T_a} V_I(t) dt} \right\} | n \rangle$$

$$= \langle n | \left(1 - \frac{1}{2} \int_{T_b}^{T_a} \int_{T_b}^{T_a} dt_1 dt_2 T \{ V_I(t_2) V_I(t_1) \} \right) | n \rangle$$

- Can be used in Euclidean space if unphysical $e^{+(E_n - E_{n'})(T_a - T_b)}$ are removed

$K \rightarrow \pi \pi$ Decay

$K \rightarrow \pi \pi$ and CP violation

- Final $\pi\pi$ states can have $I = 0$ or 2 .

$$\langle \pi \pi (I = 2) | H_w | K^0 \rangle = A_2 e^{i\delta_2} \quad \Delta I = 3/2$$

$$\langle \pi \pi (I = 0) | H_w | K^0 \rangle = A_0 e^{i\delta_0} \quad \Delta I = 1/2$$

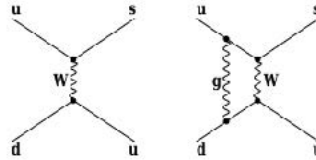
- CP symmetry requires A_0 and A_2 be real.
- Direct CP violation in this decay is characterized by:

$$\epsilon' = \frac{i e^{\delta_2 - \delta_0}}{\sqrt{2}} \left| \frac{A_2}{A_0} \right| \left(\frac{\text{Im} A_2}{\text{Re} A_2} - \frac{\text{Im} A_0}{\text{Re} A_0} \right)$$

Direct CP
violation

Local four quark operators

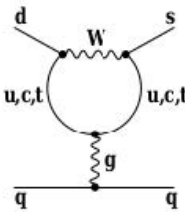
- Current-current operators**



$$Q_1 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} (\bar{u}_\beta u_\beta)_{V-A}$$

$$Q_2 \equiv (\bar{s}_\alpha d_\beta)_{V-A} (\bar{u}_\beta u_\alpha)_{V-A}$$

- QCD Penguins**



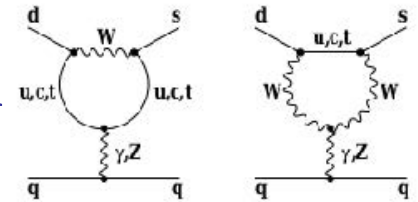
$$Q_3 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_{V-A}$$

$$Q_4 \equiv (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_{V-A}$$

$$Q_5 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_{V+A}$$

$$Q_6 \equiv (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_{V+A}$$

- Electro-Weak Penguins**



$$Q_7 \equiv \frac{3}{2} (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_{V+A}$$

$$Q_8 \equiv \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_{V+A}$$

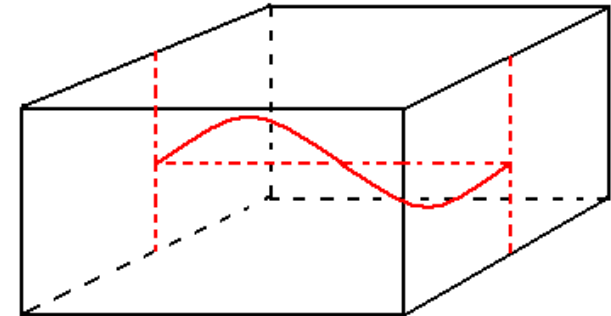
$$Q_9 \equiv \frac{3}{2} (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_{V-A}$$

$$Q_{10} \equiv \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_{V-A}$$

Physical $\pi\pi$ states

Lellouch-Lüscher

- Euclidean $e^{-H_{QCD}t}$ projects onto $|\pi\pi(\vec{p}=0)\rangle$
- Exploit finite-volume quantization.



$$p = 2\pi/L$$

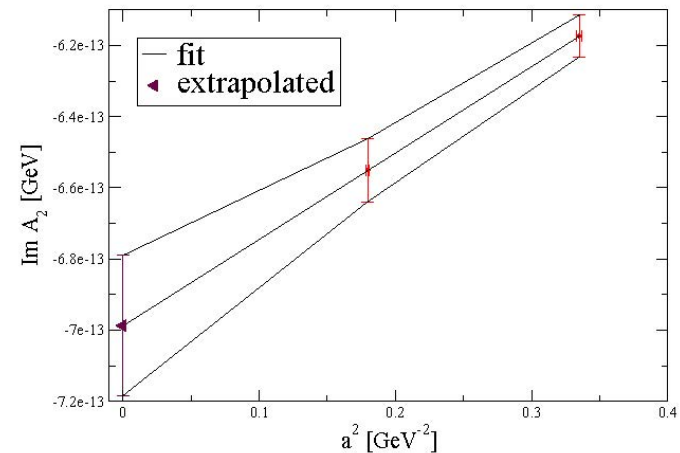
- Boundary conditions give ground state has physical $\vec{p}$
 - $\Delta I = 3/2$: impose anti-periodic BC on d quark
 - $\Delta I = 1/2$: impose G-parity BC
- Correctly include $\pi - \pi$ interactions, including normalization

Calculation of A_2

$\Delta I = 3/2$ – Continuum Results

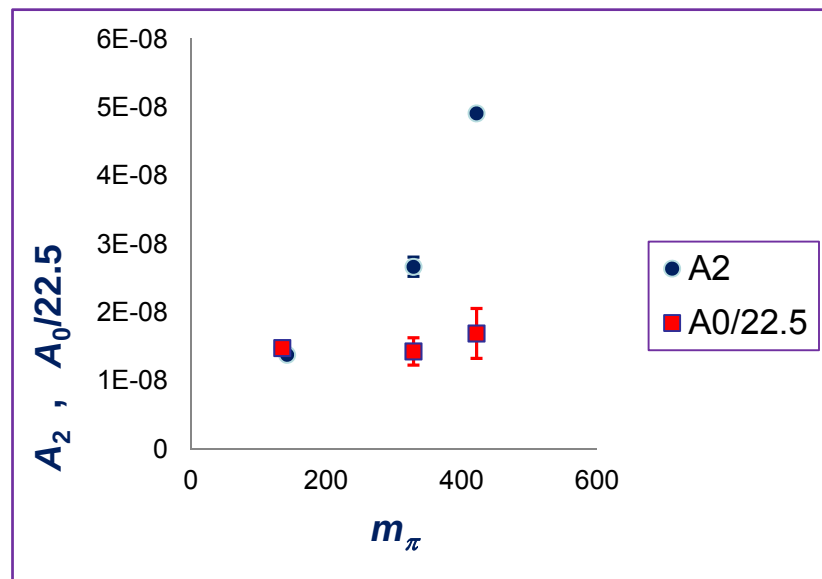
(M. Lightman, E. Goode T. Janowski)

- Use two new large ensembles to remove a^2 error ($m_\rho=135$ MeV, $L=5.4$ fm)
 - $48^3 \times 96$, $1/a=1.73$ GeV
 - $64^3 \times 128$, $1/a=2.28$ GeV
- Continuum results:
 - $\text{Re}(A_2) = 1.50(0.04_{\text{stat}}) (0.14)_{\text{syst}} \times 10^{-8} \text{ GeV}$
 - $\text{Im}(A_2) = -6.99(0.20)_{\text{stat}} (0.84)_{\text{syst}} \times 10^{-13} \text{ GeV}$
- Experiment: $\text{Re}(A_2) = 1.479(4) 10^{-8} \text{ GeV}$
- $E_{\pi\pi} \rightarrow \delta_2 = -11.6(2.5)(1.2)^\circ$
- Phys.Rev. **D91**, 074502 (2015)

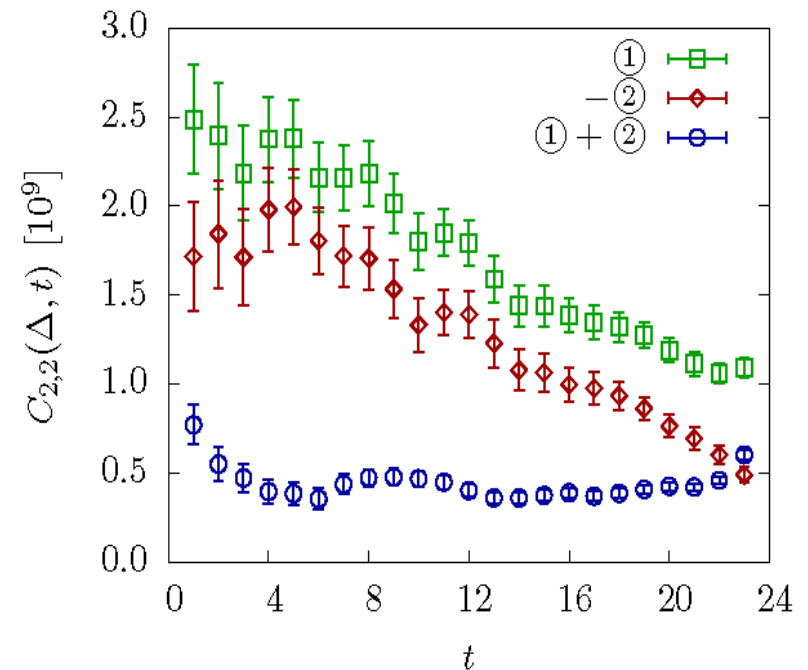


$\Delta I = 1/2$ Rule

Compare A_2 and $A_0/22.5$



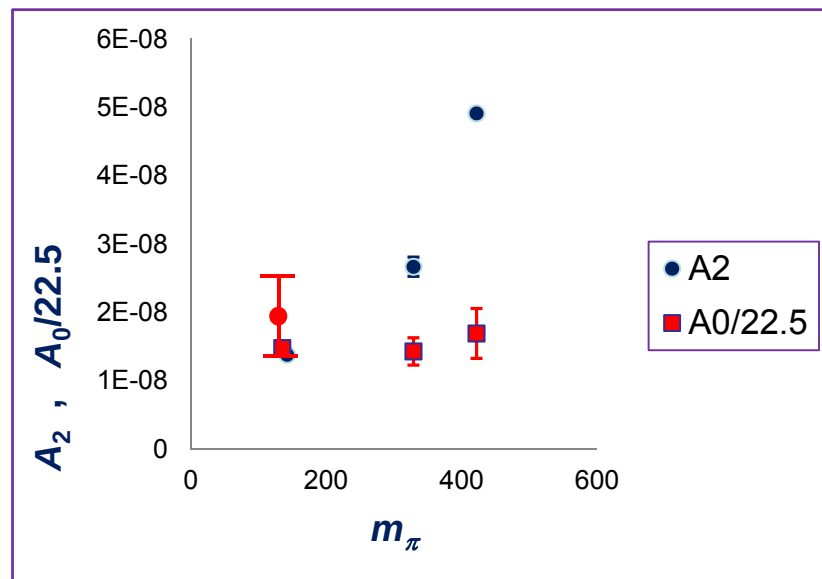
Cancellation in A_2



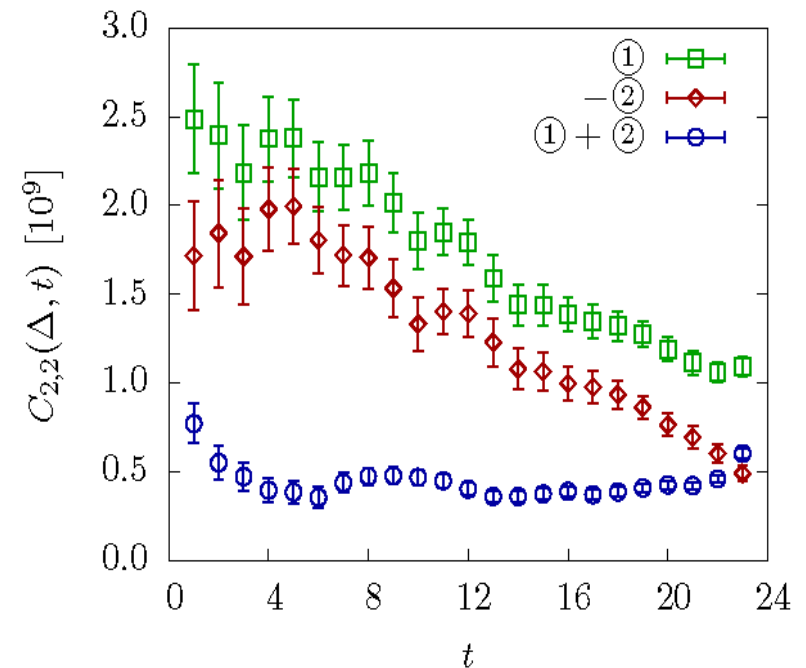
- 50 year puzzle resolved!
- A dynamical QCD effect – no more explanation needed?

$\Delta I = 1/2$ Rule

Compare A_2 and $A_0/22.5$



Cancellation in A_2



- 50 year puzzle resolved!
- A dynamical QCD effect – no more explanation needed?

Calculation of A_0 and ε'

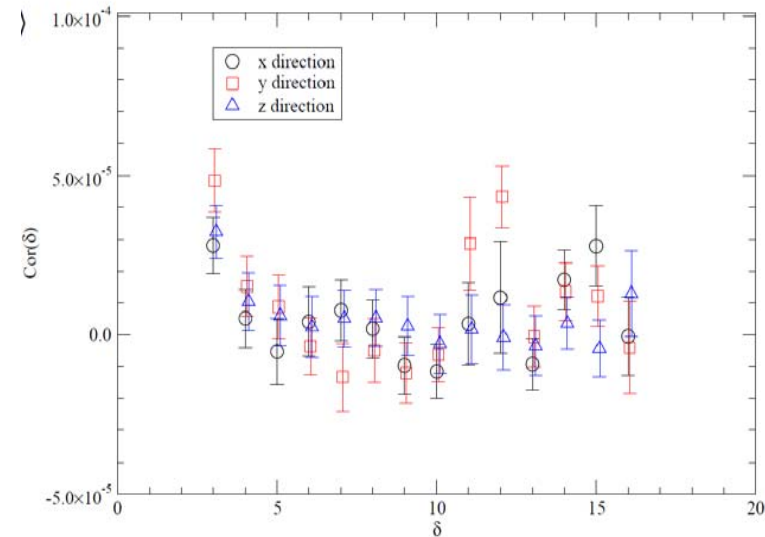
Overview of calculation

(Chris Kelly, Daiqian Zhang)

- Use $32^3 \times 64$ ensemble
 - $1/a = 1.3784(68)$ GeV, $L = 4.53$ fm.
 - 216 configurations separated by 4 time units
 - 900 low modes for all-to-all propagators
 - Solve for $\pi\pi$ and kaon sources on each of 64 time slices

Overview of calculation

- Achieve essentially physical kinematics:
 - $M_\pi = 143.1(2.0)$
 - $M_K = 490.6(2.2)$ MeV
 - $E_{\pi\pi} = 498(11)$ MeV
 - $m_{res} = 0.001842(7)$
- Error in ensemble generation
(u and d quark forces computed from the same random numbers after shift by 12 in y-direction)



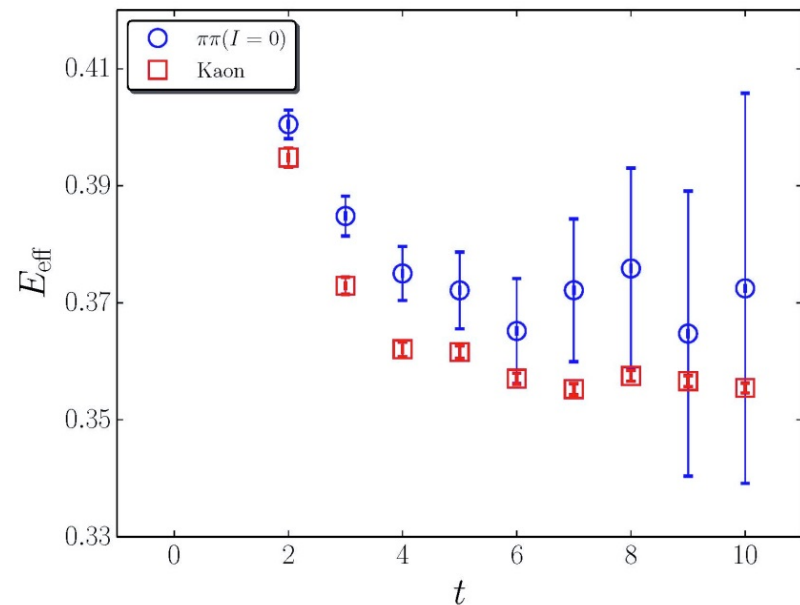
Average plaquette

Correct ensemble 0.512239(3)(7)

Incorrect ensemble 0.512239(6)

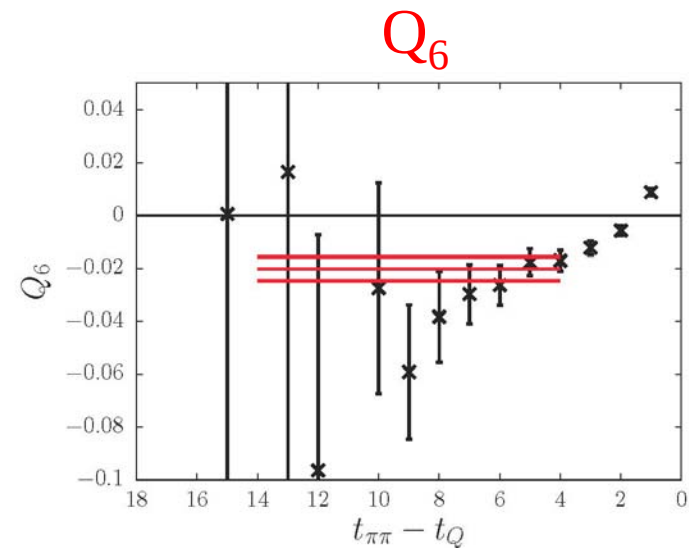
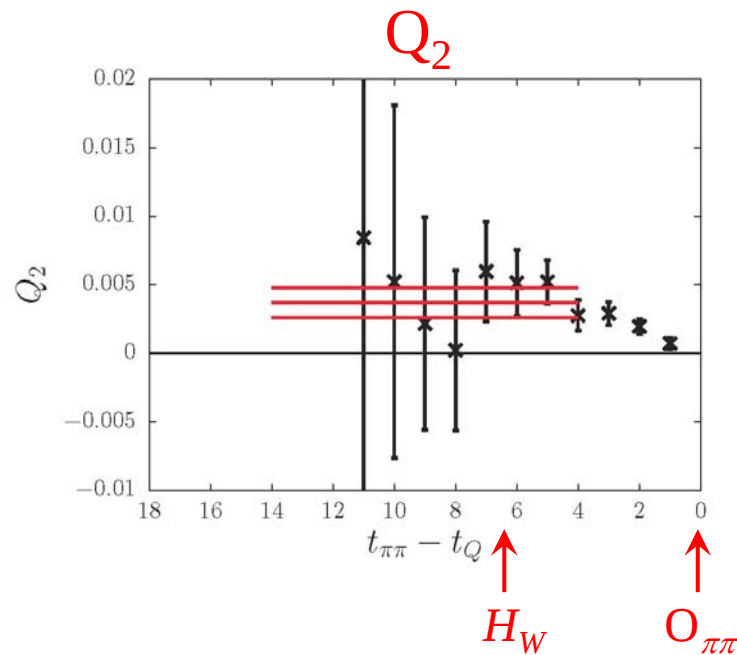
$I = 0$, $\pi\pi - \pi\pi$ correlator

- Determine normalization of $\pi\pi$ interpolating operator
- Determine energy of finite volume, $I = 0$, $\pi\pi$ state:
 $E_{\pi\pi} = 498(11) \text{ MeV}$
- Determine $I = 0$ $\pi\pi$ phase shift: $\delta_0 = 23.8(4.9)(2.2)^\circ$
- Phenomenological result:
 $\delta_0 = 38.0(1.3)^\circ$ [G. Colangelo]



$\Delta I = 1/2 \quad K \rightarrow \pi \pi$ matrix elements

- Vary time separation between H_W and $\pi\pi$ operator.
- Show data for all $K - H_W$ separations $t_Q - t_K \geq 6$ and $t_{\pi\pi} - t_K = 10, 12, 14, 16$ and 18.
- Fit correlators with $t_{\pi\pi} - t_Q \geq 4$
- Obtain consistent results for $t_{\pi\pi} - t_Q \geq 3$ or 5



Results

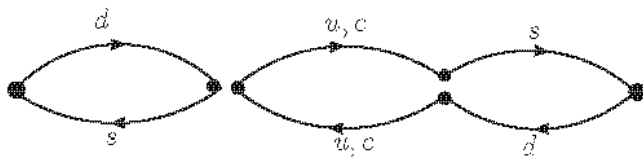
- Determine the complex $\Delta I=1/2$ amplitude A_0
 - $\text{Re}(A_0) = (4.66 \pm 1.00_{\text{stat}} \pm 1.26_{\text{sys}}) \times 10^{-7} \text{ GeV}$
 - Expt: $(3.3201 \pm 0.0018) \times 10^{-7} \text{ GeV}$
 - $\text{Im}(A_0) = (-1.90 \pm 1.23_{\text{stat}} \pm 1.08_{\text{sys}}) \times 10^{-11} \text{ GeV}$
- Calculate $\text{Re}(\varepsilon'/\varepsilon)$:
- $\text{Re}(\varepsilon'/\varepsilon) = (1.38 \pm 5.15_{\text{stat}} \pm 4.59_{\text{sys}}) \times 10^{-4}$
 - Expt.: $(16.6 \pm 2.3) \times 10^{-4}$
 - 2.1σ difference

$K_L - K_S$ mass difference

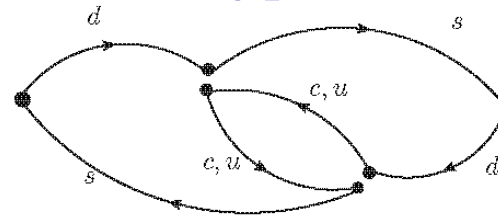
$K_L - K_S$ mass difference

- Perturbative result integrates out charm and shows poor convergence (Brod and Gorbahn)
 - NNLO is 36% of LO
 - Large μ_c dependence
- Lattice must include charm quark (GIM)

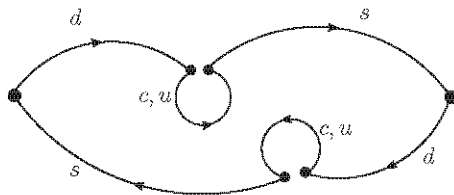
Type 1



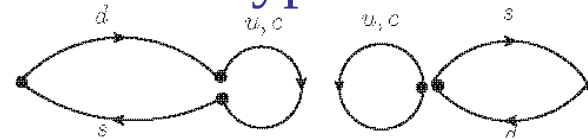
Type 2



Type 3

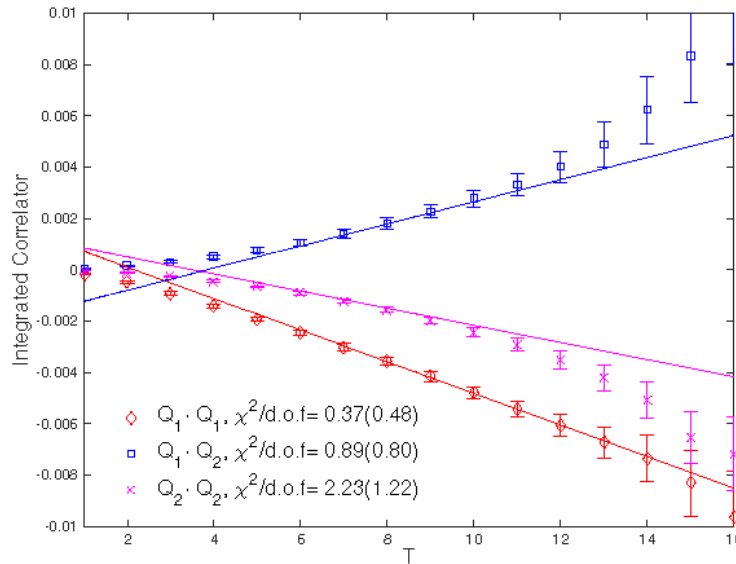


Type 4



ΔM_K Present Results

(Ziyuan Bai)



	$\Delta M_K \times 10^{+12} \text{ MeV}$
Types 1-4	5.76(73)
Types 1-2	4.19(15)
η	0
π	0.27(14)
$\pi\pi, I=0$	-0.097(49)
$\pi\pi, I=2$	$-6.56(6) \times 10^{-4}$
Δ_{FV}	0.029(19)
Expt.	3.483(6)

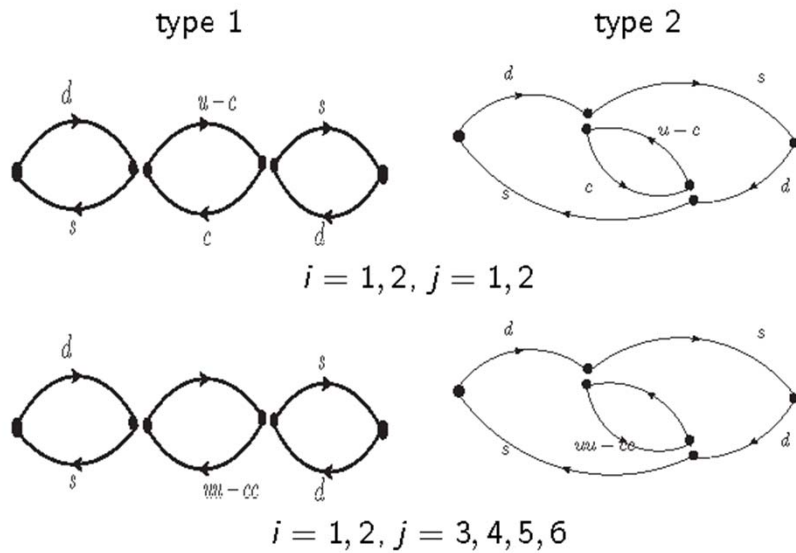
- $m_c = 750 \text{ MeV}$, $M_\pi = 170 \text{ MeV}$
- Disconnected contribution small
- $\pi\pi$ contribution $\sim 2\%$ and FV correction $\sim 0.5\%$
- New $64^3 \times 128$, $1/a = 2.38 \text{ GeV}$, $m_c = 1.2 \text{ GeV}$, $M_\pi = 140 \text{ MeV}$
26 configs: $\Delta M_K = 0.6(3.3) \times 10^{-12} \text{ MeV}$ (26 $\rightarrow$ 400?)

**Long distance
part of ε_K**

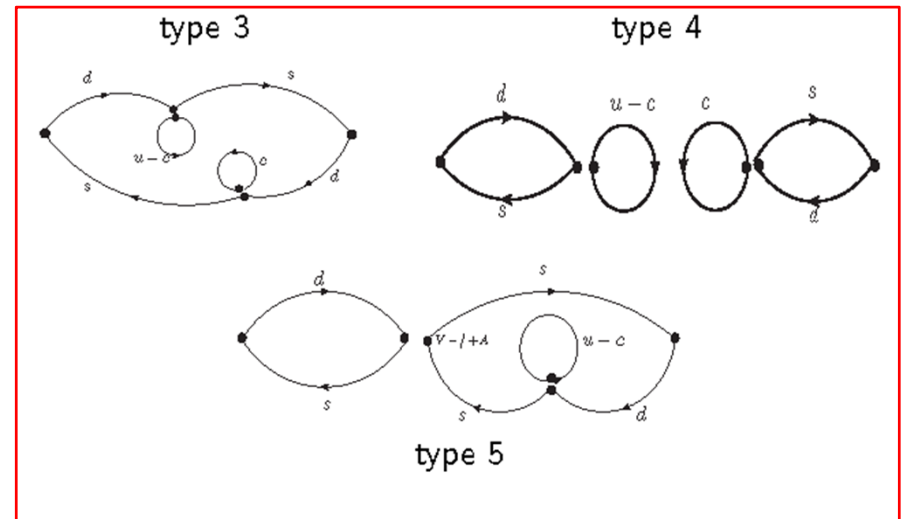
Diagrams for $\lambda_t \lambda_u$ contribution to ε_K

(Ziyuan Bai)

- Identify five types of diagrams

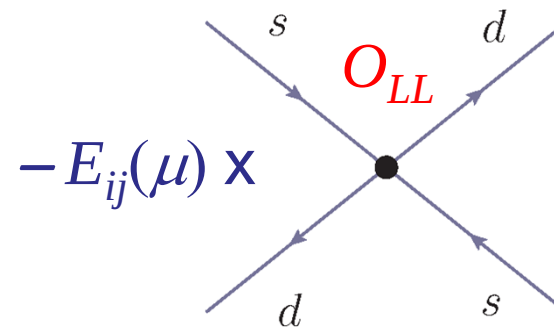
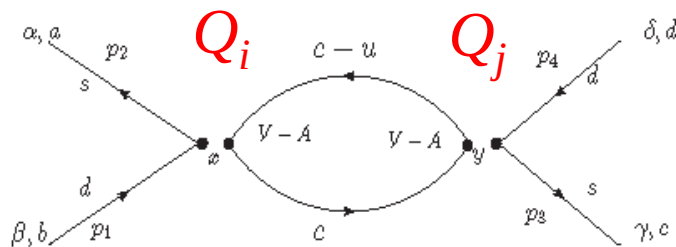


Omit from 1st study



New $\Delta S = 2$ counter term

(Ziyuan Bai)



- Subtract $E_{ij}(\mu) (\bar{s}\gamma^\nu(1-\gamma^5)d) (\bar{s}\gamma^\nu(1-\gamma^5)d)$ to make off-shell Greens function vanish at $p_i^2 = \mu_{RI}^2$
- Define infrared-safe Rome-Southampton normalization for bi-local operator.

Progress toward long-distance part of ε_K

(Ziyuan Bai)

- Examine only type 1 and 2 diagrams
- Only current-current operators
- Compute NLO (one-loop) conversion from bilocal RI to $\overline{\text{MS}}$
- Dependence on μ_{RI}
- Preliminary
- $|\varepsilon_K| = 2.228(11) \times 10^{-3}$ expt.

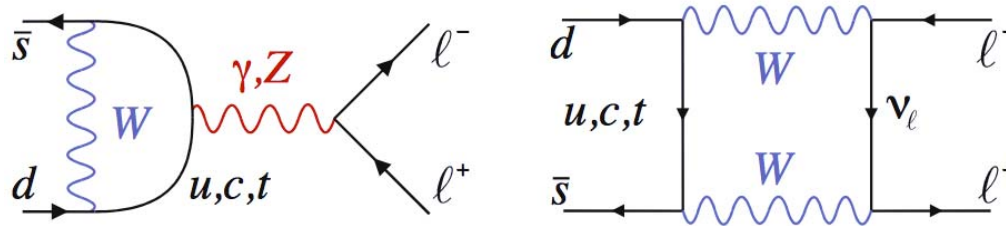
μ_{RI}	$\delta \varepsilon_K $
1.54	0.1384×10^{-3}
1.92	0.1483×10^{-3}
2.11	0.1473×10^{-3}
2.31	0.1405×10^{-3}
2.56	0.1246×10^{-3}

Rare Kaon Decays

$$K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

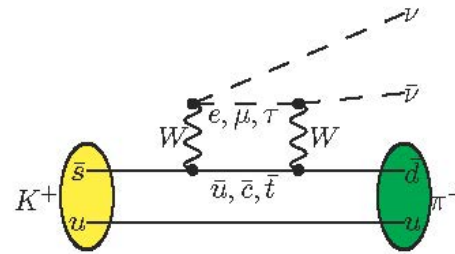
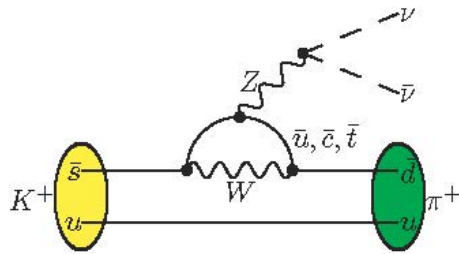
Rare Kaon Decays

(Xu Feng, Andrew Lawson, Antonin Portelli)



- $K_L \rightarrow \pi^0 + l^+ + \bar{l}^-$: determine the sign of the indirect CP violating amplitude.
- $K^+ \rightarrow \pi^+ + \nu + \bar{\nu}$: calculate the long distance ($l \geq 1/m_c$) part of charm contribution. Estimated to be small ($\approx 4\%$) but should be verified

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ (Xu Feng)



Both are bilocal products with short-distance singularity

- Estimate 3 contributions: top : charm- sd : charm- ld
[Cirigliano et.al. Rev. Mod. Phys.]

$$\lambda_t \frac{m_t^2}{M_W^2} : \lambda_c \frac{m_c^2}{M_W^2} \ln \frac{M_W}{m_c} : \lambda_u \frac{\Lambda_{\text{QCD}}^2}{M_W^2} = 68\% : 29\% : 3\%$$

- However
 - Charm contributes 50% of $K^+ \rightarrow \pi^+ + \nu \bar{\nu}$ BR
 - 50% charm contribution comes from $p \sim m_c$

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ - results

(Xu Feng)

- Present phenomenological treatment [Buras, et al. hep-ph/1503.02693]
 - Integrate out charm & represent by local $(\bar{s}d)_L(\bar{\nu}\nu)_L$ operator
 - Add correction for up quark
 - Add $(\Lambda_{\text{QCD}}/m_c)^2$ correction for dim-8 in OPE
 - $\delta P_{cu} = 0.04(2)$ [Isidori et.al, hep-ph/0503107]
- Lattice result
 - $\delta P_{cu} = -0.007(2)$ (large systematic errors)
 - Includes all physics at energies below $\mu_{RI} = 2 \text{ GeV}$

Outlook

- Lattice QCD now reaches sufficient far into Minkowski space to allow 1st-principles calculation of:
 - $K \rightarrow \pi \pi$, $\Delta I = 3/2$ and $1/2$, ε'/ε
 - $M_{K_L} - M_{K_S}$, long dist. contribution to ε
 - Long distance parts of $K \rightarrow \pi l \bar{l}$, $K \rightarrow \pi \nu \bar{\nu}$
- First realistic calculation of ΔM_K underway
- Must wait for next generation of computers to accurately include charm