

# Naturalness of Electroweak Symmetry Breaking in the LHC Era

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- Dark matter
- Cosmological baryon asymmetry
- Neutrino oscillation
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- Dark matter
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But none of these necessarily points to LHC scales, few $\times$ 100 GeV – few TeV.

What points to LHC is **naturalness problem**

(aka. hierarchy problem,  
fine tuning " )

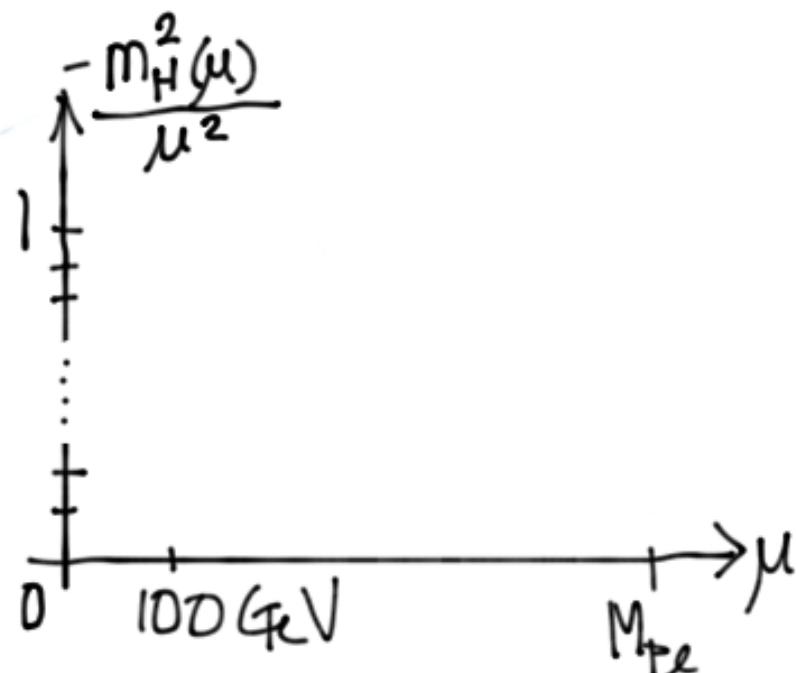
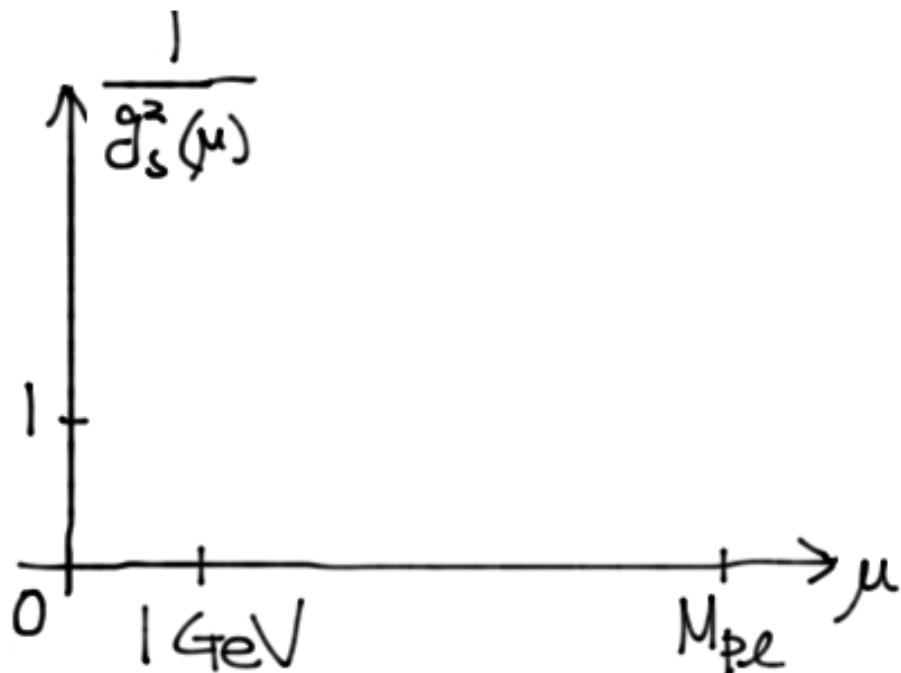
# Naturalness Problem of EW scale

Two scales in SM:

$$\Lambda_{\text{QCD}} \sim 1 \text{ GeV} \ll M_{\text{Pl}} \quad m_H^2 \sim -(100 \text{ GeV})^2 \ll M_{\text{Pl}}^2$$

$$\mathcal{L} = \frac{1}{g_s^2} G_{\mu\nu} G_{\mu\nu} + \dots$$

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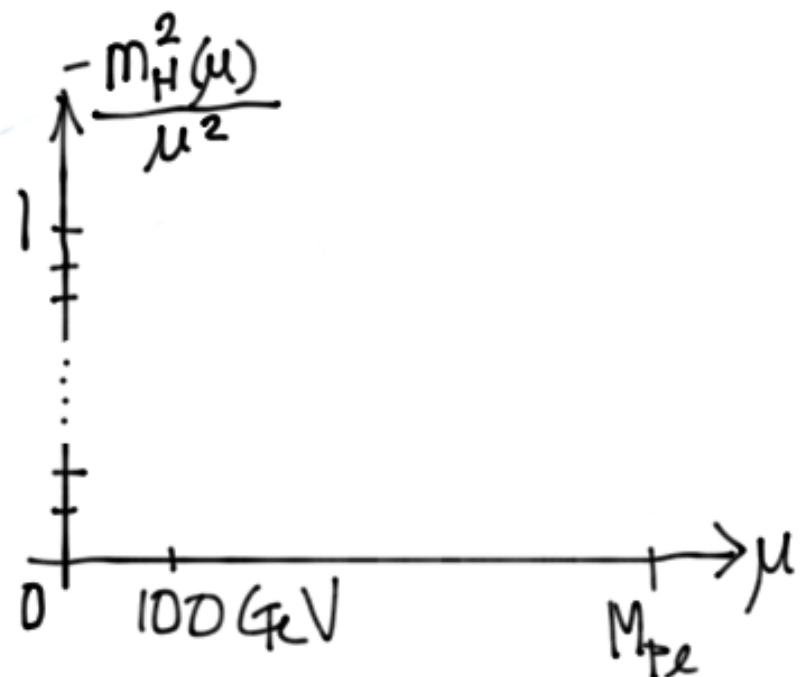
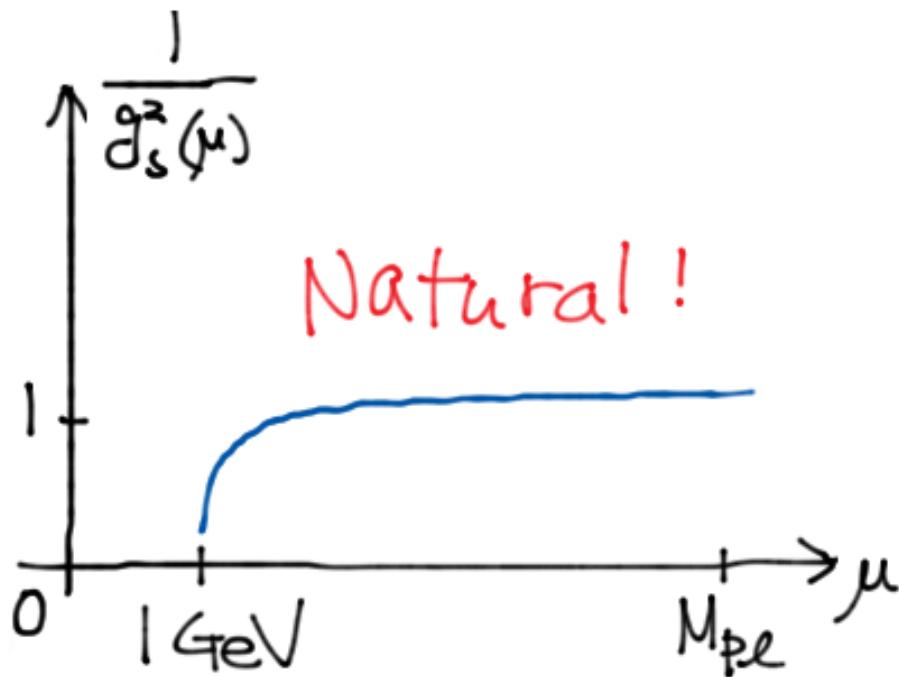
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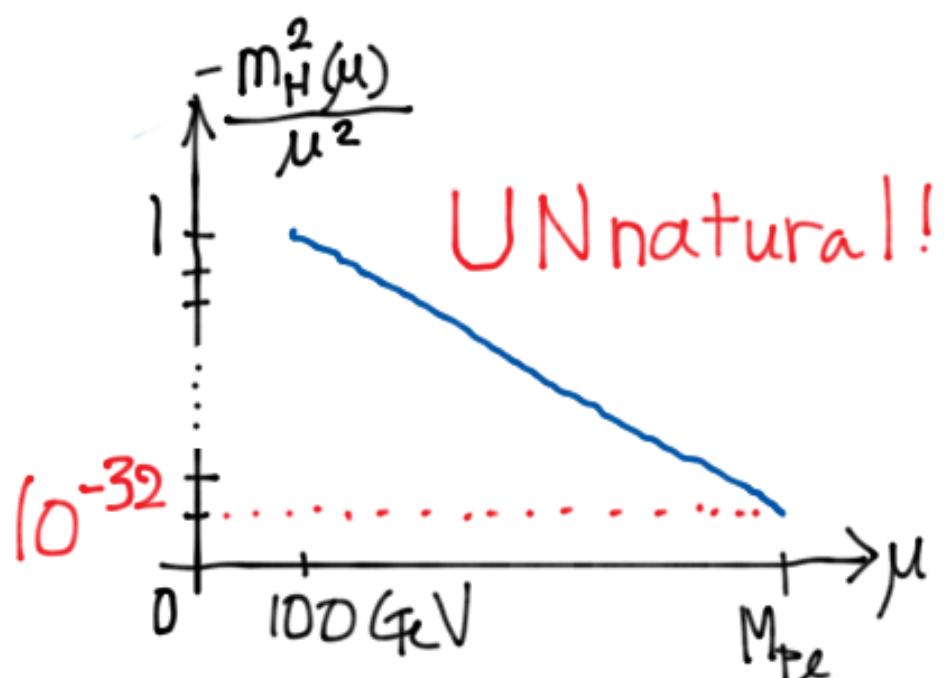
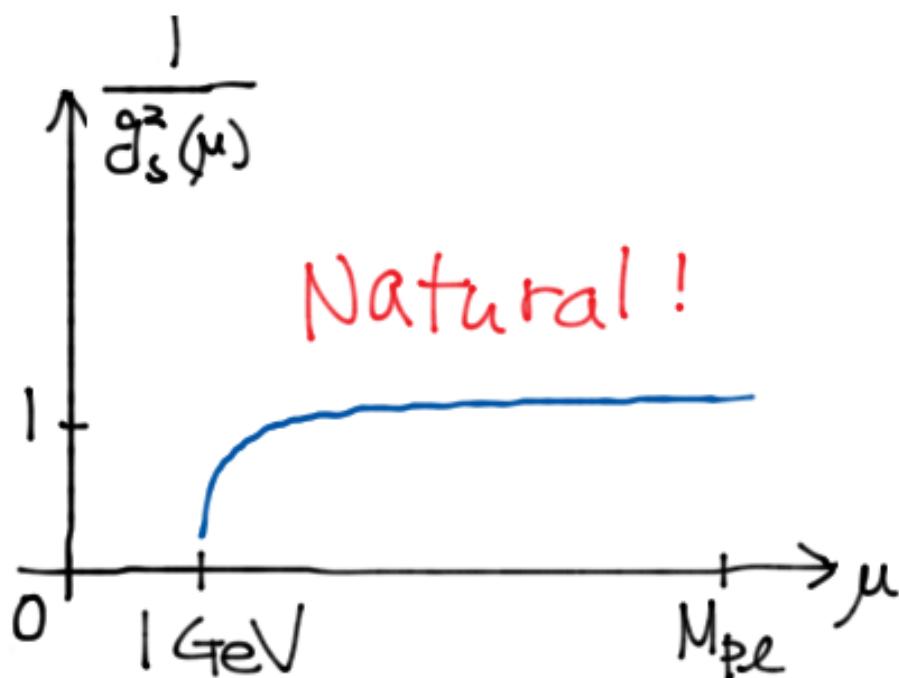
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$$\mathcal{L} = \frac{m_H^2}{2} H^\dagger H + \dots$$

dim.  $\simeq 2$



Heart of problem:

(A)  $\dim[H^\dagger H] = 2 = 4 - 2$

(B)  $H^\dagger H$  allowed by  
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Quadratically  
Sensitive to UV!

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$$m_{H,\text{UV}}^2 H^\dagger H + \Phi^\dagger \Phi H^\dagger H$$

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$$\langle \Phi \rangle = \Lambda \downarrow$$

$$m_{H,IR}^2 = \Lambda^2 + m_{H,UV}^2$$

Must be tuned  
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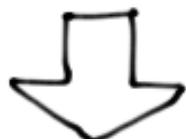
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(B)  $H^\dagger H$  allowed by all symmetries of SM  $\Rightarrow$  Quadratically sensitive to UV!

(e.g.)

$$m_{H,UV}^2 H^\dagger H + \Phi^\dagger \Phi H^\dagger H$$



Solutions: Negate (A) or (B)!

(A)  ~~$H$~~   $\rightarrow D_H$

with  $\underline{[D_H^\dagger D_H]} > 4$

(e.g.)

$$D_H = \bar{\Psi}_T \Psi_T \text{ "Technicolor"}$$

$$\langle \Phi \rangle = \Lambda$$

$$m_{H,IR}^2 = \Lambda^2 + m_{H,UV}^2$$

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(B) Non symmetry to forbid  $H^\dagger H$

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  - $m_h \ll m_{\text{others in Higgs sector}}$   
(e.g. MSSM  $\rightarrow m_h \ll m_{H^0, H^\pm, A^0}$ )

$\Rightarrow$  BSM Higgs  $\simeq$  SM Higgs unless  
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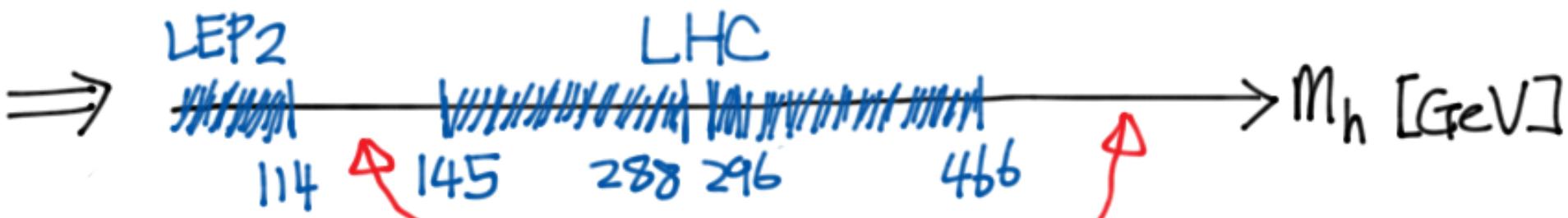


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Precisely Light Higgs (B) or No Higgs (A) !

# Naturally Light Higgs

## Plan

(i) Extend SM by a new symm. to kill  $H^+H^-$ ,

(ii) Break the new symm. in a natural way.

Typically a separate "module"  
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Focus on (i).

# Naturally Light Higgs

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Typically a separate "module"  
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Focus on (i).

Break down (i) in 2 steps:

(i-A) Find a symm. that gives some kind of massless particle  $X$ ,

(i-B) Relate  $X$  to  $H$ .

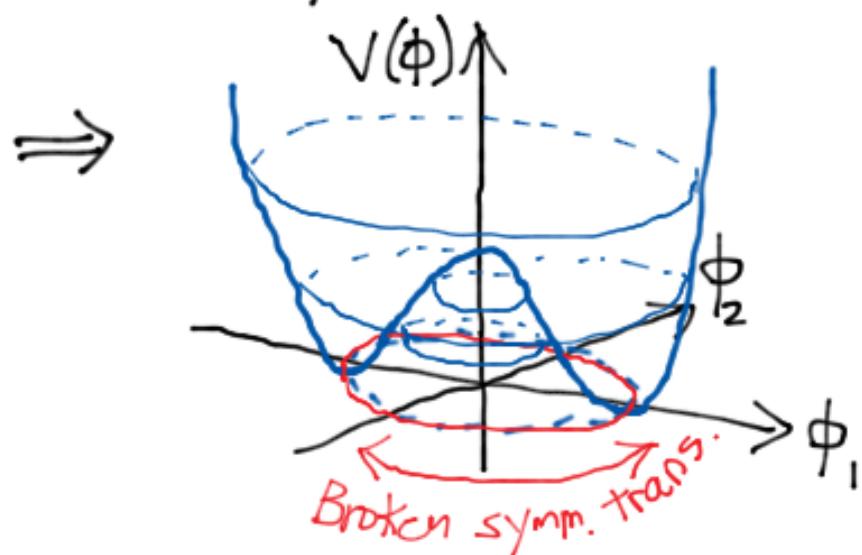
# Massless particles in QFT

spin-0 Nambu-Goldstone boson (NGB)

Unbroken symm.  $\equiv$  Rotates particles

$\Rightarrow m_{\text{particle 1}} = m_{\text{particle 2}}, \text{ etc.}$

Spontaneously broken symm.  $\equiv$  Rotates vacua



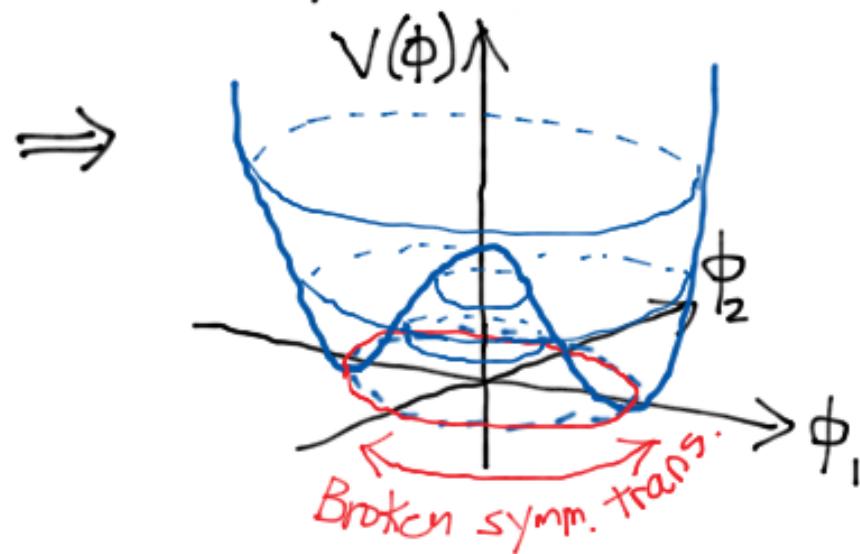
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Excitations along    
= Massless!

$H = \text{NGB of a new broken symm} \Rightarrow m_H = 0 !$

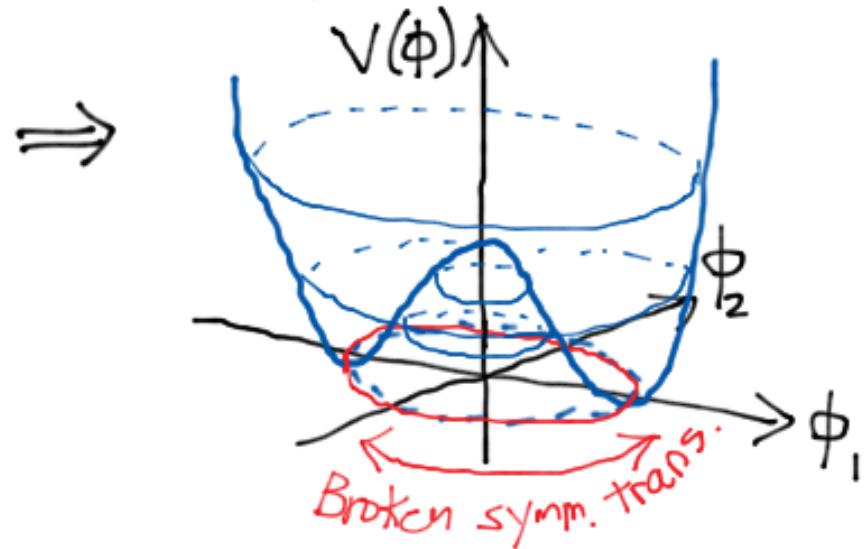
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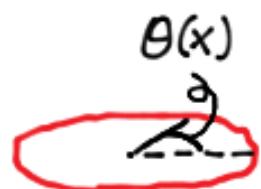
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Excitations along   
 $\theta(x)$   
= Massless!

"Shift symm."  
 $\theta(x) \rightarrow \theta(x) + \epsilon$

$H = \text{NGB of a new broken symm} \Rightarrow m_H = 0 !$

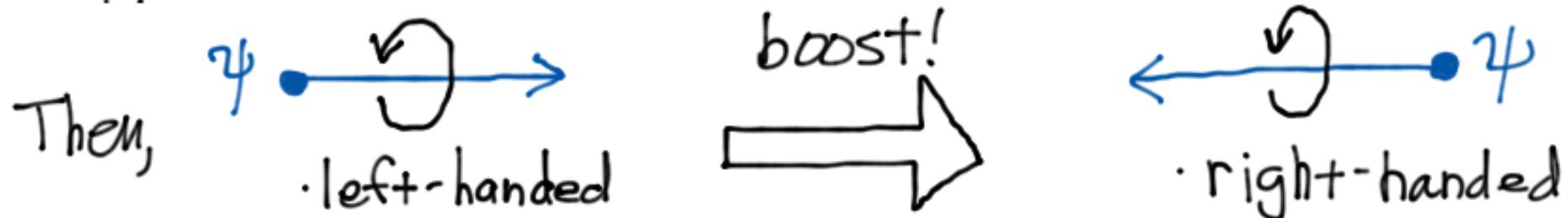
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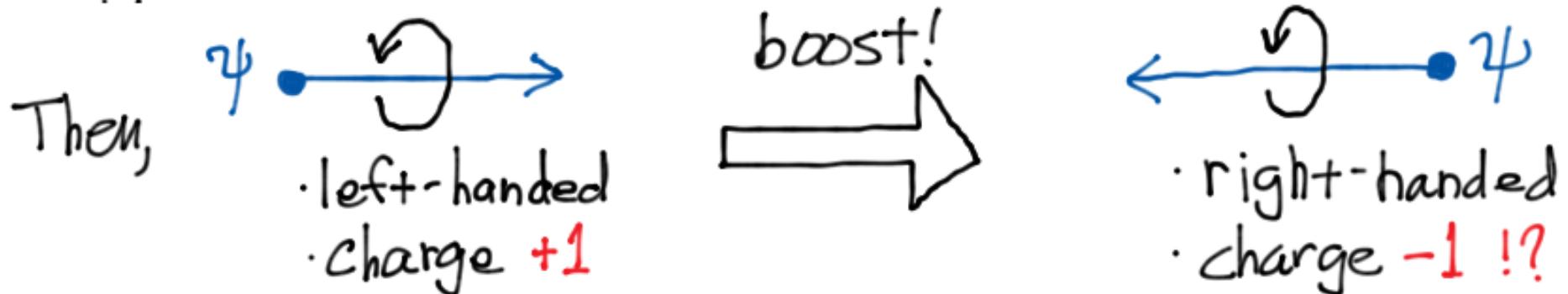
Suppose  $m_\psi \neq 0$ .



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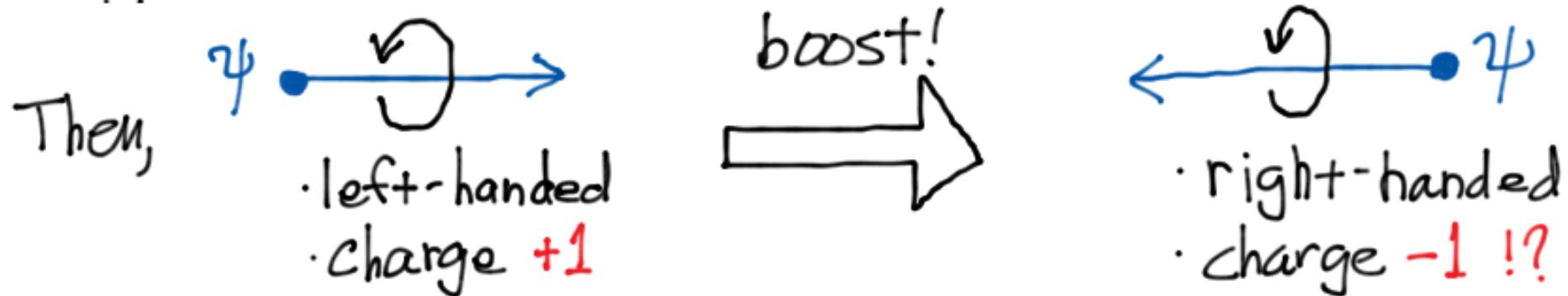
By contradiction,  $m_\psi = 0 !$

Boost shouldn't change charge!

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By contradiction,  $m_\psi = 0 !$

Must relate  $\Psi$  ( $S=\frac{1}{2}$ ) to  $H$  ( $S=0$ ) to get  $m_H=0$ .  
 $\Rightarrow$  Supersymmetry !

spin-1

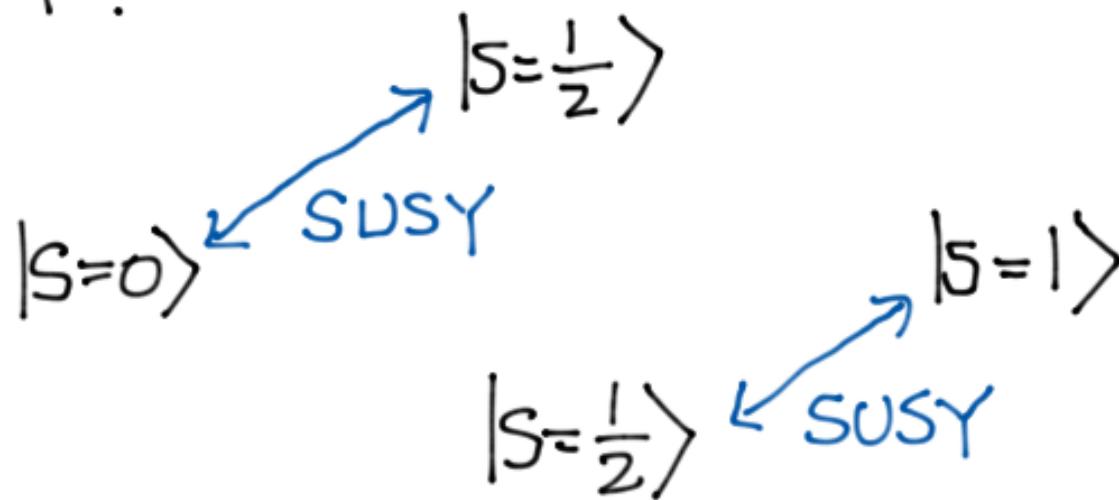
Gauge symm  $A_\mu \rightarrow A_\mu + \partial_\mu \epsilon$   $\Rightarrow$   ~~$m_A^2 A_\mu A^\mu$~~   
 $m_A = 0!$

## spin-1

Gauge symm  $A_\mu \rightarrow A_\mu + \partial_\mu \epsilon$   $\Rightarrow \cancel{m_A^z A_\mu A^\mu}$   
 $m_A = 0!$

Must relate spin-1 to spin-0 for  $M_H = 0$ .

SUSY?



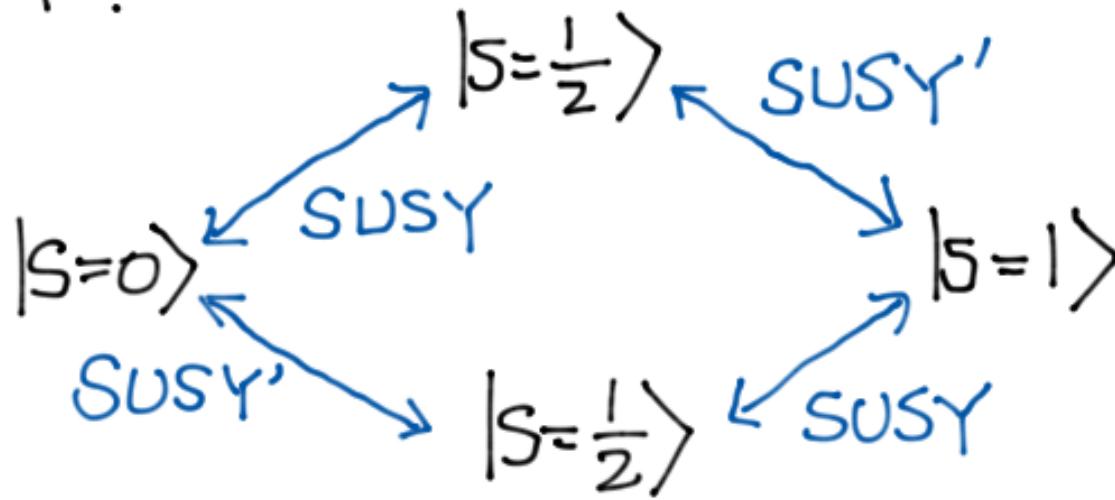
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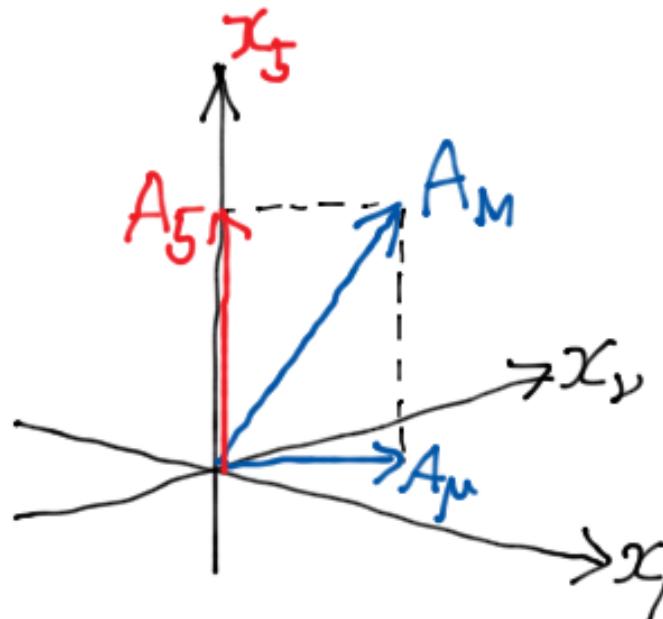
SUSY?



"N=2 SUSY"  
Not chiral  
 $(SUSY = SUSY_L)$   
 $(SUSY' = SUSY_R)$

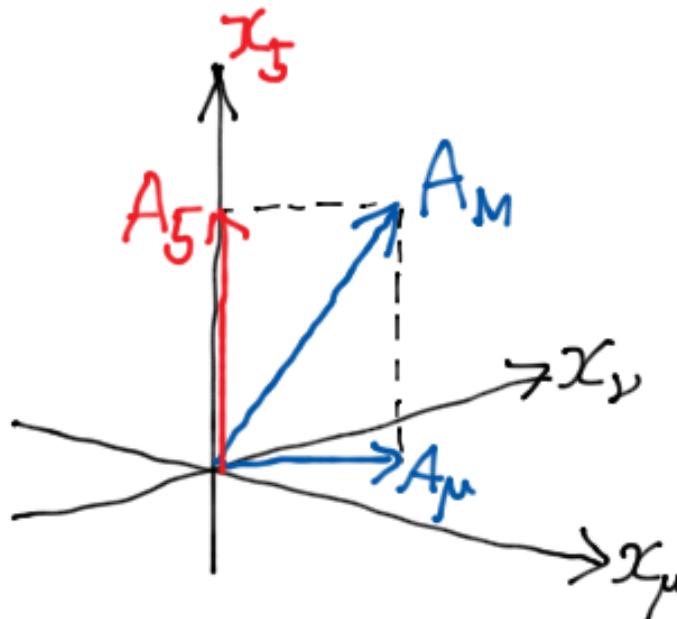
No!

Need an extra dimension!



$$(\mu, \nu = 0, 1, 2, 3; M = \mu, 5)$$

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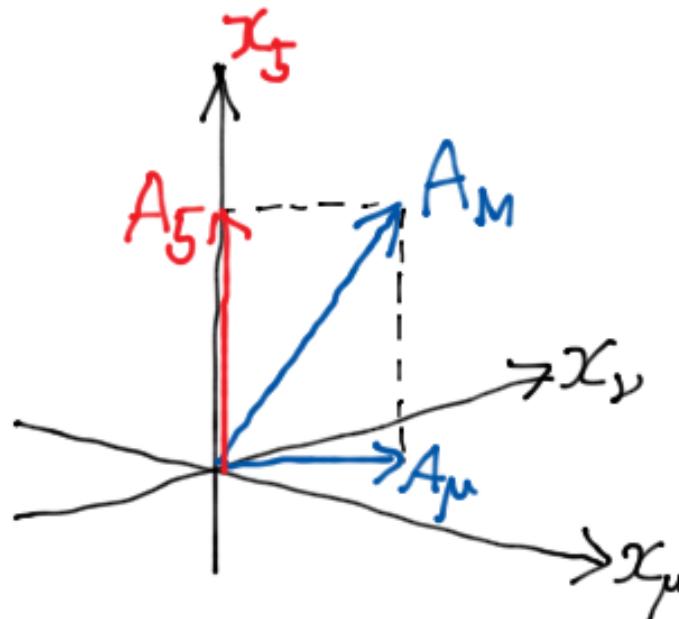
$$(\mu, \nu = 0, 1, 2, 3; M = \mu, 5)$$

$A_5$  ... unchanged  
under  $x_\mu - x_\nu$   
rotations

|||  
Scalar in 4d !

$$H = A_5 \Rightarrow m_H = 0 !$$

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$$H = A_5 \Rightarrow m_H = 0 !$$

But

5d gauge trans:  $A_5 \rightarrow A_5 + \partial_5 \epsilon \Leftrightarrow H \text{ shifts!}$

Effectively same as  $H = NGB$  !

Spin-3/2

(No known working model)

Need SUSY (& extra dim.) to relate to H.

(e.g.)  $S=0 \xleftrightarrow{\text{SUSY}} S=\frac{1}{2} \xleftrightarrow{\times D} S=\frac{3}{2}$

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Spin-2

(No known working model)

$m=0$

"Gauge" symm:  $h_{\mu\nu} \rightarrow h_{\mu\nu} + \partial_\mu \varepsilon_\nu + \partial_\nu \varepsilon_\mu \Rightarrow$

(or GR!)

$\cancel{h_{\mu\nu}}$   $\cancel{h_{\nu}^{\mu}}$   
 $\cancel{h_{\mu}^{\nu}}$   $\cancel{h_{\nu}^{\nu}}$

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$$\begin{array}{c} \cancel{h_{\mu\nu}} \\ \cancel{h_\mu} \cancel{h_\nu} \end{array}$$

Need extra dimensions to get scalar

$h_{55}, h_{56}, \text{etc.} = \text{scalars in 4d} \sim H$

But H shifts under symm.

Same story as spin-1     $H=NGB$  !

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Spin > 2

No interacting massless particles.

## Summary so far

Only 2 ways to have  $m_H = 0$  naturally !

**SUSY**

$$H \xleftrightarrow[\text{(S=0)}]{}^{\text{SUSY}} \tilde{H} \quad (\text{S}=\frac{1}{2})$$

$$m_H = 0 \xleftrightarrow[\text{SUSY}]{} m_{\tilde{H}} = 0$$

↑  
chiral  
symm.

**$H = NGB$**

Shift symm.

$$H \rightarrow H + \epsilon$$

$$\Rightarrow \cancel{m_H^2 H^\dagger H}$$

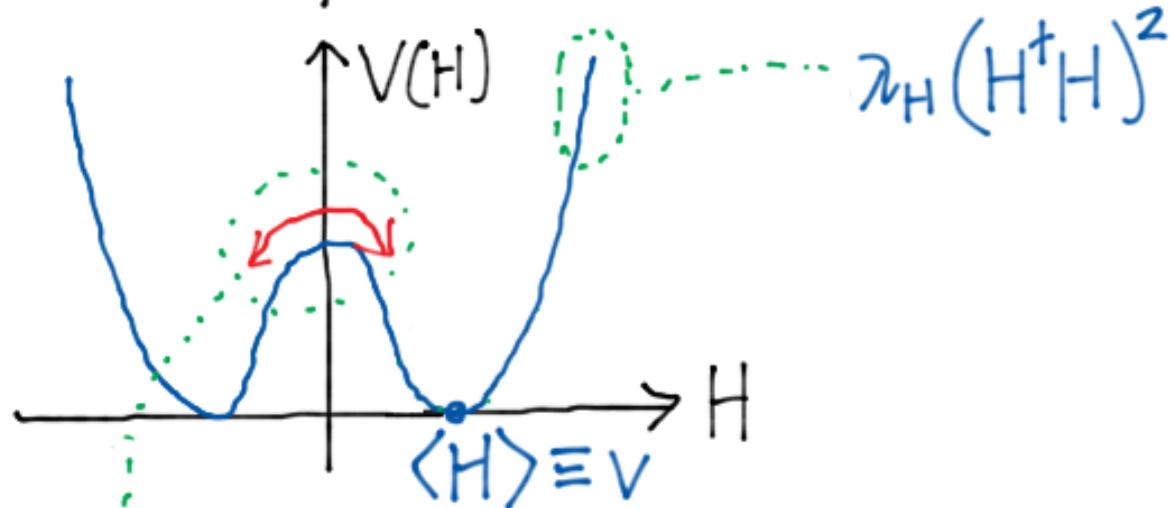
e.g. Composite Higgs  
Little Higgs  
Twin Higgs  
Higgs as  $A_5$

# EWSB and Higgs mass

Break the  $m_H=0$  symm. to get  $m_H^2 < 0$  for EWSB.

We want

$$V(H) =$$



$$m_H^2 \sim - (100 \text{ GeV})^2$$

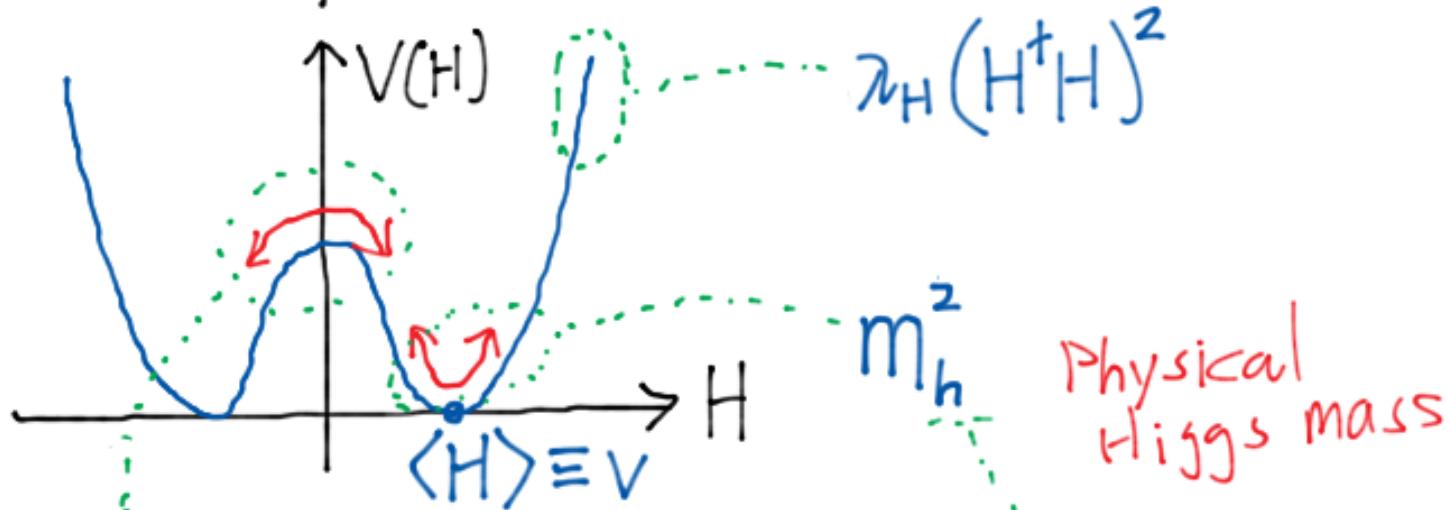
$$\sim \lambda_H (H^+ H - v^2)^2$$

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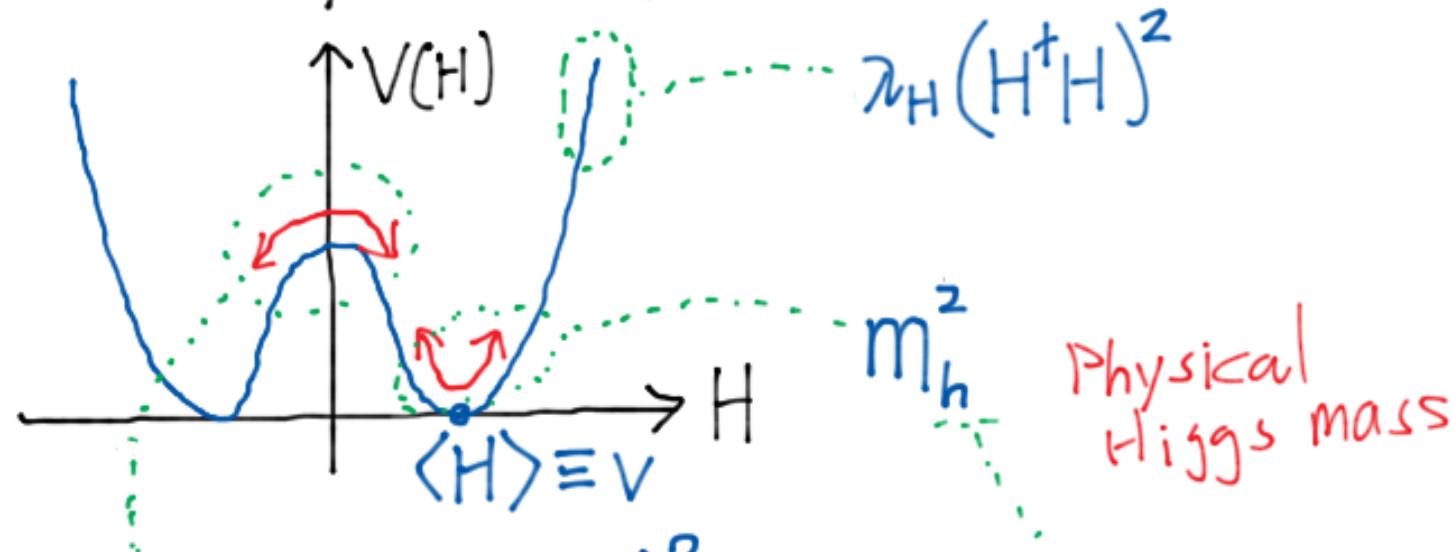
Physical Higgs

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$$\sim \lambda_H (H^\dagger H - v^2)^2$$

$$\sim \lambda_H v^2 h^2 + \dots$$

$$H = v + h \quad \text{Physical Higgs}$$

$$\Rightarrow \underline{m_h^2 \sim \lambda_H v^2}$$

$$v = (\text{measured}) \approx 174 \text{ GeV}$$

What's  $m_h$ ? = What's  $\lambda_H$ ?

# $\lambda_H$ in SUSY

(i) SUSY contributions

 Kills  $H^\dagger H$

(ii) SUSY contributions

# $\lambda_H$ in SUSY

(i) SUSY contributions

→ Kills  $H^\dagger H$

and predicts  $\lambda_H \sim g_1^2 + g_2^2$  !

$$\Rightarrow m_h^2 \sim \lambda_H V^2 \sim m_Z^2$$

(More precisely  $m_h \leq m_Z$ )

(ii) ~~SUSY~~ contributions

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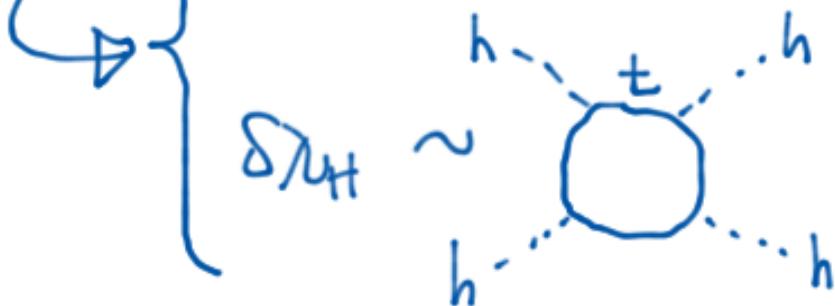
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## (ii) SUSY contributions

$$\left\{ \begin{array}{l} \delta m_H^2 \sim -(100 \text{ GeV})^2 \rightarrow \text{EWSB} \\ \delta \lambda_H \sim \dots \end{array} \right.$$



$$\sim \frac{3 y_t^4}{4 \pi^2} \log \frac{m_{\tilde{t}}^2}{m_t^2}$$

Scalar  
top mass

# $\lambda_H$ in SUSY

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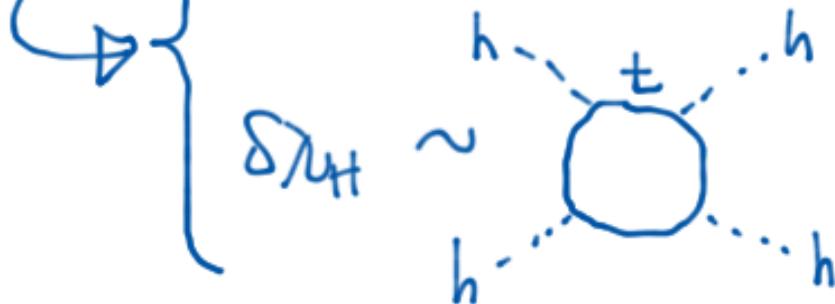
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$$\sim \frac{3 y_t^4}{4 \pi^2} \log \frac{m_{\tilde{t}}^2}{m_t^2}$$

scalar  
top mass  
||  
cutoff  
for SM!

Naturalness  $\Rightarrow m_{\tilde{t}} \lesssim 1 \text{ TeV}$

$$\Rightarrow \underline{m_h^2 \lesssim m_Z^2 + (90 \text{ GeV})^2 \simeq (130 \text{ GeV})^2}$$

# $\lambda_H$ in $H=NGB$

(i) NGB contributions

6  $H \rightarrow H + \varepsilon$  symm.  $\Rightarrow V(H) = \underline{\text{const.}}$

No  $H^\dagger H$   
No  $(H^\dagger H)^2$ !

(ii)  $H \rightarrow H + \varepsilon$  violating contributions

# $\lambda_H$ in $H=NGB$

## (i) NGB contributions

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## (ii) $H \rightarrow H + \varepsilon$ violating contributions

$\sum$  Largest violation = top Yukawa  $\frac{y_t}{\tilde{O}(1)} H \bar{Q}_{3L} t_R$

$$\Rightarrow \delta \lambda_H \sim \text{Diagram} \sim \frac{3 y_t^4}{4 \pi^2} \left( \log \frac{\Lambda_t}{m_t} + \tilde{O}(1) \right)$$

The Feynman diagram shows a central loop of three higgs bosons (h) meeting at a central top quark loop (t). The higgs bosons are represented by dashed lines, and the top quark loop is represented by a solid circle.

$$\Lambda_t = M_t, M_{t \text{ comp}}, \text{etc.}$$

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$$\Rightarrow \delta \lambda_H \sim \text{loop diagram with } h, t, h \sim \frac{3 y_t^4}{4 \pi^2} \left( \log \frac{\Lambda_t}{m_t} + \tilde{O}(1) \right)$$

$$\Lambda_t = M_t, M_{t \text{ comp}}, \text{etc.}$$

$$\Rightarrow \underline{m_h^2 \lesssim \tilde{O}(1) \cdot (90 \text{ GeV})^2}$$

At most  $\sim 100-200 \text{ GeV}$   
like SUSY case!

# Summary

